

PRECISION FIELD THEORY FOR CONDENSED MATTER SYSTEMS

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Outline

- 1 Introduction
- 2 Precision anomalous elasticity in flat membranes
- 3 New fixed point in quenched disordered flat membranes
- 4 Metal-insulator transition in graphene and super-graphene
- 5 (Bonus) Precision optical conductivity in graphene and super-graphene
- 6 Outro

Introduction

TOPIC: **Perturbative** and **non-perturbative** approaches to **condensed matter systems**.

EXPERTISE: Higher-orders renormalization-group techniques.

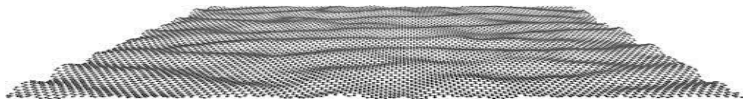
BACKGROUND:

- 2025: Post-Doctoral position - LAPTh - Annecy
- 2023-2024: Post-Doctoral position - INPAC - Shanghai Jiao Tong University
- 2020-2023: PhD thesis - LPTHE - SU - Supervisor: S. Teber

TODAY'S GOALS:

- Precision interaction effects
- Benchmark less controlled approaches
- **New physics beyond leading order**

TODAY'S SUBJECT: higher-order **elastic** and **electronic** interactions in membranes.



Outline

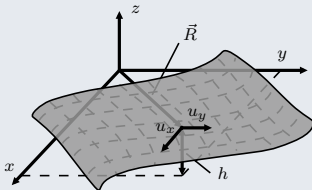
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Model – Flat crystalline membranes

Low-temperature action: A scalar derivative field theory

[Nelson, Peliti '87]

In-plane (phonon) $\vec{u}(\vec{x})$ and out-of-plane (flexuron) $\vec{h}(\vec{x})$ fields



$$S = \int d^d x \left[\frac{\kappa}{2} (\partial^2 h_\alpha)^2 + \mu T_{ab}^2 + \frac{\lambda}{2} T_{aa}^2 \right],$$

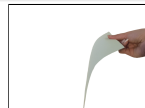
$$\text{Strain tensor} \equiv T_{ab} \approx \frac{1}{2} (\partial_a u_b + \partial_b u_a + \partial_a h_\alpha \partial_b h_\alpha)$$

$\kappa \equiv$ Bending rigidity, $\mu \equiv$ Shear modulus, $\lambda \equiv 2^{\text{nd}}$ Lamé coef.

Renormalization: Anomalous elasticity exponent η

Flexuron: $\langle hh \rangle \sim p^{-(4-\eta)}$, Lamé coeffs: $\lambda \sim \mu \sim p^{4-d-2\eta}$,

Phonon: $\langle uu \rangle \sim p^{-(6-d-2\eta)}$, Bending rigidity: $\kappa \sim p^{-\eta}$.



No deformations ($\eta = 0$)



With deformations ($\eta > 0$)

X-ray scattering experiments:

$$\eta_{\text{(blood-cells)}}^{\text{[Schmidt et al. '93]}} = 0.70(20),$$

$$\eta_{\text{(graphene)}}^{\text{[Lopez-Polin et al. '15]}} \approx 0.82,$$

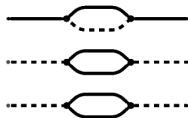
$$\eta_{\text{(amphiphilic-films)}}^{\text{[Gourier, Daillant '97]}} = 0.70(20),$$

$$\eta_{\text{(fly-egg)}}^{\text{[Jackson et al. '23]}} \approx 0.8.$$

1 & 2 loop approach

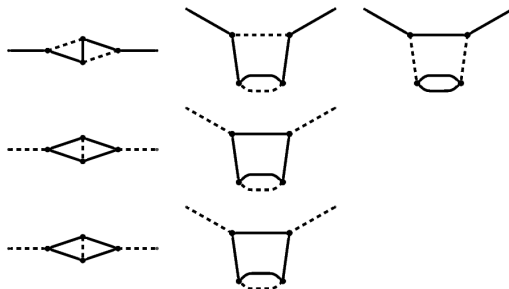
One loop: 3 diagrams, ~ 10 integrals, can be done by hand

[Aronovitz, Lubensky '88]



