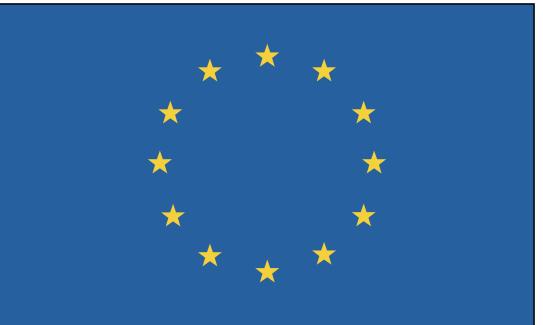




PSL



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Quantum coherent effects in the hydrodynamics of low-temperature 1D gases

Stefano Scopa

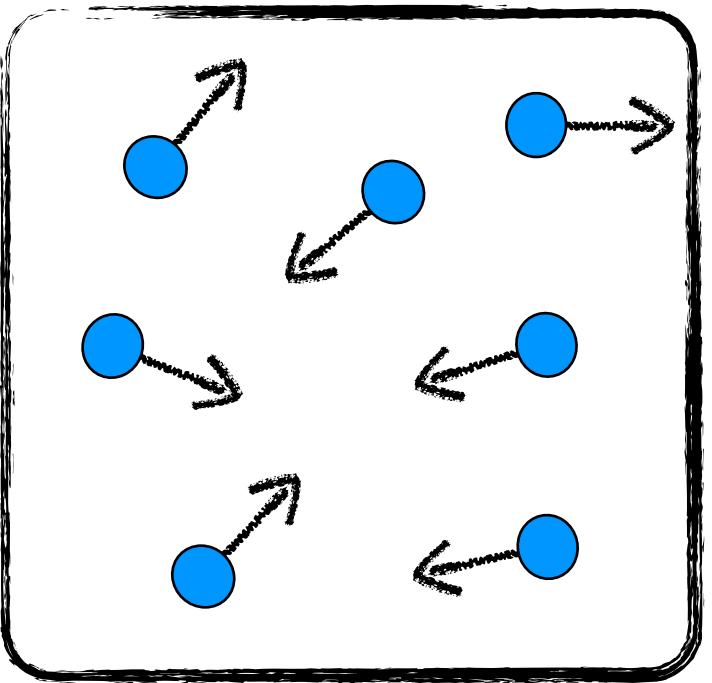
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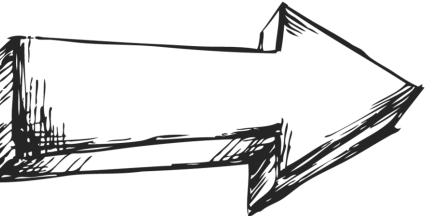
1 April 2025 – LPTMC, Paris

Invitation to nonequilibrium physics



$|\psi\rangle$

Statistical Mechanics



Equilibrium Physics:

Statistical Ensembles



$$\hat{\rho} \propto e^{-\sum_n \beta_n \hat{Q}_n}$$

↑
Conserved quantities

*describe many-body **quantum systems**



exponential complexity $\sim e^N$

1 spin: $|\uparrow\rangle$ or $|\downarrow\rangle$ states

N spins on a lattice $\sim 2^N$ states

$$2^{20} = 1\ 048\ 576$$

Invitation to nonequilibrium physics

Nonequilibrium Physics: There is no unique framework for nonequilibrium processes



Quantum quench:

- Prepare the system at equilibrium of $\hat{H}(V_0)$
- Sudden change of $\hat{H}(V_0) \rightarrow \hat{H}(V_1)$

$$|GS\rangle \mapsto |\Psi_t\rangle = e^{-it\hat{H}}|GS\rangle$$



exponential complexity

How ?

general mechanisms



Effective descriptions

- No microscopic laws
- Universality

Emergent Hydrodynamics

Look for an **effective description** at large space & time scales

I. Scales separation hypothesis

Assume local stat phys ensembles



$$\hat{\rho} \propto e^{-\sum_n \beta_n(x,t) \hat{Q}_n}$$

$$\langle \hat{o}(x, t) \rangle = \langle \hat{o} \rangle_{\{\beta_n(x, t)\}}$$

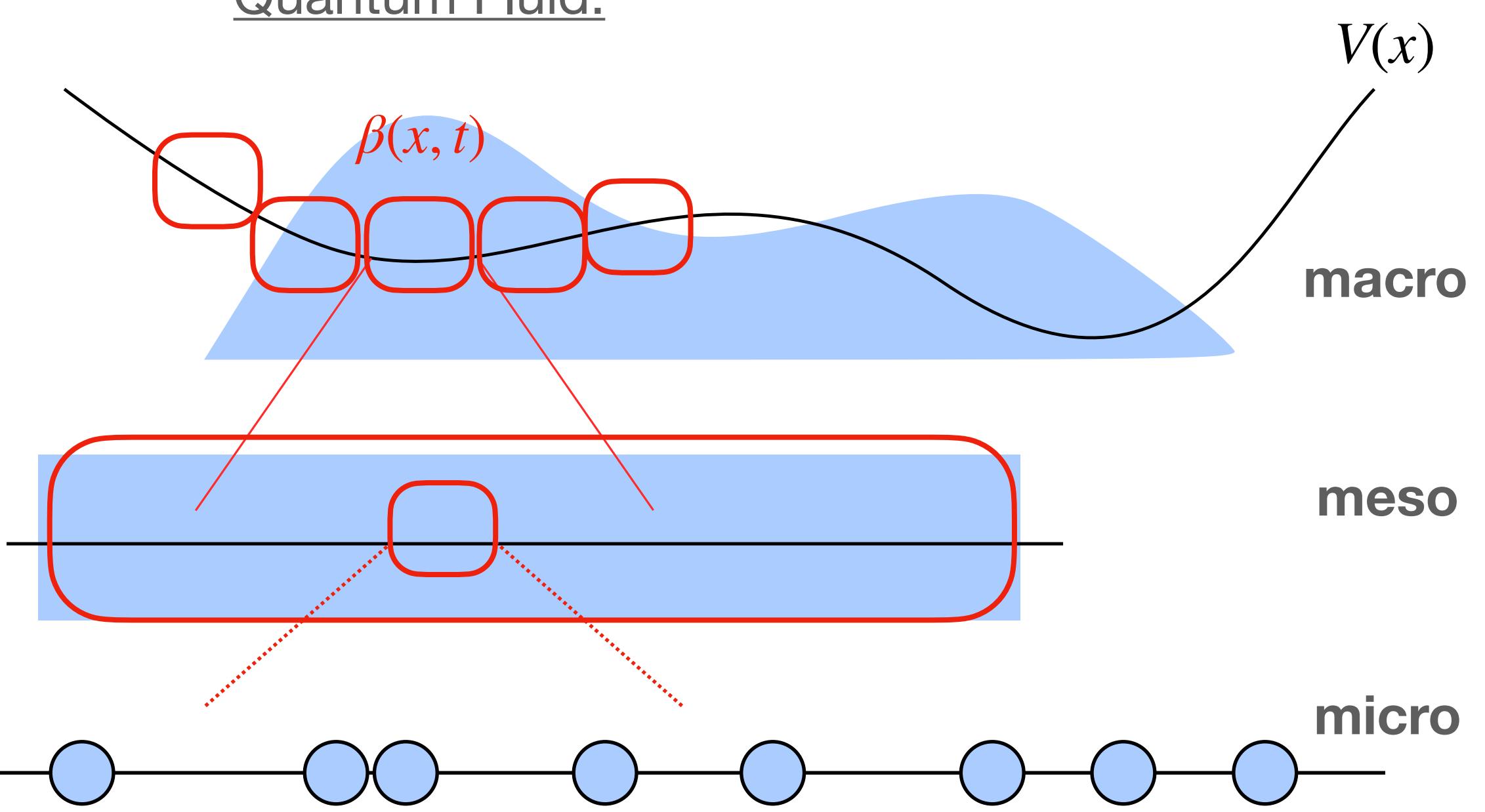
II. Local charges

$$\langle \hat{Q}_n \rangle = \int dx q_n(x, t)$$

III. Current (*need eq of state*)

$$j_n = \mathbb{F}[\{q\}, t]$$

Quantum Fluid:



IV. Continuity equation

$$\partial_t q_n + \partial_x j_n = 0$$

Integrable models

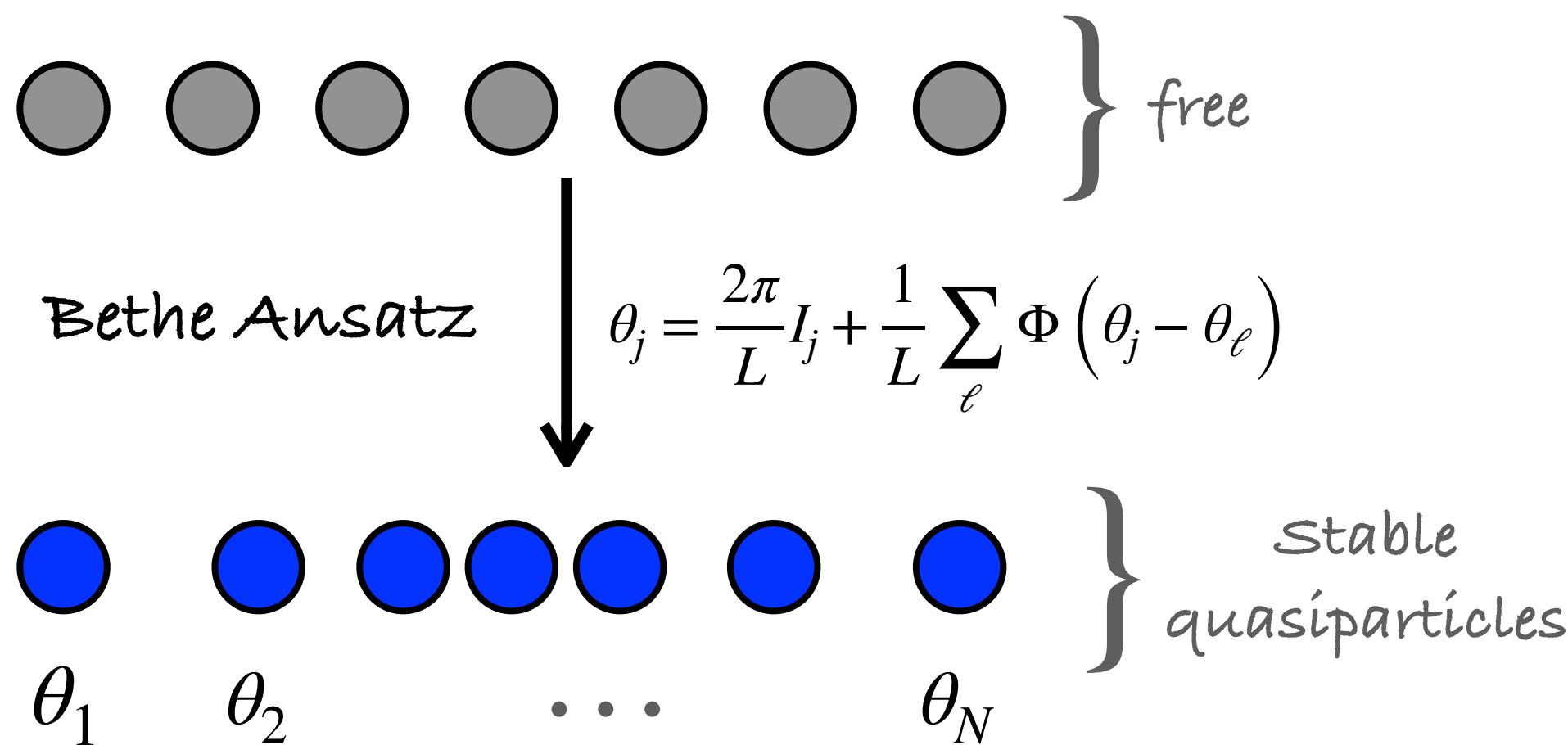


Importance of exact solutions \leftrightarrow infer general mechanisms on nonequilibrium physics

