

Molecular Dynamics: basic methods

Pascal Viot

September 20, 2020

Simple Liquids

- m mass of the particle, σ diameter, and ϵ energy scale of the interaction potential: typical time

$$\tau = \sigma \sqrt{\frac{m}{\epsilon}}$$

Simple Liquids

- m mass of the particle, σ diameter, and ϵ energy scale of the interaction potential: typical time

$$\tau = \sigma \sqrt{\frac{m}{\epsilon}}$$

- Argon:., $\sigma = 3\text{\AA}$, $m = 6.63 \cdot 10^{-23} \text{ kg}$ and $\epsilon = 1.64 \cdot 10^{-20} \text{ J}$, which gives

$$\tau = 2.8 \cdot 10^{-14} \text{ s}$$

Simple Liquids

- m mass of the particle, σ diameter, and ϵ energy scale of the interaction potential: typical time

$$\tau = \sigma \sqrt{\frac{m}{\epsilon}}$$

- Argon: $\sigma = 3\text{\AA}$, $m = 6.63 \cdot 10^{-23} \text{ kg}$ and $\epsilon = 1.64 \cdot 10^{-20} \text{ J}$, which gives

$$\tau = 2.8 \cdot 10^{-14} \text{ s}$$

- Typically $\Delta t = 10^{-15} \text{ s}$, even smaller.

Simple Liquids

- m mass of the particle, σ diameter, and ϵ energy scale of the interaction potential: typical time

$$\tau = \sigma \sqrt{\frac{m}{\epsilon}}$$

- Argon:., $\sigma = 3\text{\AA}$, $m = 6.63 \cdot 10^{-23} \text{kg}$ and $\epsilon = 1.64 \cdot 10^{-20} \text{J}$, which gives

$$\tau = 2.8 \cdot 10^{-14} \text{s}$$

- Typically $\Delta t = 10^{-15} \text{s}$, even smaller.
- The total number of steps performed in a run is typically of order of magnitude is 10^7 ; the simulation duration is

$$10^{-8} \text{s}$$

From exact equations to discretization



$$m \frac{d^2 \mathbf{r}_i}{dt^2} = - \sum_{j \neq i} \nabla_{\mathbf{r}_i} u(\mathbf{r}_{ij}).$$

where m is the particle mass.

From exact equations to discretization



$$m \frac{d^2 \mathbf{r}_i}{dt^2} = - \sum_{j \neq i} \nabla_{\mathbf{r}_i} u(\mathbf{r}_{ij}).$$

where m is the particle mass.

- Force calculation acting on particle i involves the computation of $(N - 1)$ elementary forces between each particle and particle i

From exact equations to discretization

-

$$m \frac{d^2 \mathbf{r}_i}{dt^2} = - \sum_{j \neq i} \nabla_{\mathbf{r}_i} u(\mathbf{r}_{ij}).$$

where m is the particle mass.

- Force calculation acting on particle i involves the computation of $O(N - 1)$ elementary forces between each particle and particle i
- For the sake of simplicity

$$m \frac{d^2 \mathbf{r}}{dt^2} = \mathbf{f}(\mathbf{r}(t))$$

From exact equations to discretization

- A time series expansion gives

$$r(t + \Delta t) = r(t) + v(t)\Delta t + \frac{f(r(t))}{2m}(\Delta t)^2 + \frac{d^3r}{dt^3}(\Delta t)^3 + \mathcal{O}((\Delta t)^4)$$

From exact equations to discretization

- A time series expansion gives

$$r(t + \Delta t) = r(t) + v(t)\Delta t + \frac{f(r(t))}{2m}(\Delta t)^2 + \frac{d^3r}{dt^3}(\Delta t)^3 + \mathcal{O}((\Delta t)^4)$$

-

$$r(t - \Delta t) = r(t) - v(t)\Delta t + \frac{f(r(t))}{2m}(\Delta t)^2 - \frac{d^3r}{dt^3}(\Delta t)^3 + \mathcal{O}((\Delta t)^4).$$

From exact equations to discretization

- A time series expansion gives

$$r(t + \Delta t) = r(t) + v(t)\Delta t + \frac{f(r(t))}{2m}(\Delta t)^2 + \frac{d^3r}{dt^3}(\Delta t)^3 + \mathcal{O}((\Delta t)^4)$$

-

$$r(t - \Delta t) = r(t) - v(t)\Delta t + \frac{f(r(t))}{2m}(\Delta t)^2 - \frac{d^3r}{dt^3}(\Delta t)^3 + \mathcal{O}((\Delta t)^4).$$

- Adding these above equations, one obtains

$$r(t + \Delta t) + r(t - \Delta t) = 2r(t) + \frac{f(r(t))}{m}(\Delta t)^2 + \mathcal{O}((\Delta t)^4).$$

From exact equations to discretization

- The trajectories are calculated with an accuracy of $(\Delta t)^4$.