Two-loops: 7 diagrams, ~ 1000 integrals, need some automatization

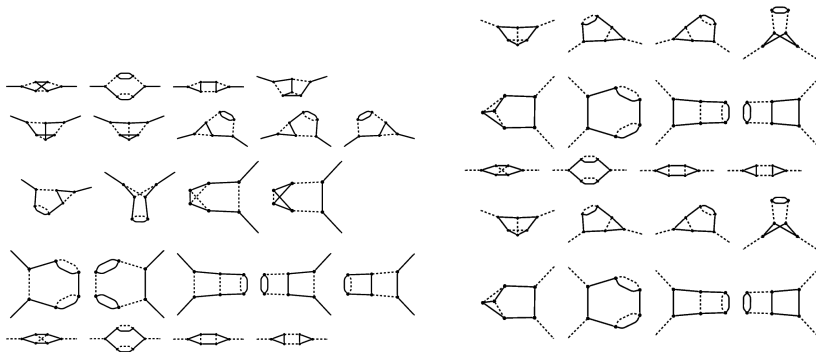
[Coquand, Mouhanna, Teber '20]



3 loop approach

Three loops: 42 diagrams, ~200000 integrals, need full automatization

[SM, Mouhanna, Teber '21]

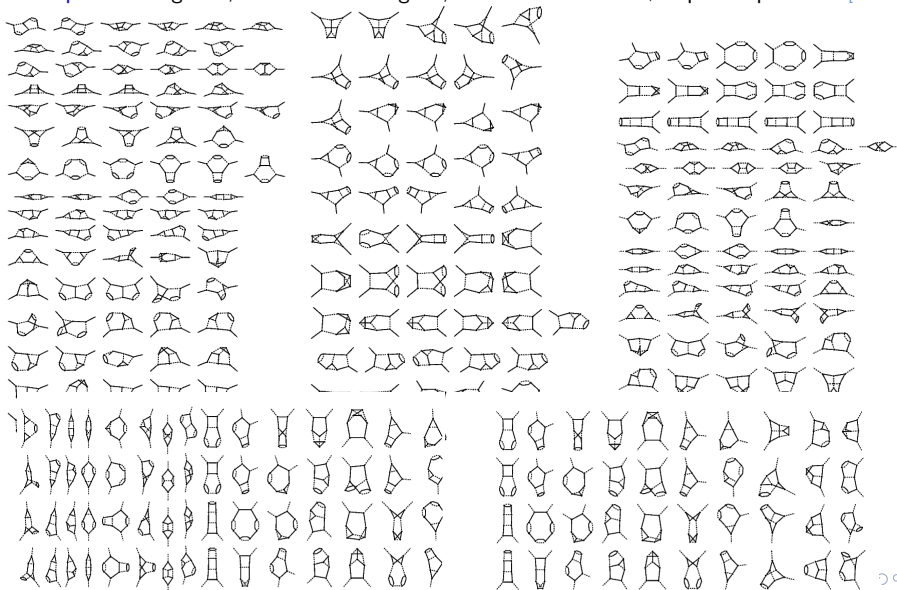


Highly automated computation using software from high-energy physics:
Qgraf [Nogueira '93], Fire [Smirnov '16], LiteRed [Lee '12], Form [Vermaseren '89] ...

4 loop approach

Four loops: 337 diagrams, ~60 million integrals, full automatization + supercomputer

[SM '24]



Results

Analytical result at the stable fixed point in $d = 4 - 2\varepsilon$:

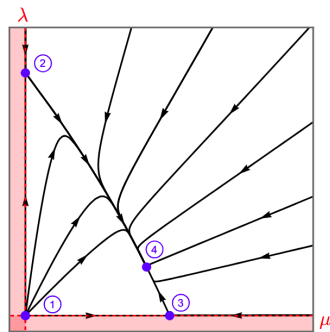
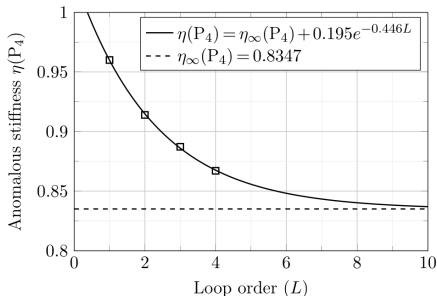
[SM '24]

$$\eta = \underbrace{\frac{24\varepsilon}{25}}_{\text{1-loop}} - \underbrace{\frac{144\varepsilon^2}{3125}}_{\text{2-loop}} - \underbrace{\frac{4(1286928\zeta_3 - 568241)\varepsilon^3}{146484375}}_{\text{3-loop}} - \underbrace{\frac{4(139409079893 + \dots)\varepsilon^4}{54931640625}}_{\text{4-loop}} + O(\varepsilon^5)$$