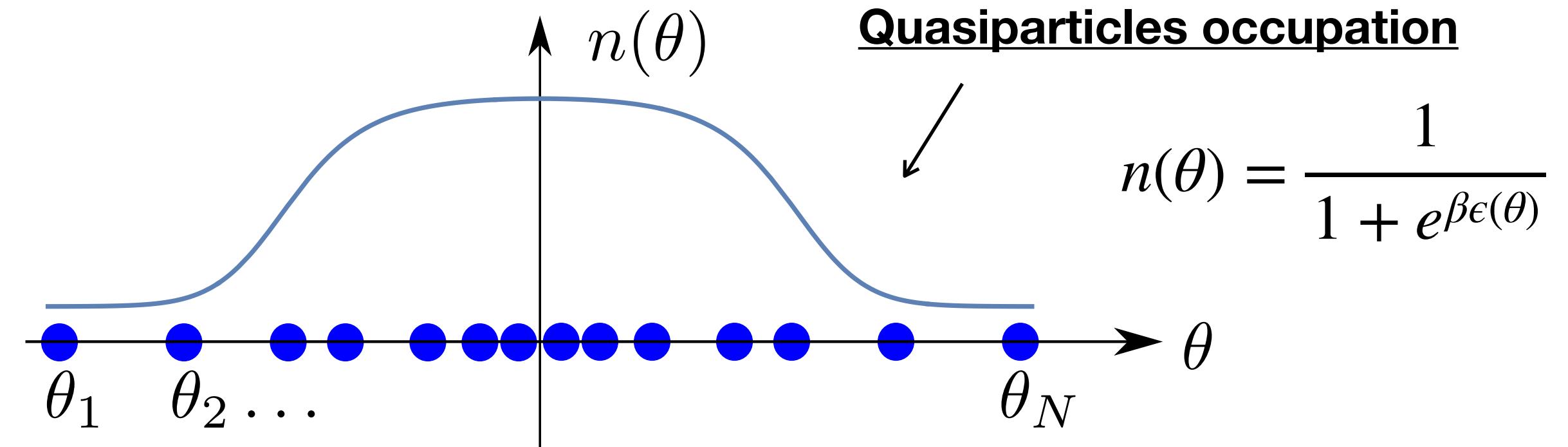
❖ Diffusive transport

❖ Ballistic transport \rightarrow **Integrable models**

- Models featuring ∞ cons. laws $[\hat{H}, \hat{Q}_n] = 0$
- Many-body wavefunction: $|\Psi\rangle = |\theta_1, \dots, \theta_N\rangle$



Thermodynamic limit:



Realizations in Cold-Atom Experiments

Ultracold atom gases

- Quasi-1D geometries
- tunable interaction and trapping potentials
- temperatures up to nK

Reference model:

Lieb-Liniger model

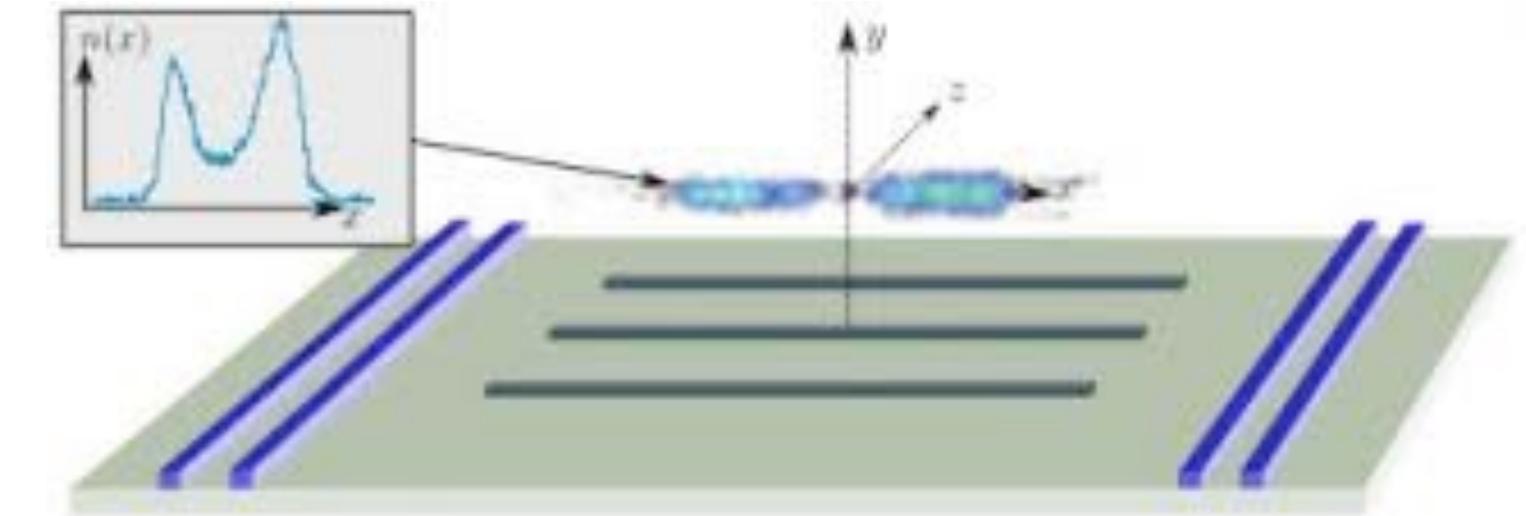
[Lieb, Liniger 1963]

*Describe Bose gases in 1D with contact (repulsive) interactions

$$\hat{H} = -\frac{\hbar^2}{2m} \sum_{i=1}^N \frac{\partial^2}{\partial x_i^2} + g \sum_{1 \leq i < j \leq N} \delta(x_i - x_j) + \sum_{i=1}^N V(x_i)$$

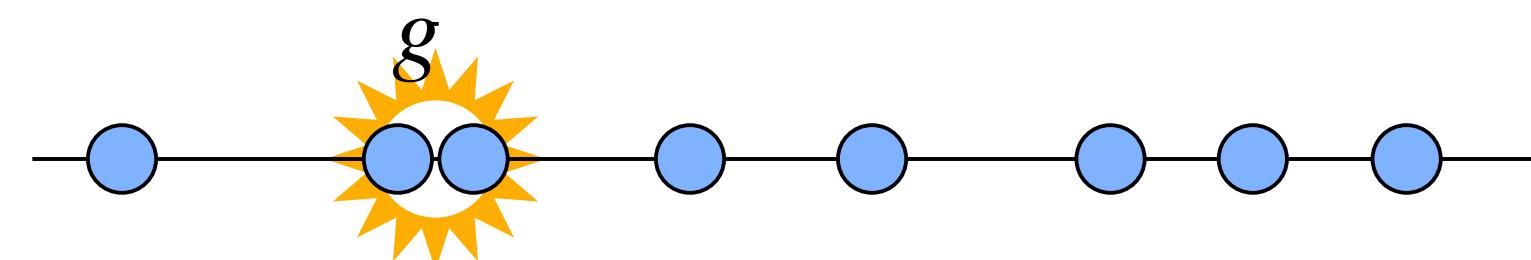
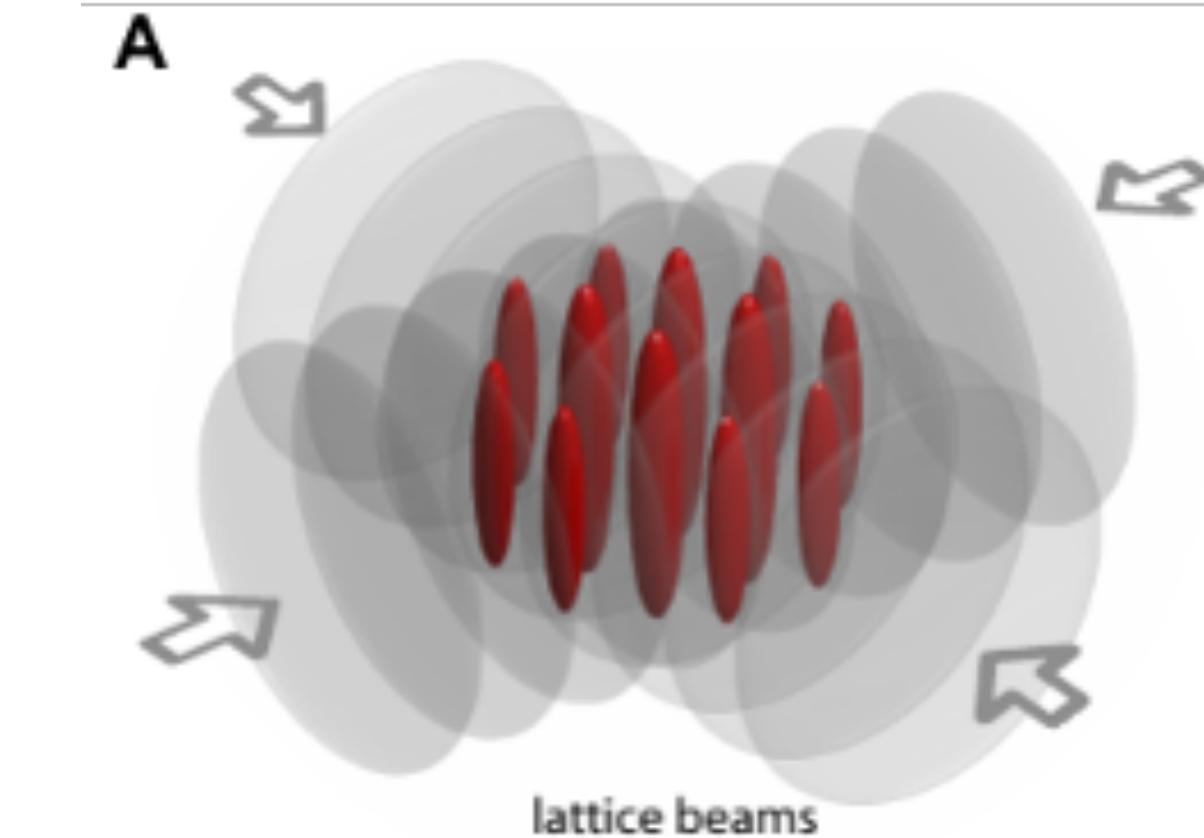
- Atom chips

magnetic fields generated by micro-fabricated wire patterns on a chip to trap and manipulate cold atoms very close to the chip surface