From exact equations to discretization

- The trajectories are calculated with an accuracy of $(\Delta t)^4$.

- $$v(t) = \frac{r(t + \Delta t) - r(t - \Delta t)}{2\Delta t} + \mathcal{O}((\Delta t)^2)$$

From exact equations to discretization

- The trajectories are calculated with an accuracy of $(\Delta t)^4$.

- $$v(t) = \frac{r(t + \Delta t) - r(t - \Delta t)}{2\Delta t} + \mathcal{O}((\Delta t)^2)$$

- The Verlet's algorithm preserves the time symmetry. This is, if we change $\Delta t \rightarrow -\Delta t$, discrete equation of motion are the same

From exact equations to discretization

- The trajectories are calculated with an accuracy of $(\Delta t)^4$.

-

$$v(t) = \frac{r(t + \Delta t) - r(t - \Delta t)}{2\Delta t} + \mathcal{O}((\Delta t)^2)$$

- The Verlet's algorithm preserves the time symmetry. This is, if we change $\Delta t \rightarrow -\Delta t$, discrete equation of motion are the same
- Leapfrog algorithm

$$v(t + \Delta t/2) = v(t - \Delta t/2) + \frac{f(r(t))}{m} \Delta t$$

$$r(t + \Delta t) = r(t) + v(t + \Delta t/2) \Delta t$$

Event-driven algorithm

- Hamiltonian

$$\mathcal{H}_N = \sum_i^N \frac{1}{2} m v_i^2 + \frac{1}{2} \sum_{i \neq j} u(r_i - r_j),$$

where the interaction potential is

$$u(r_i - r_j) = \begin{cases} +\infty & |r_i - r_j| \leq \sigma \\ 0 & |r_i - r_j| > \sigma \end{cases}$$

Event-driven algorithm

- Hamiltonian

$$\mathcal{H}_N = \sum_i^N \frac{1}{2} m v_i^2 + \frac{1}{2} \sum_{i \neq j} u(\mathbf{r}_i - \mathbf{r}_j),$$

where the interaction potential is

$$u(\mathbf{r}_i - \mathbf{r}_j) = \begin{cases} +\infty & |\mathbf{r}_i - \mathbf{r}_j| \leq \sigma \\ 0 & |\mathbf{r}_i - \mathbf{r}_j| > \sigma \end{cases}$$

- $\exp(-\beta u(\mathbf{r}_i - \mathbf{r}_j))$ is a binary variable: 0 (when two spheres overlap) or 1 otherwise.

Event-driven algorithm

- Hamiltonian

$$\mathcal{H}_N = \sum_i^N \frac{1}{2} m v_i^2 + \frac{1}{2} \sum_{i \neq j} u(\mathbf{r}_i - \mathbf{r}_j),$$

where the interaction potential is

$$u(\mathbf{r}_i - \mathbf{r}_j) = \begin{cases} +\infty & |\mathbf{r}_i - \mathbf{r}_j| \leq \sigma \\ 0 & |\mathbf{r}_i - \mathbf{r}_j| > \sigma \end{cases}$$

- $\exp(-\beta u(\mathbf{r}_i - \mathbf{r}_j))$ is a binary variable: 0 (when two spheres overlap) or 1 otherwise.
- The excess thermodynamic quantities do not depend on the temperature.