Numerically:

$$\eta = 0.96\varepsilon - 0.046\varepsilon^2 - 0.027\varepsilon^3 - 0.020\varepsilon^4 + \dots \underset{\varepsilon \rightarrow 1}{\approx} \mathbf{0.867}$$

Smaller and smaller corrections! (unexpected)



Benchmark less controlled methods

NPRG and SCSA provide non-perturbative results:

$$\text{NPRG: } \frac{(d-1)(d+8)(d-\eta+4)(4-d-2\eta)}{2\eta(d-\eta+8)} - d_c = 0 \quad [\text{Kownacki, Mouhanna '09}]$$

$$\text{SCSA: } \frac{d(d-1)\Gamma(2-\eta)\Gamma(2-\eta/2)\Gamma(\eta/2)\Gamma(\eta+d)}{\Gamma(2-\eta-d/2)\Gamma((4-\eta+d)/2)\Gamma((\eta+d)/2)\Gamma(\eta+d/2)} - d_c = 0 \quad [\text{Le Doussal, Radzihovsky '92}]$$

And non-perturbative answers $\eta_{\text{NPRG}} = 0.849$ and $\eta_{\text{SCSA}} = 0.821$. Can we trust them?

We can benchmark numerically in $d = 4 - 2\epsilon$:

$$\eta_{4\text{-loop}} = 0.96\epsilon - 0.046\epsilon^2 - 0.027\epsilon^3 - 0.020\epsilon^4 + \mathcal{O}(\epsilon^5), \quad [\text{SM '24}]$$

$$\eta_{\text{NPRG}} = 0.96\epsilon - 0.037\epsilon^2 - 0.027\epsilon^3 - 0.018\epsilon^4 + \mathcal{O}(\epsilon^5), \quad [\text{Kownacki, Mouhanna '09}]$$

$$\eta_{\text{SCSA}} = 0.96\epsilon - 0.048\epsilon^2 - 0.028\epsilon^3 - 0.018\epsilon^4 + \mathcal{O}(\epsilon^5), \quad [\text{Le Doussal, Radzihovsky '92}]$$

They are numerically **extremely successful** for this model!

Literature – A 30 years old story

	η	Method	Year/ref
Simulations ↓	≈ 0.66	Monte Carlo (membrane)	1990 Abraham, Nelson
	0.667	Large D (LO)	1988 Gutter, <i>et al.</i>
	≈ 0.7	Monte Carlo (vesicles)	1991 Komura, Baumgärtner
	0.72 ± 0.04	Monte Carlo (membrane)	1989 Leibler, Maggs
	0.75 ± 0.05	Monte Carlo (membrane)	1990 Gutter <i>et al.</i>
	0.750(5)	Monte Carlo (membrane)	1996 Bowick <i>et al.</i>
	0.789	SCSA (large- d_c NLO, semi-numerical)	2009 Gazit
	0.795(10)	Monte Carlo (graphene)	2013 Tröster
	0.81(3)	Monte Carlo (membrane)	1993 Zhang <i>et al.</i>
	≈ 0.82	Molecular dynamics simulations	1996 Zhang <i>et al.</i>
	≈ 0.82	SCSA (LO, semi-numerical)	2010 Zakharchenko <i>et al.</i>
	0.821	SCSA (LO, analytical)	1992 Le Doussal, Radzihovsky
Theory ↑	0.835	1 to 4-loop (extrapolation)	2024 SM
	0.849	NPRG (analytical)	2009 Kownacki, Mouhanna
	≈ 0.85	NPRG (semi-numerical)	2009 Braghin, Hasselmann
	≈ 0.85	Monte Carlo (graphene)	2009 Los <i>et al.</i>
	0.867	4-loop	2021 Pikelner
	0.887	3-loop	2021 SM <i>et al.</i>
	0.9 ± 0.04	Molecular dynamics simulations	1993 Petsche, Grest
	0.914	2-loop	2020 Coquand <i>et al.</i>
	0.960	1-loop	1988 Aronovitz, Lubensky
	1	Mean field	1987 Nelson, Peliti