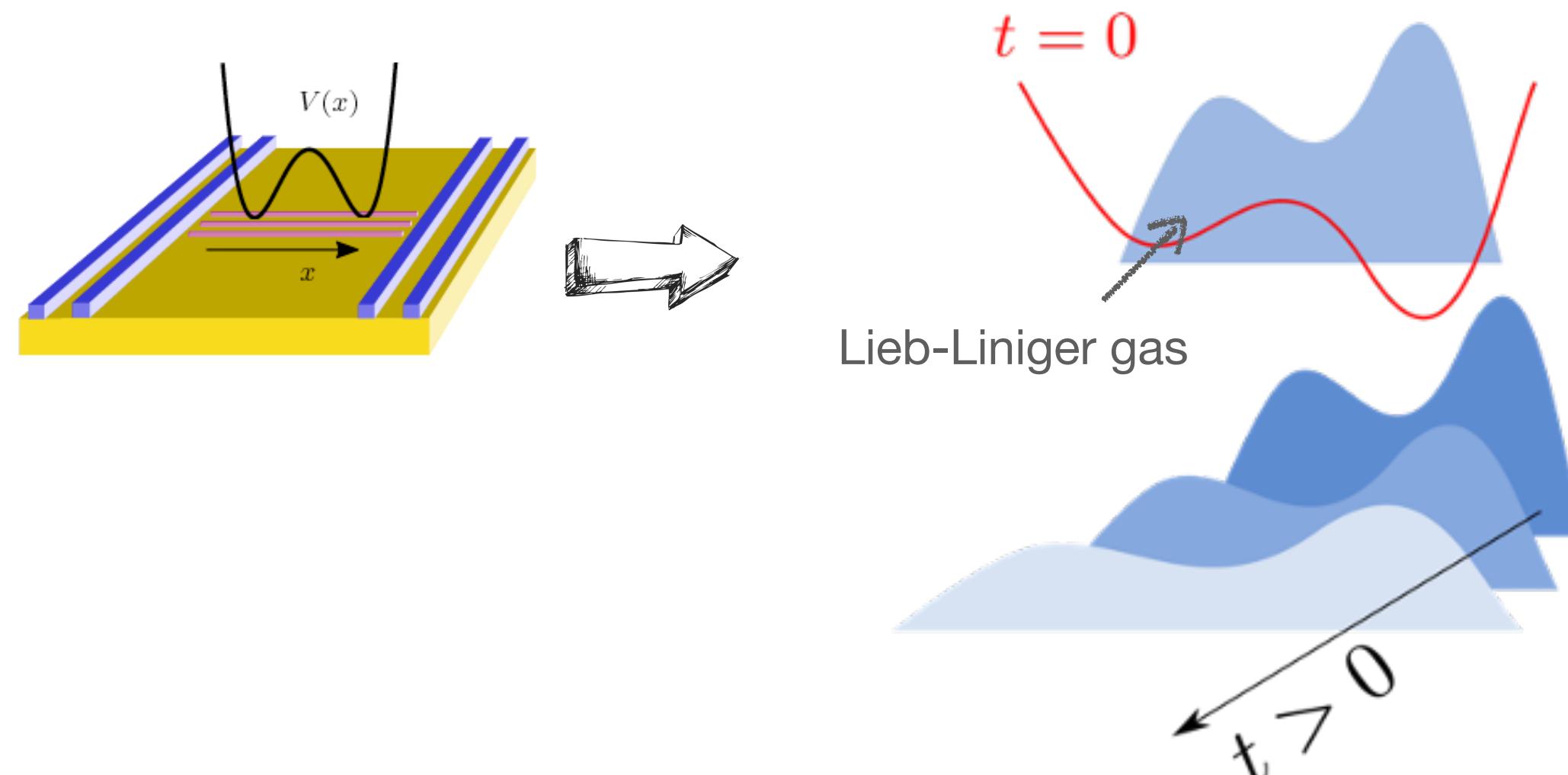
- Optical lattices

interfering laser beams used to create periodic potentials that trap neutral atoms at the (anti)nodes of the standing wave. This generates an array of quasi 1D tubes



Integrable ($V = 0$)

Setup of the problem



Protocol:

- Prepare the system at equilibrium with some potential $V_0(x)$
- Low temperature:** $|\Psi(0)\rangle \approx |\text{GS}\rangle$
- At $t = 0$ sudden change of the confining trap

$$|\text{GS}\rangle \mapsto |\Psi(t)\rangle = e^{-it\hat{H}}|\text{GS}\rangle$$

nonequilibrium
dynamics

Goal:

- ▶ Local quantities (e.g. particle density)
- ▶ **Quantum coherent effects** (e.g. one-body correlations and entanglement)

why?

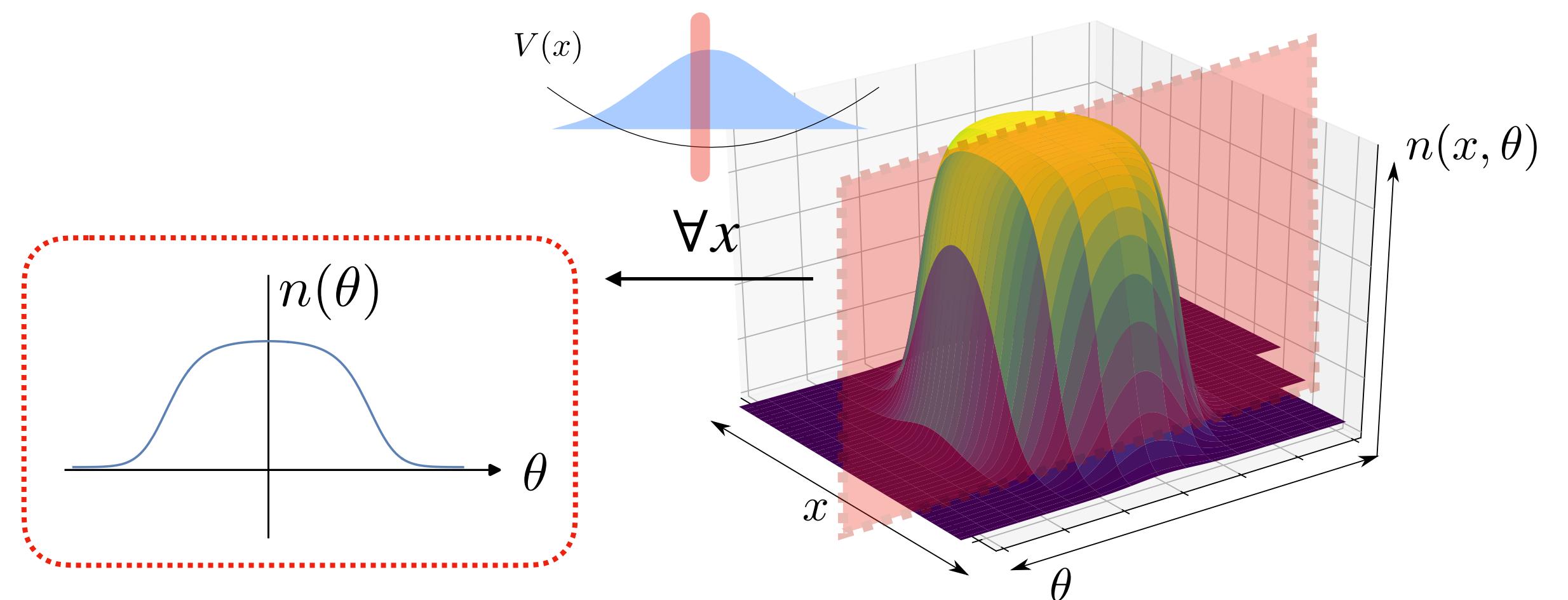
- ▶ **Theory prediction** for some experimental measurment. (beyond current hydrodynamic results)
- ▶ Better understand the nonequilibrium dynamics of low-temperature 1D gases

Hydrodynamics of Integrable models

X Integrability is lost $V \neq 0 \rightarrow \text{Emergent Hydrodynamics}$

! ∞ cons. laws: $\partial_t q_n + \partial_x j_n = 0$ not a good strategy!

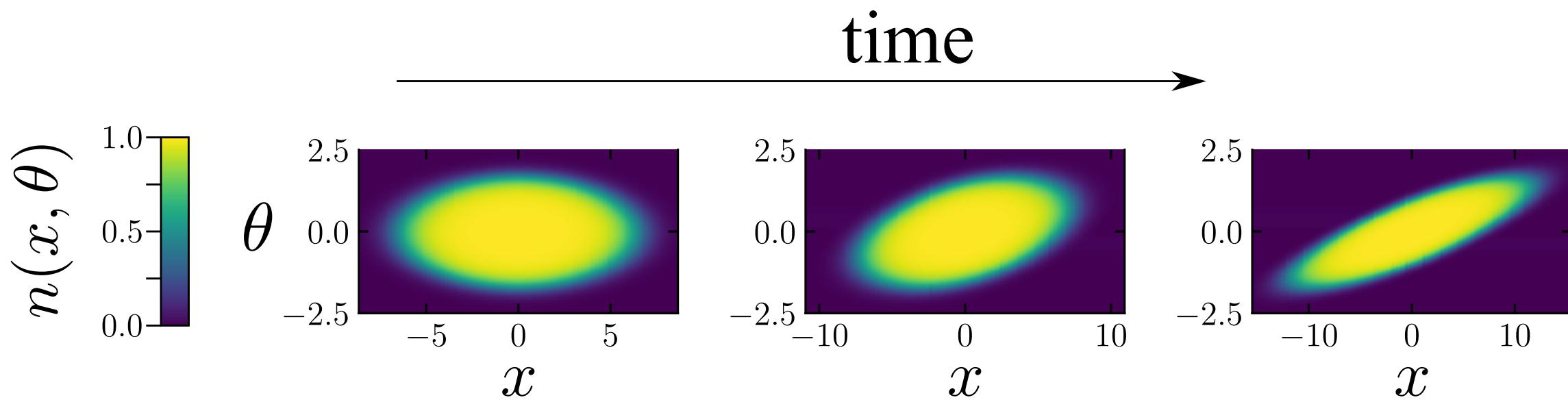
⌚ $t = 0$: Phase-space distribution of quasiparticles



Phase-space dynamics \rightarrow **Euler equation**

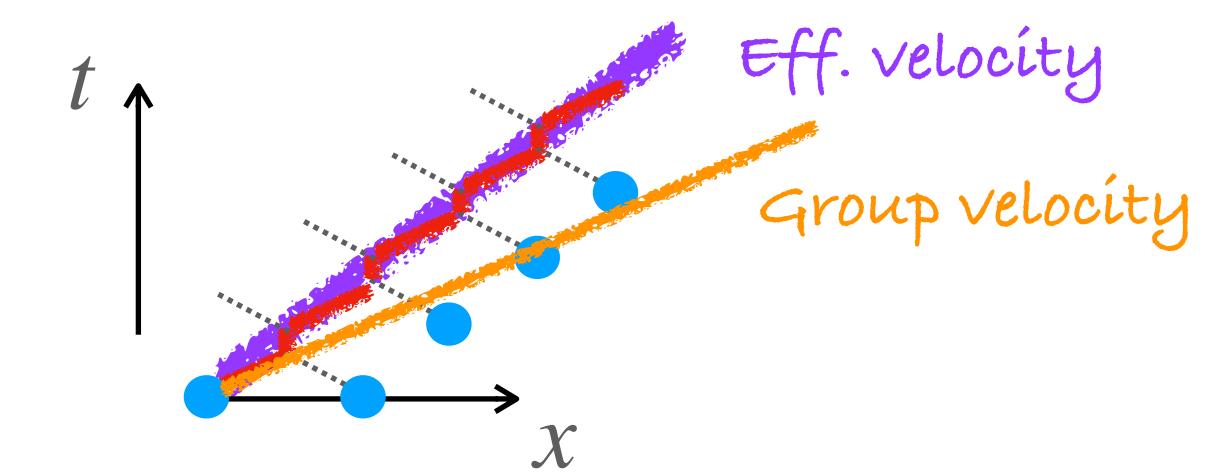
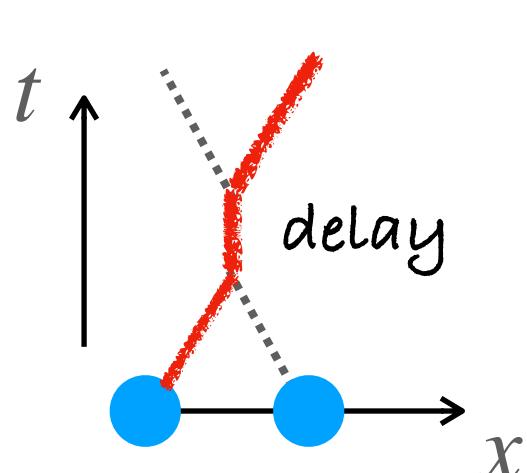
$$(\partial_t + v^{\text{eff}}[n] \partial_x) n_t(x, \theta) = 0$$