Event-driven algorithm

- Galilean motion between binary collisions

Event-driven algorithm

- Galilean motion between binary collisions
- Exact molecular dynamics: event-driven algorithm

Event-driven algorithm

- Galilean motion between binary collisions
- Exact molecular dynamics: event-driven algorithm
- Consider two spheres of identical mass: during the collision, the total momentum is conserved

$$v_1 + v_2 = v'_1 + v'_2$$

Event-driven algorithm

- Galilean motion between binary collisions
- Exact molecular dynamics: event-driven algorithm
- Consider two spheres of identical mass: during the collision, the total momentum is conserved

$$\mathbf{v}_1 + \mathbf{v}_2 = \mathbf{v}'_1 + \mathbf{v}'_2$$

- Elastic collision

$$(\mathbf{v}'_1 - \mathbf{v}'_2) \cdot \mathbf{n} = -(\mathbf{v}_1 - \mathbf{v}_2) \cdot \mathbf{n}$$

$$(\mathbf{v}'_1 - \mathbf{v}'_2) \cdot \mathbf{t} = (\mathbf{v}_1 - \mathbf{v}_2) \cdot \mathbf{t}$$

Event-driven algorithm



$$v'_1 = v_1 + \frac{(v_2 - v_1) \cdot (r_1 - r_2)}{(r_1 - r_2)^2} (r_1 - r_2)$$

$$v'_2 = v_2 + \frac{(v_1 - v_2) \cdot (r_1 - r_2)}{(r_1 - r_2)^2} (r_1 - r_2)$$

Event-driven algorithm

- $$v'_1 = v_1 + \frac{(v_2 - v_1) \cdot (r_1 - r_2)}{(r_1 - r_2)^2} (r_1 - r_2)$$
$$v'_2 = v_2 + \frac{(v_1 - v_2) \cdot (r_1 - r_2)}{(r_1 - r_2)^2} (r_1 - r_2)$$

- Between collisions, the positions of i and j are given by

$$r_i = r_i^0 + v_i t$$

$$r_j = r_j^0 + v_j t$$

Event-driven algorithm

- $$v'_1 = v_1 + \frac{(v_2 - v_1) \cdot (r_1 - r_2)}{(r_1 - r_2)^2} (r_1 - r_2)$$
$$v'_2 = v_2 + \frac{(v_1 - v_2) \cdot (r_1 - r_2)}{(r_1 - r_2)^2} (r_1 - r_2)$$

- Between collisions, the positions of i and j are given by

$$r_i = r_i^0 + v_i t$$

$$r_j = r_j^0 + v_j t$$

- The contact condition between two particles is given by relation

$$\sigma^2 = (r_i - r_j)^2 = (r_i^0 - r_j^0)^2 + (v_i - v_j)^2 t^2 + 2(r_i^0 - r_j^0)(v_i - v_j)t$$

Event-driven algorithm

- The collision time is a solution of a quadratic equation,

Event-driven algorithm

- The collision time is a solution of a quadratic equation,
- ① Complex roots: the collision time is set to a very large value in simulation

Event-driven algorithm

- The collision time is a solution of a quadratic equation,
- ① Complex roots: the collision time is set to a very large value in simulation
- ② Real negative roots: : the collision time is set to a very large value in simulation

Event-driven algorithm

- The collision time is a solution of a quadratic equation,
- - 1 Complex roots: the collision time is set to a very large value in simulation
 - 2 Real negative roots: : the collision time is set to a very large value in simulation
 - 3 Positive and negative roots: the collision time corresponds to the positive value

Event-driven algorithm

- The collision time is a solution of a quadratic equation,
 - 1 Complex roots: the collision time is set to a very large value in simulation
 - 2 Real negative roots: : the collision time is set to a very large value in simulation
 - 3 Positive and negative roots: the collision time corresponds to the positive value
 - 4 Two positive roots: the smallest one is selected as the collision time.

Event-driven algorithm

- The collision time is a solution of a quadratic equation,
 - 1 Complex roots: the collision time is set to a very large value in simulation
 - 2 Real negative roots: : the collision time is set to a very large value in simulation
 - 3 Positive and negative roots: the collision time corresponds to the positive value
 - 4 Two positive roots: the smallest one is selected as the collision time.
- A necessary condition for having a positive solution for t (collision in the future) is that

$$(\mathbf{r}_i^0 - \mathbf{r}_j^0) \cdot (\mathbf{v}_i - \mathbf{v}_j) < 0$$

Event-driven algorithm

- Once all pair of particles have been examined, the shortest time is selected, namely the first collision.

Event-driven algorithm

- Once all pair of particles have been examined, the shortest time is selected, namely the first collision.
- The trajectories of particles evolve rectilinearly until the new collision occurs (calculation of trajectories is exact)
- At the collision time, velocities of the two particles are updated, and one calculates the next collision again.