Recalling $\eta_{(\text{exp.})} \approx 0.8$ [Schmidt *et al.* '93; Gourier *et al.* '97; Lopez-Polin *et al.* '15; Jackson *et al.* '23]]

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Model: Quenched disordered membranes

The quenched disordered action

[Nelson, Radzihovsky '91]

Introduce local random curvature and local random stress. Take **quenched** Gaussian distribution with zero mean. Use **replica** trick $\Rightarrow$ fields promoted with **replica** indices $A = 1, \dots, n \rightarrow 0$.

$$S = \int d^d x \left[\frac{\kappa}{2} (\partial^2 h_\alpha^A)^2 + \frac{\lambda}{2} (T_{aa}^A)^2 + \mu (T_{ab}^A)^2 + \Delta_\kappa \partial^2 h_\alpha^A \partial^2 h_\alpha^B + \frac{\Delta_\lambda}{2} T_{aa}^A T_{bb}^A + \Delta_\mu T_{ab}^A T_{ab}^B \right].$$

Content: 2 fields, 6 couplings, 3 indices type ... **challenging**.

Renormalization

With disorder, correlation functions become tensorial in the **replica** indices:

$$\langle h^A h^B \rangle \sim \delta^{AB} p^{-(4-\eta)} + J^{AB} p^{-(4-\eta-\phi)},$$

$$\langle u^A u^B \rangle \sim \delta^{AB} p^{-(6-d-2\eta)} + J^{AB} p^{-(6-d-2\eta-2\phi)}.$$

There is now an extra disorder exponent ϕ :

$\phi > 0 \equiv$ Temperature dominates (Clean phase),

$\phi < 0 \equiv$ Disorder dominates (**Glassy phase**),

$\phi = 0 \equiv$ Temperature and disorder coexist (**Marginal phase**).

Compute η in the glassy and marginal phase!

Results: Exponents (1 & 2)

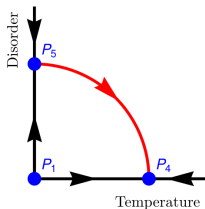
At **one loop**:

[Neslon Radzihovsky '91]

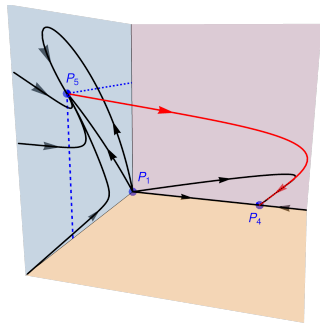
- A single disordered fixed point ($\phi=?$): $\eta_5 = \frac{6\varepsilon}{7} + O(\varepsilon^2)$

and ϕ is not resolved properly...

The RG-flow is weird:



Disorder seems irrelevant...



At **two loops**: Some progress, but the authors could not resolve the fixed points properly. [Coquand Mouhanna '21].

Results: Exponents (3)

At **three loop**, the fixed points are finally resolved

[SM, Mouhanna '22]

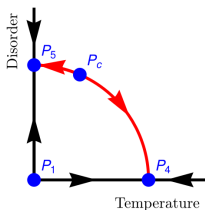
- Disorder dominated fixed point ($\phi < 0$):

$$\eta_5 = \frac{6\varepsilon}{7} - \frac{3629\varepsilon^2}{24010} - \frac{(698184144\zeta_3 - 759884263)\varepsilon^3}{823543000} + O(\varepsilon^4) \underset{\varepsilon \rightarrow 1}{\approx} \mathbf{0.610}$$

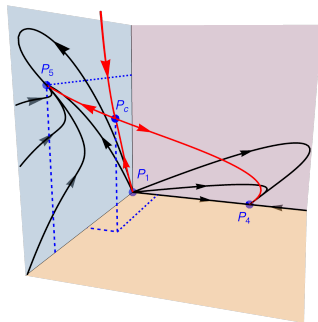
- New critical fixed point** ($\phi = 0$), finite disorder finite temperature glassy phase:

$$\eta_c = \frac{6\varepsilon}{7} - \frac{507\varepsilon^2}{3430} - \frac{(9504432\zeta_3 - 10463737)\varepsilon^3}{10084200} + O(\varepsilon^4) \underset{\varepsilon \rightarrow 1}{\approx} \mathbf{0.614}$$

New physics beyond LO. Disorder is now relevant:



But P_c is still marginally unstable (in contradiction with NPRG), and ϕ is still not resolved properly at P_5 ...
calls for 4-loop to settle everything.