⌚ $t > 0$:



effective velocity : $v^{\text{eff}}(\theta) = \theta - \int d\theta' \rho_s(\theta') n(\theta') \varphi(\theta - \theta') [v^{\text{eff}}(\theta) - v^{\text{eff}}(\theta')]$

[Bertini et al, 1605.09790]
[C.-Alvaredo et al, 1605.07331]

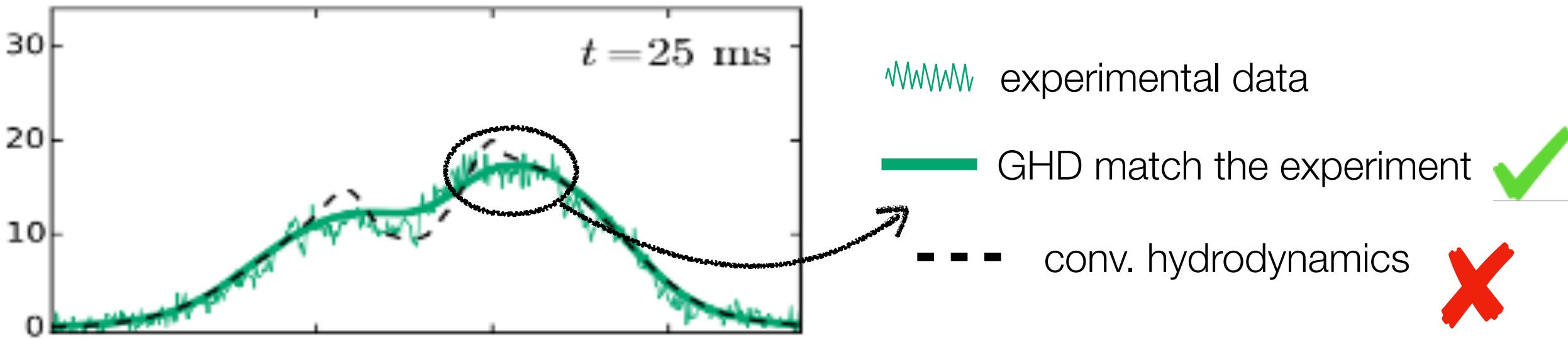


Missing Quantum Fluctuations

Quantitative results for (mean) charges/currents

[Bouchoule group, 1810.07170]

+many more Weiss, Schmiedmayer,
Lev, Nagerl, etc



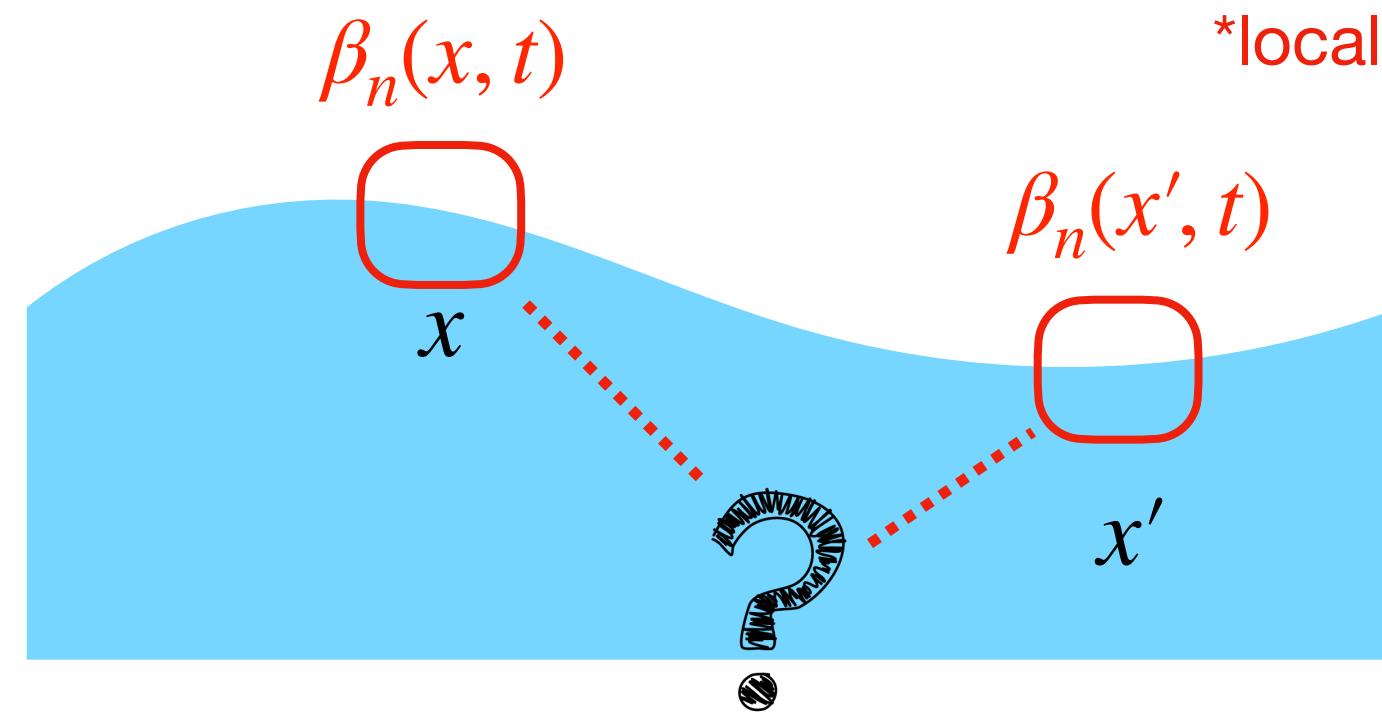
Missing fluctuations in the hydrodynamic theory

fluctuations — thermal, quantum, noise, ...

zero-temperature
limit

❖ What we miss...

- Zero-temperature entanglement
- Equal-time correlation functions
- Quantum corrections to charge/currents profiles



*local GGE description
*absence of long-range correlations
between x and x'

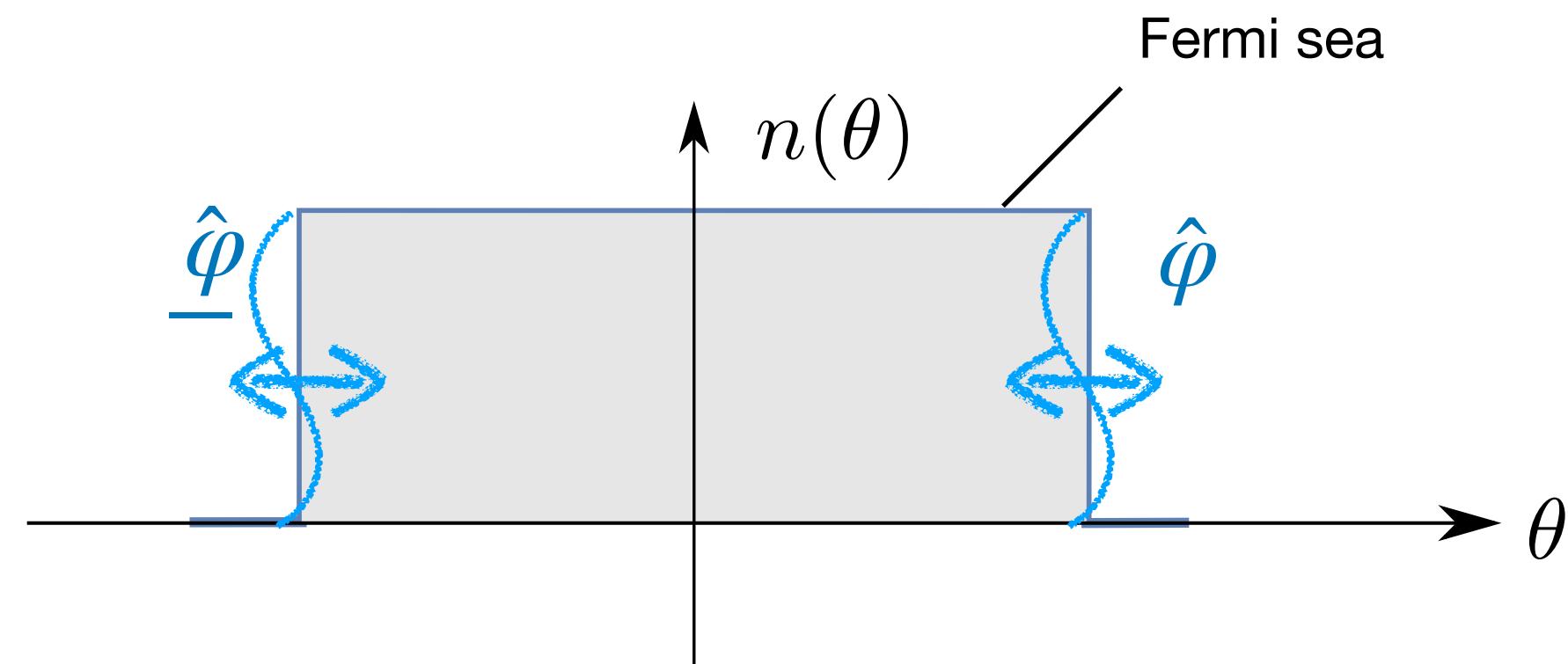
Re-quantization of the theory?

Incorporating Quantum Fluctuations

❖ back to **equilibrium**:

Ground state of integrable models

- Fermionic quasiparticles → **Fermi sea**
- Relevant quantum fluctuations are **particle-hole excitations**
- **effective description** with fluctuating fields at the Fermi edges



► Bosonization

$$\hat{\psi}^\dagger = B_0 : \exp\left(\frac{i\hat{\phi} + i\bar{\hat{\phi}}}{2\sqrt{K}}\right) : + \dots$$

K = compressibility of the gas

$\hat{\phi}, \bar{\hat{\phi}}$ = fluctuating fields with 1D effective Hamiltonian

► Tomonaga-Luttinger liquid

$$\hat{\mathcal{H}} = \int dx \frac{v_s}{4\pi K} \left((\partial_x \hat{\phi})^2 + (\partial_x \bar{\hat{\phi}})^2 \right)$$

GFF

❖ $\langle \hat{\phi}(x)\hat{\phi}(x') \rangle = \log[-i(x-x')] :$ 2pt function fixed by CFT

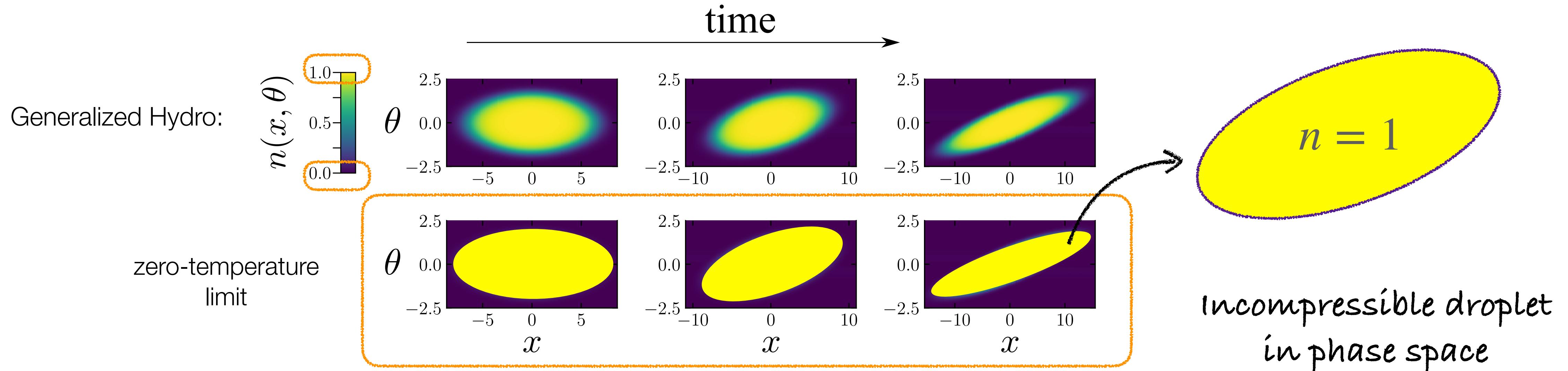
+ Wick's theorem:

$$g_1(x) = \langle \hat{\psi}^\dagger(x)\hat{\psi}(0) \rangle \sim |x|^{-\frac{1}{2K}}$$

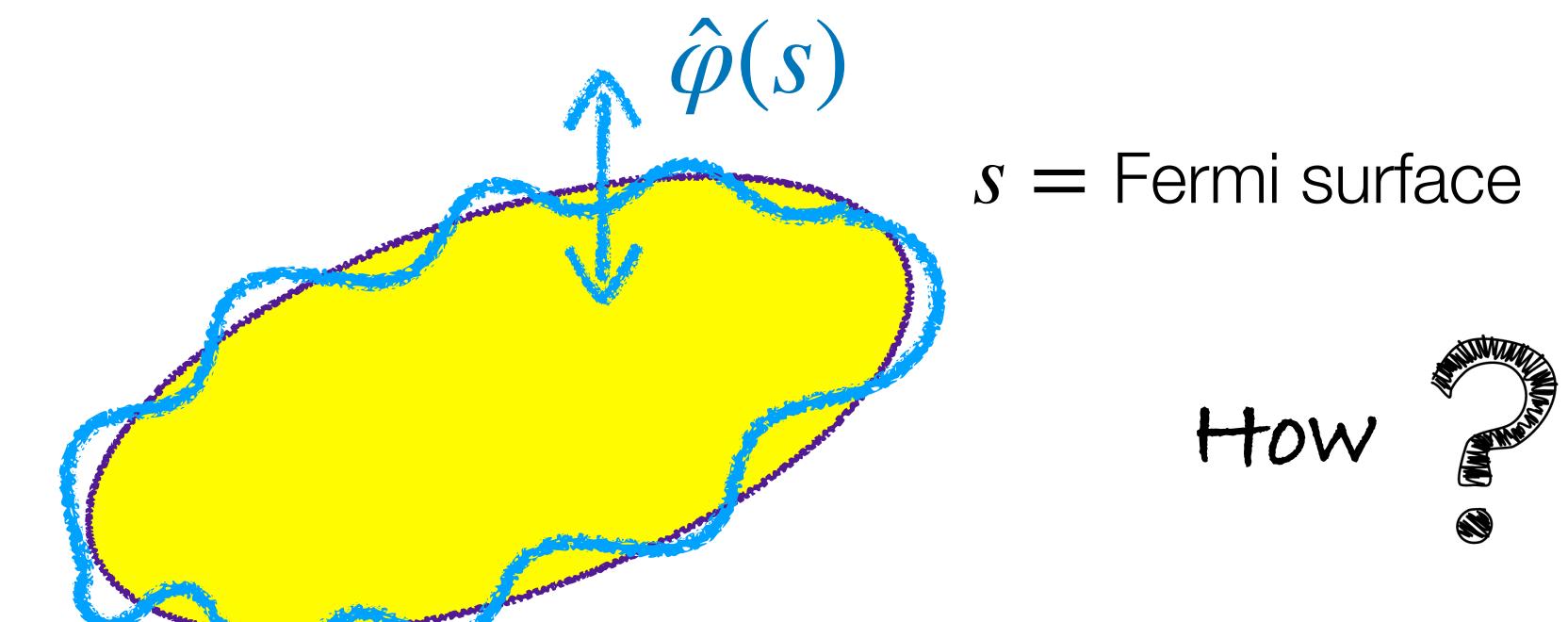
*algebraic decay of correlations

Incorporating Quantum Fluctuations

- ❖ Out of equilibrium?



- ❖ Quantum Fluctuating Hydrodynamics

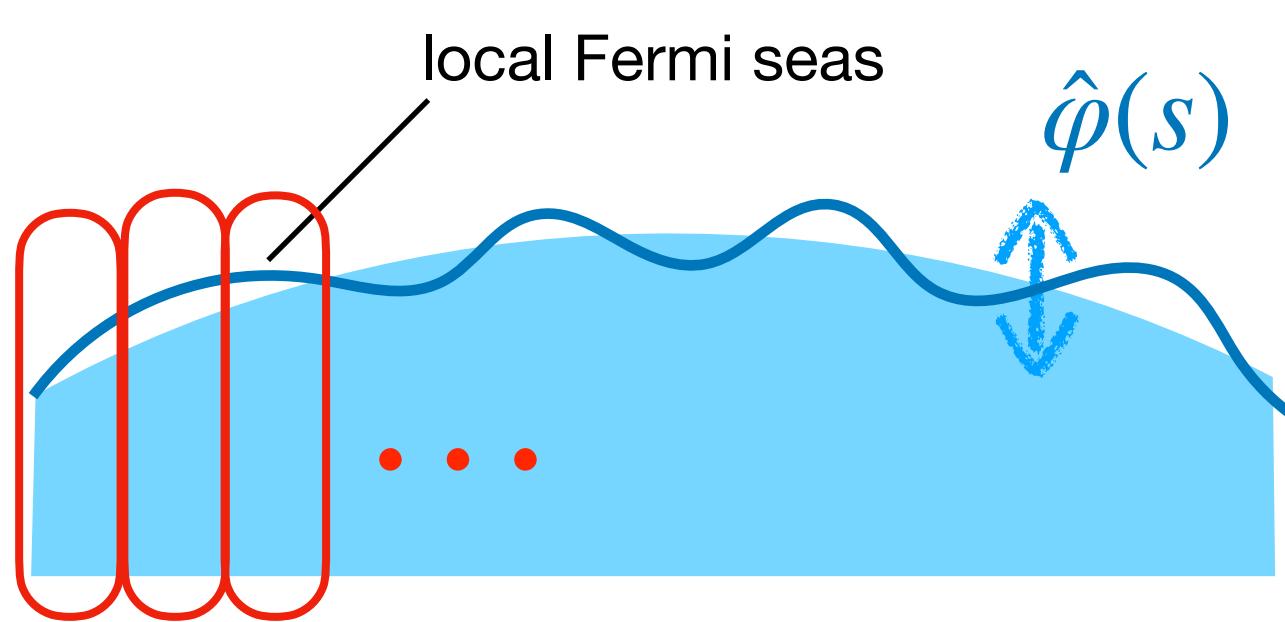


$s = \text{Fermi surface } (x_F(s); \theta_F(s))$

How ?

Incorporating Quantum Fluctuations

 $t = 0$: Ground state of inhomogeneous gas



 $K = K[\rho(x)]$ conformal invariance is broken

2D Electrostatic analogy:

$$[d^2\mathbf{x} = dx^2 + v_s^2(x)d\tau^2 \propto ds^2 + d\tau^2]$$

$$\nabla 4K(\mathbf{x}) \nabla G_0^{[\Theta\Theta]}(\mathbf{x}, \mathbf{y}) = 4\pi\delta^{(2)}(\mathbf{x} - \mathbf{y})$$

Finite interaction,
Finite (small) temperature

$$g_1(x, y) = e^{-G_0^{[\Theta\Theta]}(s_x, s_y)} \prod_{z=x,y} \frac{|B_0(z)|^2 \rho(z)^{\frac{2K(z)-1}{4K(z)}}}{(v_z)^{\frac{1}{4K(z)}}} e^{\frac{1}{2}G_0^{[\Theta\Theta]}(s_z)}$$

Inhomogeneous GFF

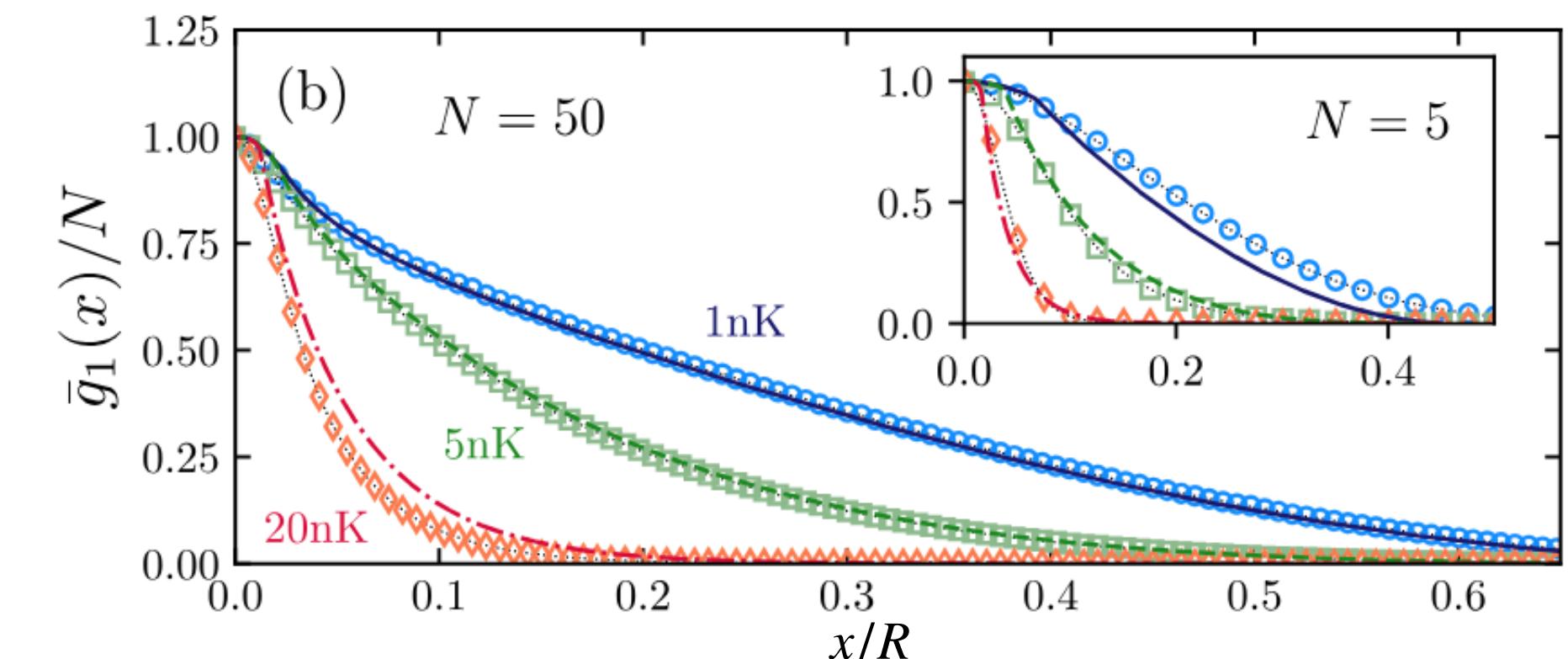
$$\hat{\mathcal{H}} = \int dx \frac{v_s(x)}{2\pi} \left(\frac{(\partial_x \hat{\Phi})^2}{K(x)} + K(x)(\pi \hat{\Pi})^2 \right)$$

numerical input   **Scopa et al, 2005.10214**

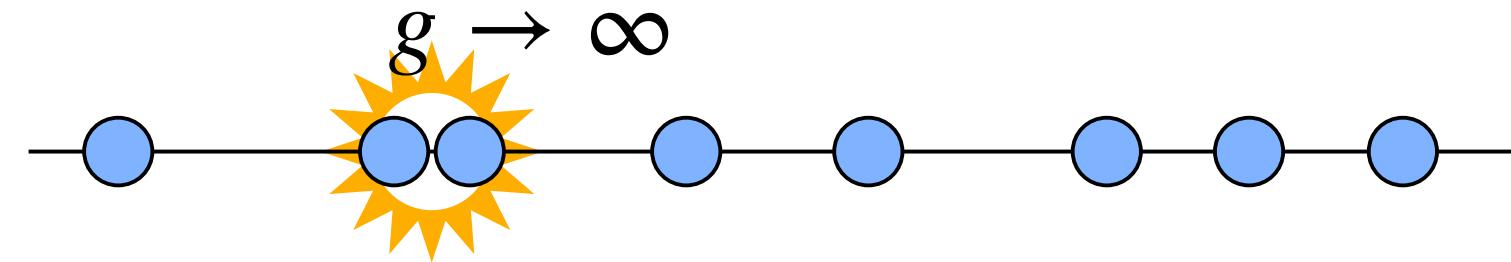
$$C_0 = \begin{pmatrix} G_0^{[\Phi\Phi]} & G_0^{[\Phi\Theta]} \\ G_0^{[\Theta\Phi]} & G_0^{[\Theta\Theta]} \end{pmatrix}$$