Literature comparison

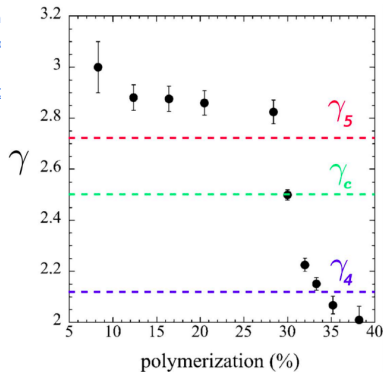
$\varepsilon = 1$	This work			Other approaches		
	1-loop	2-loop	3-loop	NPRG	SCSA	large- d_c
η_5	0.857 ^a	0.706 ^b	0.610 ^c	0.449 ^d	0.449 ^e	0.228 ^f
η_c	0.857 ^a	0.709 ^b	0.614 ^c	0.492 ^d		
η_4	0.960 ^g	0.914 ^h	0.887 ⁱ	0.849 ^j	0.821 ^k	1.68 ^l

The NPRG is the only other approach to distinguish P_5 and P_c !

SCSA (and large- d_c) can't distinguish $P_{c/5}$ at LO.

^a[Morse, Lubensky, 92'], ^b[Coquand, Mouhanna, 21'], ^c[Metayer, Mouhanna, 18'], ^e[Radzihovsky, Le Doussal, 92'], ^f[Saykin, Kachorovskii, Burmistrov, Teber, 20'], ⁱ[Metayer, Mouhanna, Teber, 21'], ^j[Kownacki, Doussal, Radzihovsky, 92'], ^l[Saykin, Gornyi, Kachorovskii, Burmistrov, 21']

New fixed point P_c observed in partial polymerization experiments [Chaieb *et al.* '06] ($\gamma = 3 - \eta$), first seen theoretically in NPRG [Coquand, Essafi, Kownacki, Mouhanna, 18'].



Conclusion

Takeway:

- Apparently convergent ε -series.
- Good benchmark for NPRG, SCSA, large- d_c ...
- $\exists$ a new non-trivial, finite-temperature, finite-disorder phase transition towards a glassy phase in polymerized membranes.

Ongoing projects:

- ϕ not completely resolved even at 3-loop, **4-loop needed** for a final answer.
- **SCSA beyond LO** might distinguish η_5 and η_c

Model perspectives:

- Crumpled-to-flat phase transition model
- Surface growth models (KPZ)
- Smectic-glass transition in liquid crystals

Techniques perspectives:

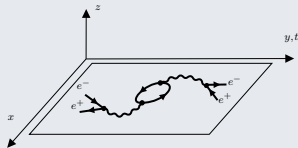
- Auxiliary field technique, big shortcut?
- NPRG
- Automatic NLO SCSA package

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Model for electronic interaction in graphene

Model – QED in three dimensions



Massless (S)QED₃ in large- N_f :

$$S = \int d^3x \left[i\bar{\psi} \not{D} \psi - \frac{1}{4} F_{\mu\nu}^2 \right] + \text{GF}_{(\xi)} + \text{SUSY}_{(\mathcal{N}=1)},$$

N_f electrons flavors (ψ) coupled (e) to photons (A^μ)

Naively superrenormalizable $[e] = 1$, but **non-trivial IR fixed point in the large- N_f limit** $\Rightarrow$
renormalizable with dimless coupling $1/N_f$

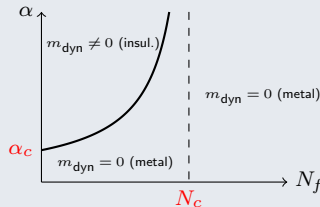
[Appelquist, Pisarski '81]

Dynamical electron mass generation

Renormalization of the (s)electron mass:

$$m_\psi \sim p^{1+\gamma_{m_\psi}}, \quad m_\phi \sim p^{1+\gamma_{m_\phi}}$$

Search for N_c . **Non-perturbative** effect!