Takacs, Zhang, Calabrese, Dubail, Rigol, **Scopa**,
2409.16929

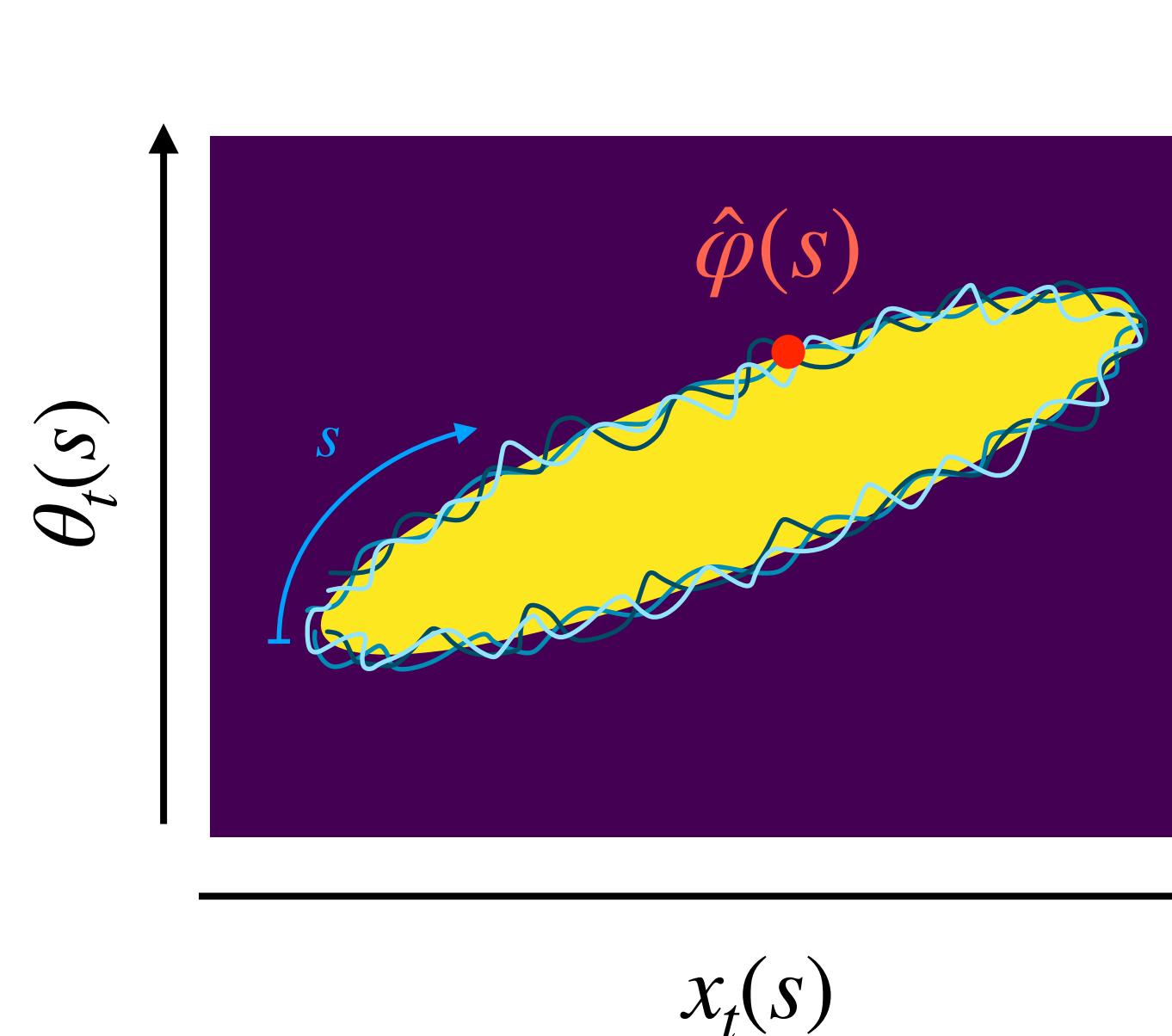


The strong repulsion limit

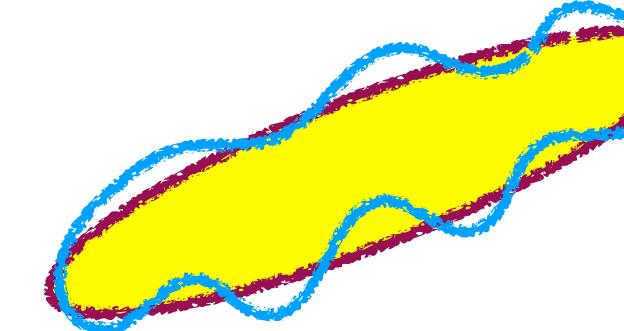
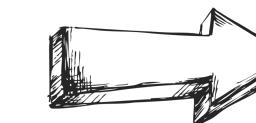
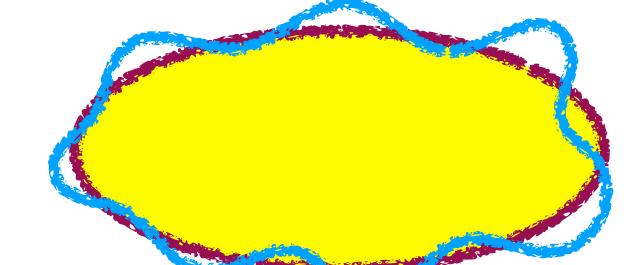


Tonks-Girardeau limit ($K \equiv 1$)

- gas of hardcore bosons
- $\langle \hat{\phi}(s)\hat{\phi}(s') \rangle = \text{boundary CFT in } s \text{ space}$
- Gaussian correlations are comovingly transported with the background:



$$\langle \hat{\phi}(s)\hat{\phi}(s') \rangle_{t=0} = \langle \hat{\phi}(s)\hat{\phi}(s') \rangle_{t>0}$$



analytical calculations

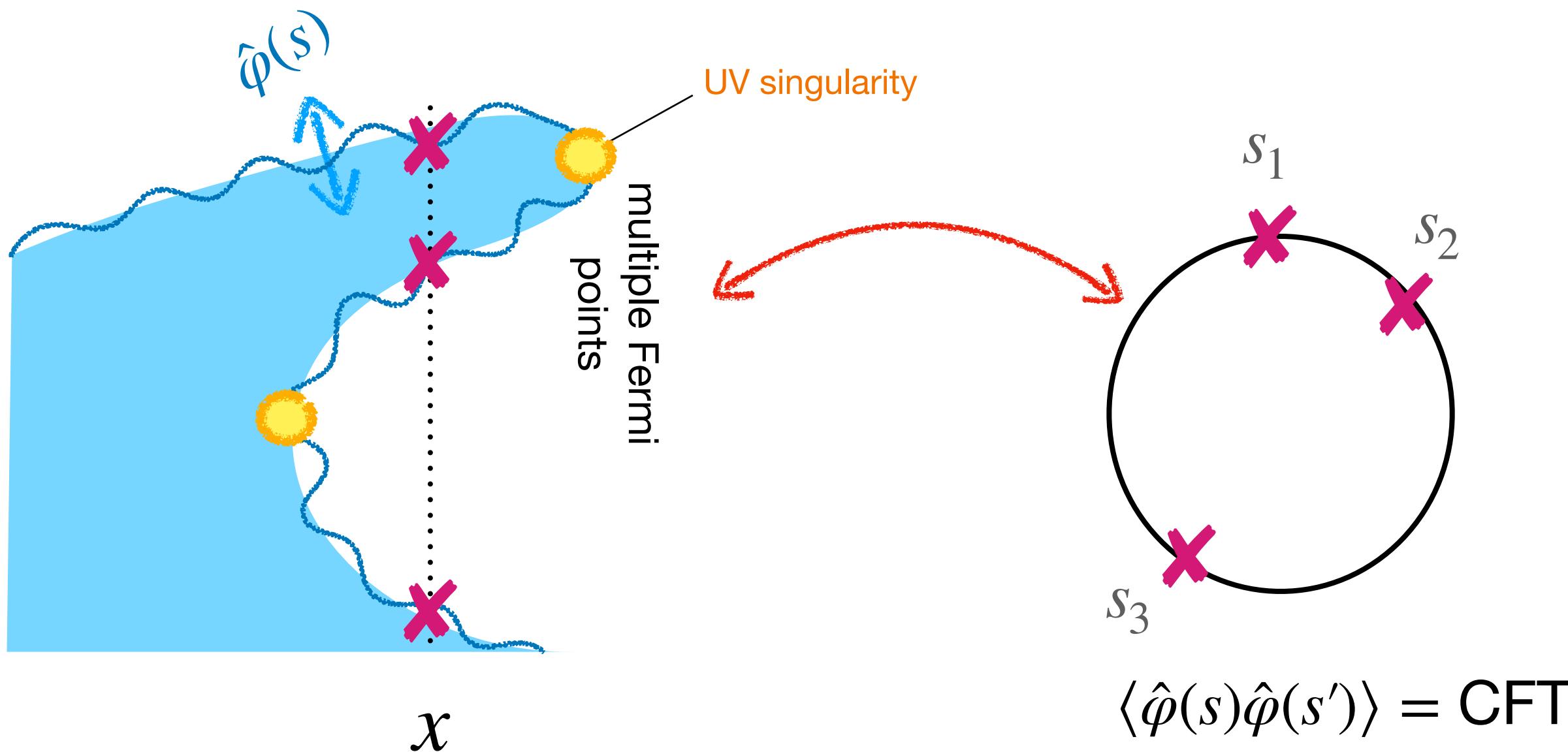


Bosonization Out of Equilibrium



Scopa et al, 2301.04094

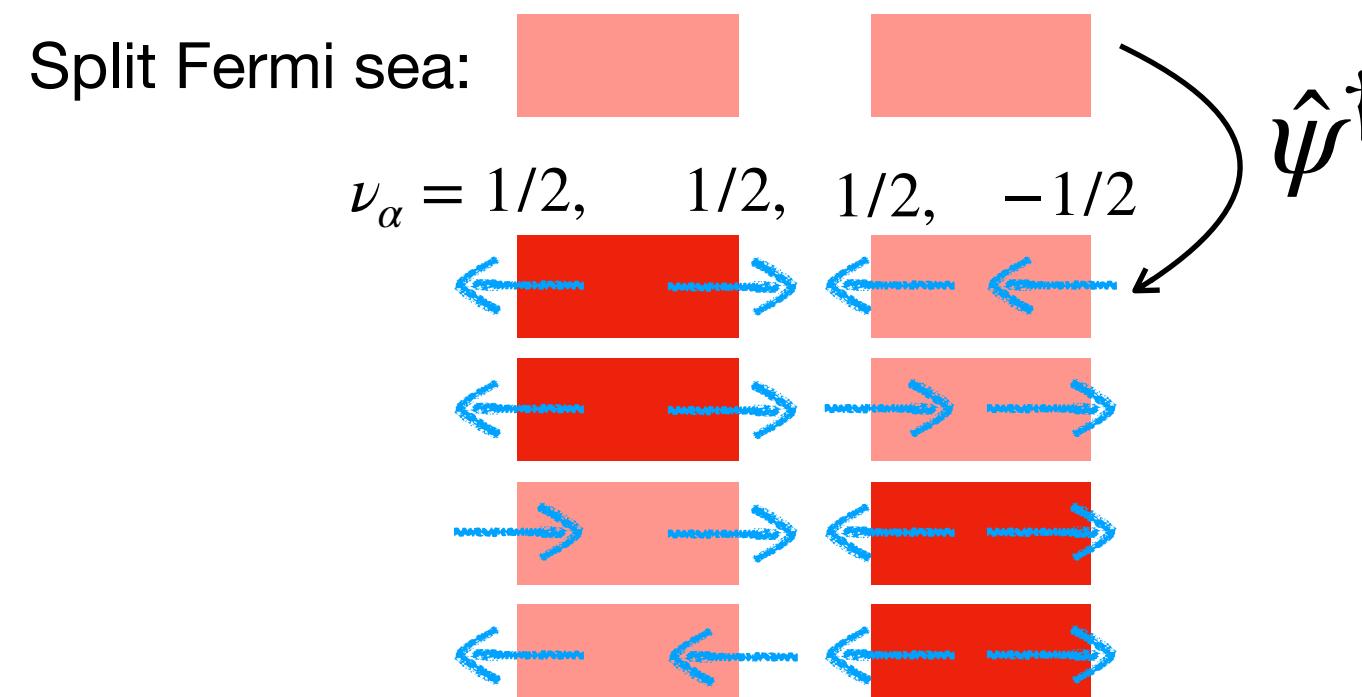
❖ Summarizing:



Non-univ. Amplitude: Semiclassical WKB phase

$$\hat{\psi}^\dagger = \sum_{\vec{\nu}} B_{\vec{\nu}}(x, t) \prod_{\alpha} e^{i\nu_{\alpha}\phi_t(x)} : e^{i\nu_{\alpha}\hat{\phi}(s_{\alpha})} : + \dots$$

Sum over configurations Gaussian fields



$$B_{\vec{\nu}}(x, t) \propto | \langle \leftarrow \rightarrow \leftarrow \leftarrow | \hat{\psi}^\dagger(0) | \text{red band} \text{ red band} \rangle |$$

One-Body Correlations and Momentum distribution

$$g_1(t; x, y) = \langle \Psi_t | \hat{\psi}^\dagger(x) \hat{\psi}(y) | \Psi_t \rangle \longrightarrow n_t(p) = \iint dx dy e^{ip(x-y)} g_1(t; x, y)$$



Scopa et al, 2301.04094

* Any trap quench $V_0(x) \rightarrow V_1(x)$

Fully analytical result:

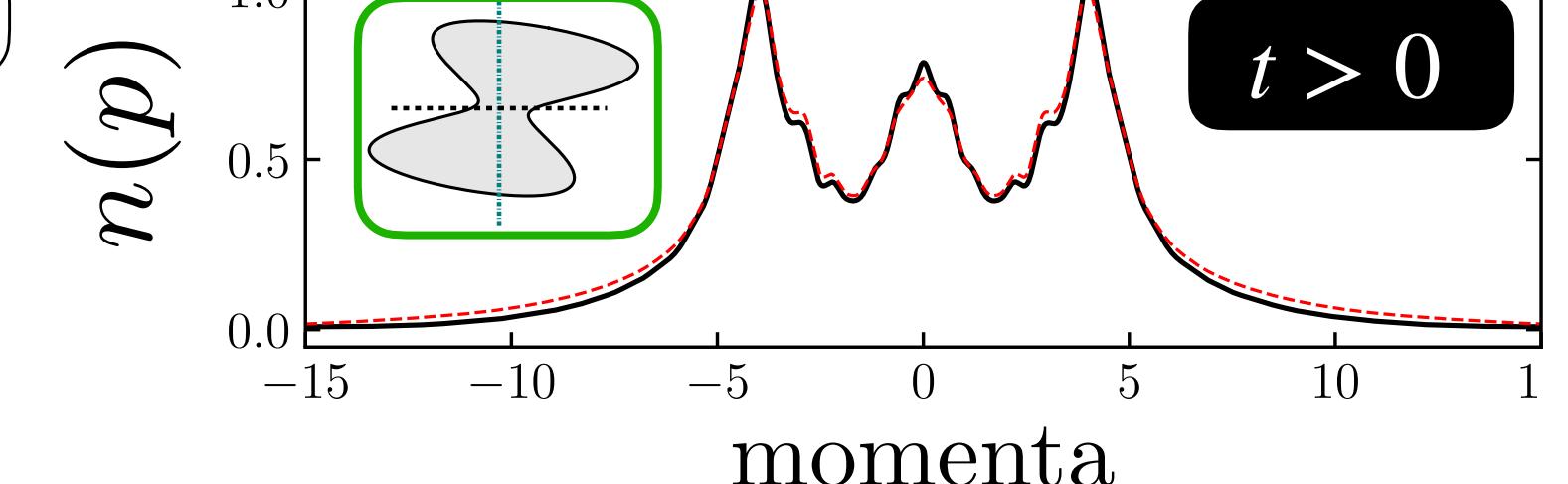
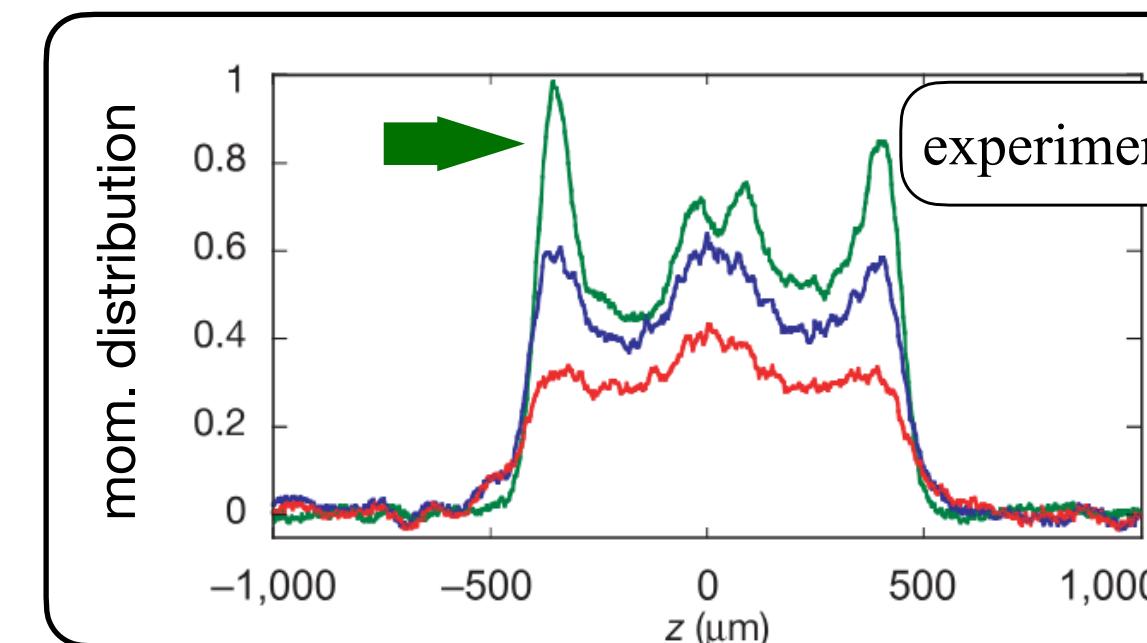
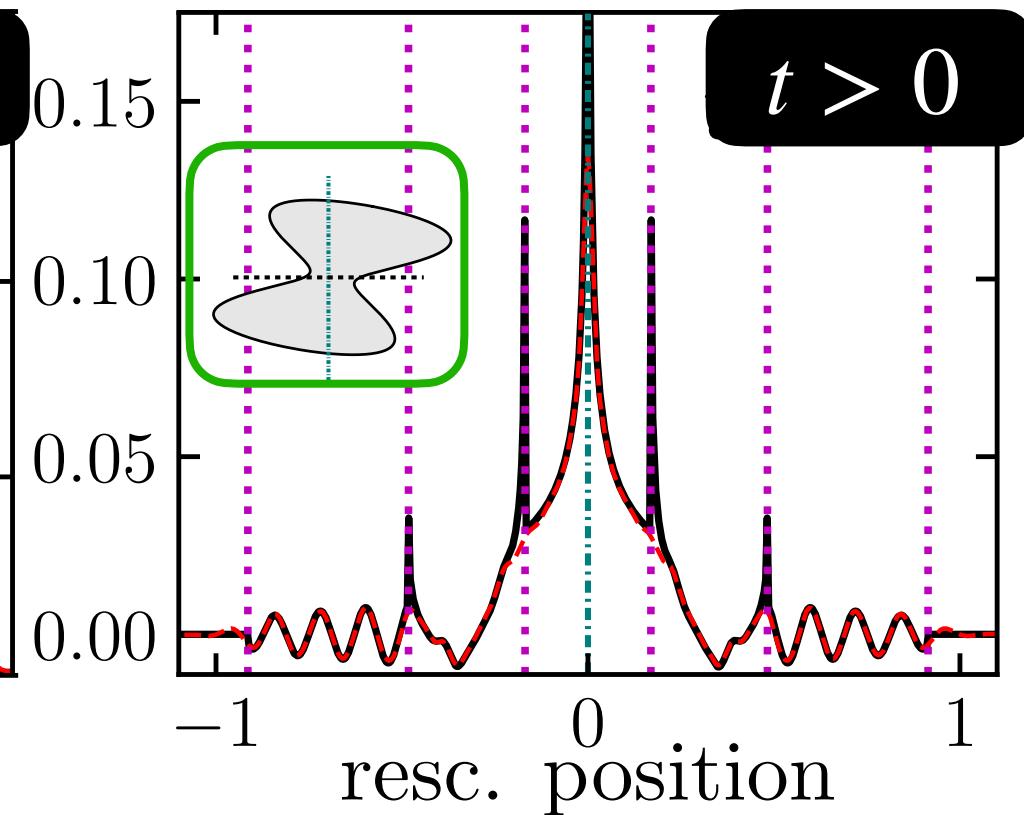
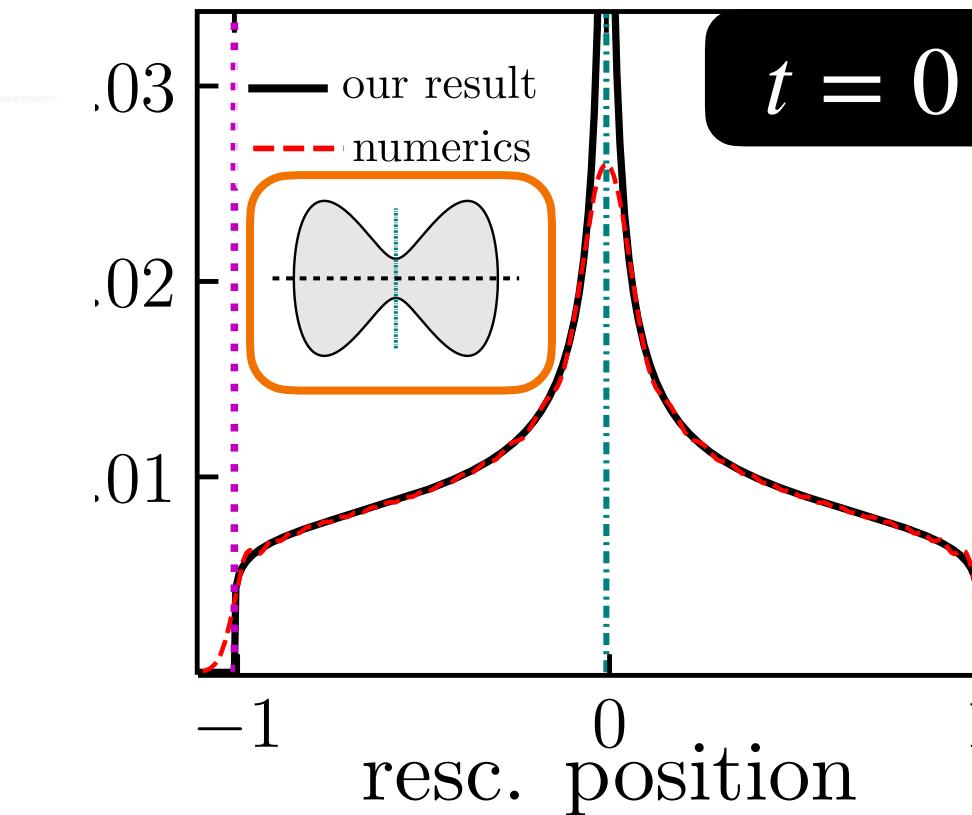
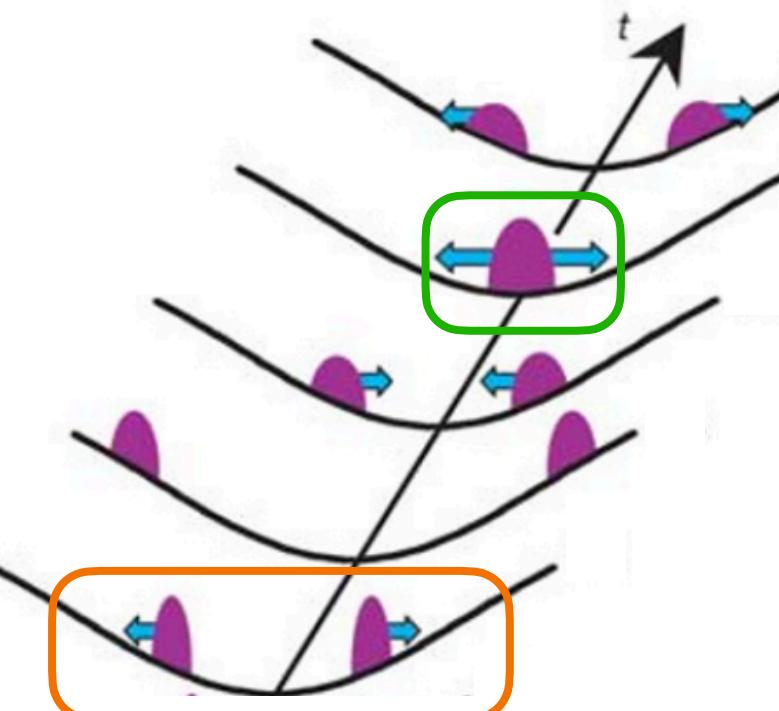
$$g_1(t; x, x') = \sum_{\vec{\nu}, \vec{\nu}'} [\mathbf{G}_t(x)]_{\vec{\nu}}^* [\mathbf{F}_t(x, x')]_{\vec{\nu}, \vec{\nu}'} [\mathbf{G}_t(x')]_{\vec{\nu}'}$$