Renormalization and dynamical mass generation

Non-perturbative effect, need to solve the SD equations self-consistently:

$$-i\Sigma(p) = \text{diagram} = \int [d^d_e k] \frac{\Gamma^\mu(k,p) \Gamma_0^{\nu} D_{\mu\nu}^{(0)}(p-k)}{(\not{k} - \Sigma(k))(1 - \Pi(p-k))},$$

$$i\Pi^{\mu\nu}(p) = \text{diagram} = -N_f \int [d^d_e k] \text{Tr} \frac{\Gamma^\mu(k, k-p) \Gamma_0^\nu}{(\not{k} - \Sigma(k))(\not{k} - \not{p} - \Sigma(k-p))},$$

$$\Gamma^\mu(p_1, p_2) = \text{diagram} = -ie\gamma^\mu - ie\Lambda^\mu(p_1, p_2),$$

$$\Lambda^\mu(p_1, p_2) = \text{diagram} = -N_f \int [d^d_e k] \frac{\Gamma^\mu(k-p_1, k-p_2) K(p_1, k-p_2, k-p_1, p_2)}{(\not{k} - \not{p}_1 - \Sigma(k-p_1))(\not{k} - \not{p}_2 - \Sigma(k-p_2))},$$

Messaging the equations around Σ , we conjecture the simple all-order **gap equation**: [SM, Teber '21]

$$(1-b)b = (1-\gamma_{m_\psi})\gamma_{m_\psi}, \quad \text{with} \quad m_{\text{dyn}} = \Sigma(p \rightarrow 0) \sim p^{-b}$$

Which depends only on $\gamma_{m_\psi} \implies$ precision needed.

We need to go beyond LO in large- N_f expansion...

Large- N_f approach

NLO (s)electron **self-energies**, 34 diagrams, ~ 500 branchut integrals

[SM, Teber '22]

$$\begin{aligned} \Sigma_\psi = & \text{diagram 1} + \text{diagram 2} + \text{diagram 3} + \text{diagram 4} + \text{diagram 5} \\ & + \text{diagram 6} + \text{diagram 7} + \text{diagram 8} + \text{diagram 9} + \text{diagram 10} \\ & + \text{diagram 11} + \text{diagram 12} + \text{diagram 13} + \text{diagram 14} + \text{diagram 15} \\ & + \text{diagram 16} + \text{diagram 17} + \text{diagram 18} + \text{diagram 19} + \text{diagram 20} + O(1/N_f^3) \\ \Sigma_\phi = & \text{diagram 21} + \text{diagram 22} + \text{diagram 23} + \text{diagram 24} + \text{diagram 25} \\ & + \text{diagram 26} + \text{diagram 27} + \text{diagram 28} + \text{diagram 29} + \text{diagram 30} \\ & + \text{diagram 31} + \text{diagram 32} + \text{diagram 33} + \text{diagram 34} + O(1/N_f^3) \end{aligned}$$

Results – (S)QED₃

$$\text{In QED}_3: \gamma_{m_\psi} = \frac{32}{3\pi^2 N_f} - \frac{64(28 - 3\pi^2)}{9\pi^4 N_f^2} + O(1/N_f^3)$$

[Gracey '93-'94]

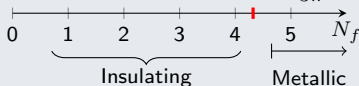
$$\text{In SQED}_3: \gamma_{m_\psi} = \gamma_{m_\phi} = \frac{4}{\pi^2 N_f} - \frac{4(14 - \pi^2)}{\pi^4 N_f^2} + O(1/N_f^3)$$

[SM, Teber '22]

Using the gap equation we derive N_c at leading order:

QED₃:

$$N_c = \frac{128}{3\pi^2} = 4.32$$



SQED₃:

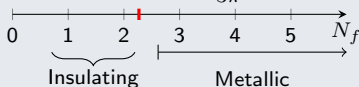
$$N_c = \frac{16}{\pi^2} = 1.62$$



And beyond leading order? N_c reduces drastically $\Rightarrow$ New physics.

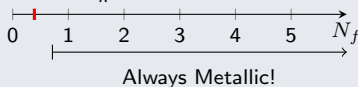
QED₃:

$$N_c = \frac{2(4 + 3\pi^2)}{3\pi^2} = 2.27$$



SQED₃:

$$N_c = \frac{\pi^2 - 6}{\pi^2} = 0.39$$



Literature – N_c in QED₃, a 40 years old debate!

In the literature, seemingly all values for N_c has been found, from 0 to ∞ ...