$$[\mathbf{F}_t(x, x')]_{\vec{\nu}, \vec{\nu}'} = \frac{\prod_{a < b} |2 \sin(\frac{s_a - s_b}{2})|^{\nu_a \nu_b} \prod_{c < d} |2 \sin(\frac{s'_c - s'_d}{2})|^{\nu'_c \nu'_d}}{\prod_{j=1}^{2Q} \prod_{j'=1}^{2Q'} |2 \sin(\frac{s_j - s'_{j'}}{2})|^{\nu_j \nu'_{j'}}$$

$$[\mathbf{G}_t(x)]_{\vec{\nu}} = \frac{1}{\sqrt{2}} \left(\frac{G(3/2)^2}{\sqrt{\pi}} \right)^Q \prod_{j=1}^{2Q} \left| \frac{ds_j}{dx} \right|^{1/8} e^{i\nu_j \Phi(s_j)} \prod_{a < b} |\theta_F(s_a) - \theta_F(s_b)|^{\nu_a \nu_b}$$

✓ measured in experiments

[Kinoshita et al, Nature 440, 900 (2006)]



Ready-to-use:

$$(x, x'; t) \rightarrow \{s_a\}, \{s'_a\}$$

Before: Only harmonic traps

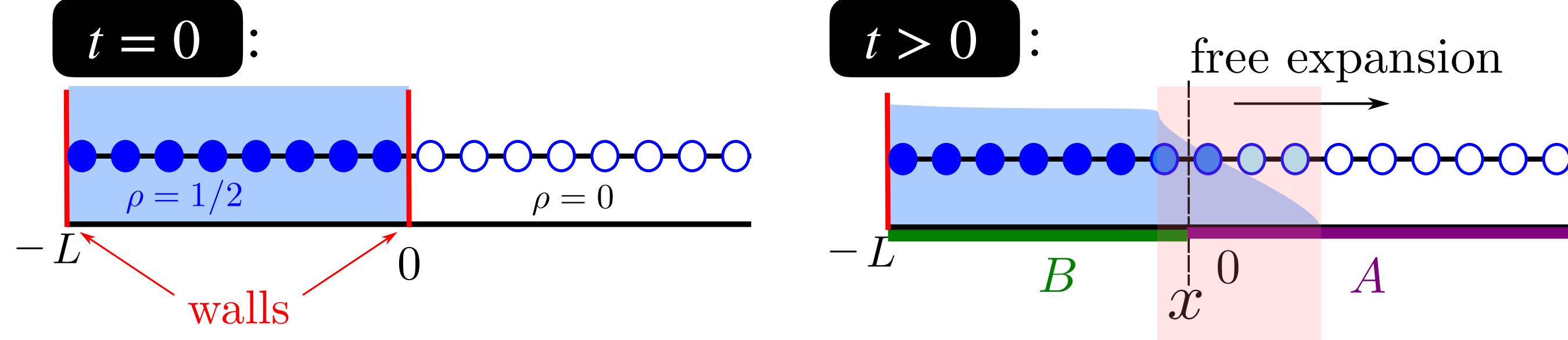
✗ [Forrester-Frankel-Garroni-Witte, 2000]

✗ [Minguzzi-Gangardt, 2006]

Entanglement dynamics



Scopa et al, 2105.05054



bi-partition of the chain at position x : $S_1(t, x) = -\text{tr} \hat{\rho}_A \log \hat{\rho}_A$

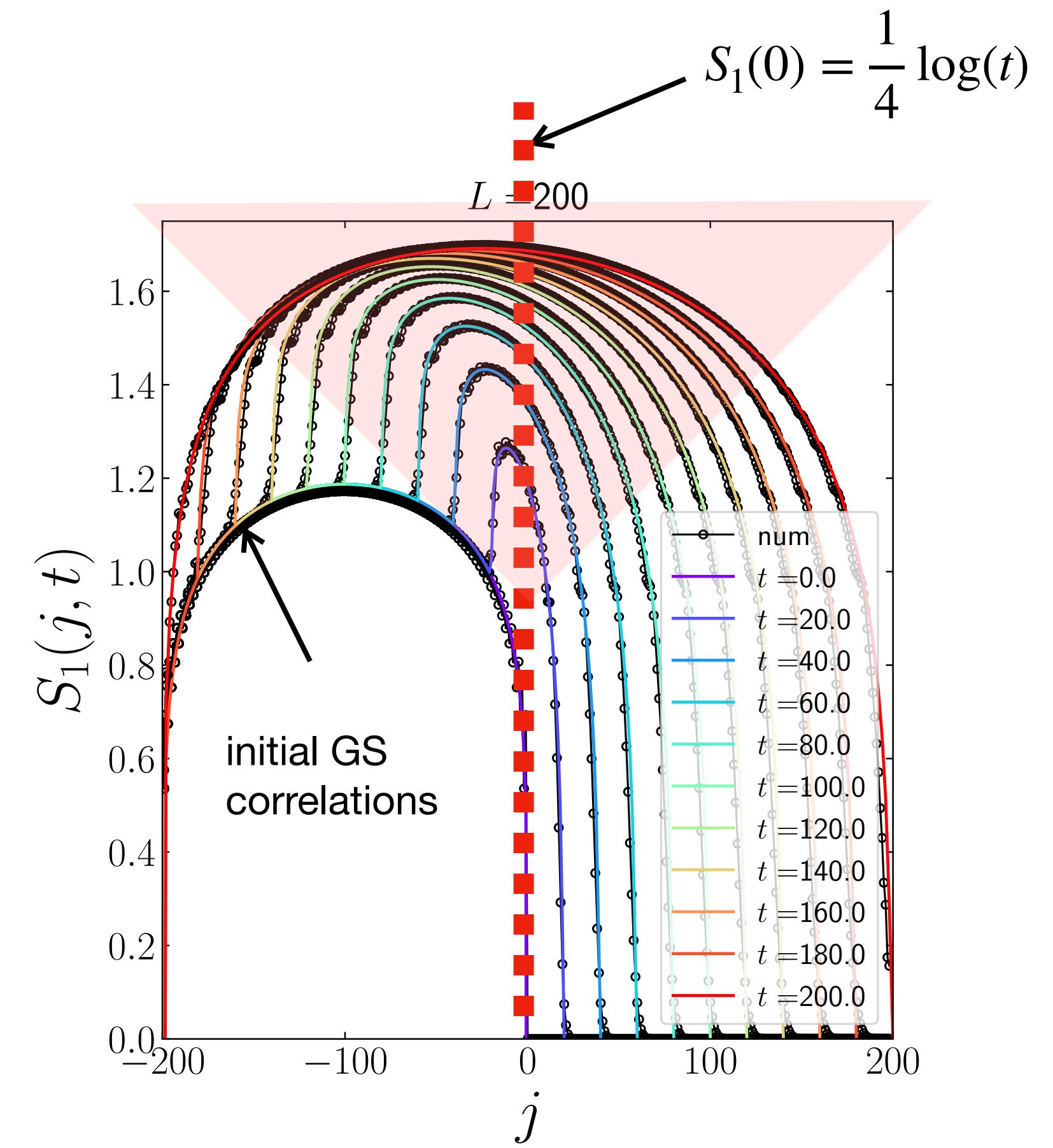
[Calabrese-Cardy, 2004]

$$S_1(t, x) = \lim_{n \rightarrow 1} \frac{1}{1-n} \log \left[\left| \prod_{\alpha} \frac{dx_t(s_{\alpha})}{ds} \right|^{-\Delta_n} \left\langle \prod_{\alpha} \hat{\tau}_n^+(s_{\alpha}) \hat{\tau}_n^-(s_{\alpha+1}) \right\rangle + f_{UV}(t, x) \right]$$

chiral twist fields

Fully analytical result:

$$S_1(|x| \leq t) = \frac{1}{6} \log \left(\frac{L}{2\pi} \sqrt{\left| \frac{x}{t} \left(1 - \frac{x^2}{t^2}\right) \right|} \left| \sqrt{1 + \sqrt{1 - \frac{x^2}{t^2}}} - \text{sgn}(x) \sqrt{1 - \sqrt{1 - \frac{x^2}{t^2}}} \sin \frac{\pi(x-t)}{2L} \right| \right) + \Upsilon$$



X Before: CFT-based conjectures

Summary and Conclusions

- ❖ Study the nonequilibrium dynamics of one-dimensional gases through hydrodynamic methods
 - **One-dimensional Bose gas:** quasi-integrable description → **Generalized Hydrodynamics**

- Quantitative results for local transport properties (density and currents)
- Non-local correlations coming from quantum fluctuations

- ❖ **Quantum Fluctuating theory** on top of the hydrodynamic background

- One-body correlations and Momentum distribution out of equilibrium
- Entanglement dynamics

Thank You!