N_c in QED ₃	Method	Year
∞	SD (LO)	1984 Pisarski
∞	SD (non-perturbative, Landau gauge)	1990, 1992 Pennington <i>et al.</i>
∞	RG study	1991 Pisarski
∞	lattice simulations	1993, 1996 Azcoiti <i>et al.</i>
< 4.4	F-theorem	2015 Giombi <i>et al.</i>
$(4/3)(32/\pi^2) = 4.32$	SD (LO, resummation)	1989 Nash
4.422	RG study (1-loop) ($N_c^{\text{conf}} \approx 6.24$)	2016 Janssen
4	functional RG ($4.1 < N_c^{\text{conf}} < 10.0$)	2014 Braun <i>et al.</i>
$3 < N_c < 4$	RG study	2001 Kubota, Terao
3.5 ± 0.5	lattice simulations	1988, 1989 Dagotto <i>et al.</i>
3.31	SD (NLO, Landau gauge)	1993 Kotikov
3.29	SD (NLO, Landau gauge)	2016 Kotikov <i>et al.</i>
$32/\pi^2 \approx 3.24$	SD (LO, Landau gauge)	1988 Appelquist <i>et al.</i>
$3.0084 - 3.0844$	SD (NLO, resummation)	2016 Kotikov, Teber
2.89	RG study (1-loop)	2016 Herbut
2.85	SD (NLO, resummation, $\forall \xi$)	2016 Gusynin <i>et al.</i> Kotikov <i>et al.</i>
$1 + \sqrt{2} = 2.41$	F-theorem	2016 Giombi <i>et al.</i>
2.27	Effective gap eq. (NLO, double resummation, $\forall \xi$)	2022 SM, Teber
$< 9/4 = 2.25$	RG study (1-loop)	2015 Di Pietro <i>et al.</i>
$< 3/2$	Free energy constraint	1999 Appelquist <i>et al.</i>
$1 < N_c < 4$	lattice simulations	2004 Hands <i>et al.</i> 2008 Strouthos <i>et al.</i>
0	SD (non-perturbative, Landau gauge)	1990 Atkinson <i>et al.</i>
0	lattice simulations	2015, 2016 Karthik, Narayanan

Recent results converges towards $N_c \in [2,3] \dots$

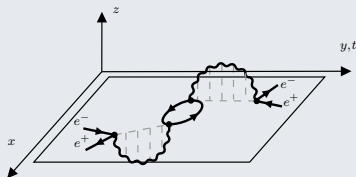
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Model – Electronic interactions in graphene

(S)QED in mixed dimensions

[Gorbar, Gusynin, Miransky '01]

Massless (S)QED_{4,3}:

$$S = i \int d^3x \bar{\psi} \not{D} \psi - \frac{1}{4} \int d^4x F_{\mu\nu}^2 + \text{GF}_{(\xi)} + \text{SUSY}_{(\mathcal{N}=1)},$$

N_f electrons flavors (ψ) in 3-dim coupled (e) to photons (A^μ) in 4-dim

Renormalizable with non-running dimensionful coupling $\alpha = e^2/4\pi$.

[Gorbar, Gusynin, Miransky '02]

Optical conductivity

The photon (perpendicular) propagator renormalize as

$$\langle A^\mu A^\nu \rangle_\perp = \frac{i}{2p} \frac{P_\perp^{\mu\nu}}{1 - \Pi_\gamma}, \quad \text{with} \quad \Pi_\gamma = -\frac{\pi N_f \alpha}{4} [1 + C_\gamma \alpha + O(\alpha^2)].$$

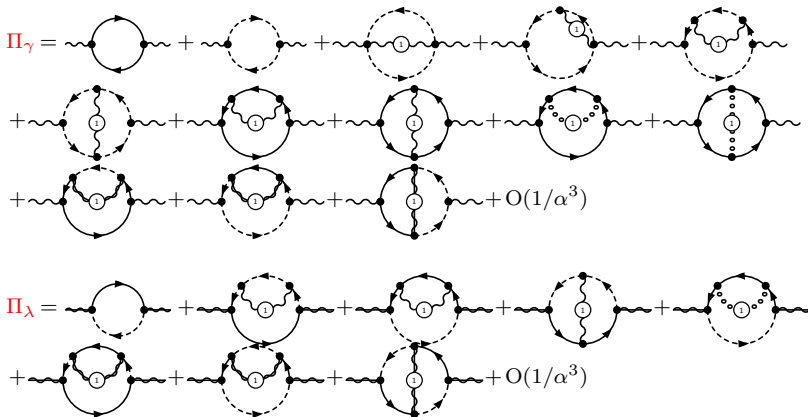
The photon polarization Π_γ is finite & gauge-independent $\Rightarrow$ physical! Need C_γ for precision.

Optical conductivity (Kubo formula):

$$\sigma(\omega) = -p \times \Pi_\gamma = \sigma_0 (1 + C_\gamma \alpha + O(\alpha^2))$$

Multi-loop approach

Two-loop photon and photino **polarizations**: 21 diagrams, ~ 500 branchut integrals [\[SM, Teber '21\]](#)



Result: $C_\gamma = \frac{92 - 9\pi^2}{18\pi} \approx 0.06$ and $C_\gamma^{\text{SUSY}} = \frac{12 - \pi^2}{2\pi} \approx 0.34$ both very small corrections!

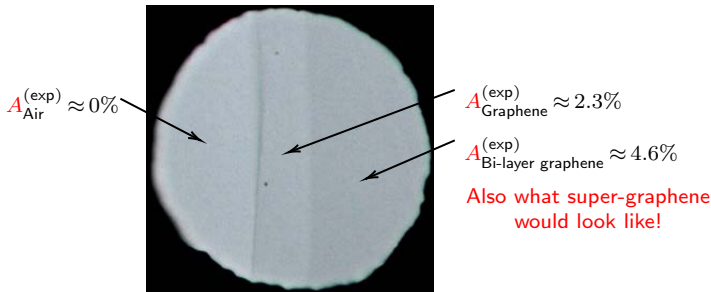
Result

The optical conductivity is directly related to the **universal optical absorbance** (A): [SM, Teber '23]

$$A_{\text{graphene}} = \pi\alpha(1 + \alpha C_\gamma + \dots) = (2.293 \pm 0.002)\%$$

$$A_{\text{super-graphene}} = 2\pi\alpha(1 + \alpha C_\gamma^{\text{SUSY}} + \dots) = (4.59 \pm 0.15)\%$$

Interestingly, **optical measurements** provides [Nair *et al.* '08]:



- SUSY **enhance** the optical absorbance.
- Bi-layer and super-graphene have **similar optical absorbance**.

Conclusion

Takeway:

- In fermionic QEDs, dynamical matter mass generation is **possible** for small N_f
- SUSY **strongly suppress** dynamical matter mass generation
- SUSY **enhances the optical absorbance** of planar materials

Ongoing projects:

- Higher orders for QCD cusp anomalous dimension matrix

Models perspectives:

- Dynamical mass generation in QCD
- Josephson junction, sine-Gordon model
- ...

Technical perspectives:

- Bootstrap methods
- Automatic NLO SD solving package

Outline

- 1 Introduction
- 2 Precision anomalous elasticity in flat membranes
- 3 New fixed point in quenched disordered flat membranes
- 4 Metal-insulator transition in graphene and super-graphene
- 5 (Bonus) Precision optical conductivity in graphene and super-graphene
- 6 **Outro**

Conclusion

Final takeaway:

- **Higher precision** for critical exponents is successful in many models
- Strong benchmark for less controlled methods
- **Higher orders** for non-perturbative methods is needed
- **Perturbative** input to access **non-perturbative** features
- **New physics beyond leading order!**

Thank you for your attention :)

Selection of seminars and posters:

- **Anomalous elasticity in polymerized membranes, analytical 4-loop result**
2024: Seminar, "journée des utilisateurs du supercalculateur MeSU" – SU
- **Electronic interaction effects in low-dimensional abelian field theories**
2024: Seminar, ShanghaiTech
- **Field theoretic approach to flat quenched disordered polymerized membranes**
2023: Seminar, International conference "48th Middle European Cooperation in Statistical Physics" (MECO48) Slovakia
- **Membranes elastic degrees of freedom, a multi-loop approach**
2022: Poster, international conference "47th Middle European Cooperation in Statistical Physics" (MECO47)
- **3-loop order approach to flat polymerized membranes**
2022: Seminar, "Journée de physique statistique" - ENS - France
- **Membranes elastic degrees of freedom, a multi-loop approach**
2021: Poster, international conference "Advanced Computing and Analysis Techniques in Physics Research" (ACAT21)
- **2-loop anomalous dimensions in reduced QED and dynamical mass generation**
2021: Poster, international conference "Relativistic Fermions in Flatland"

Publications

- **Field-theoretic approach to flat polymerized membranes**
[SM & S. Teber, 2025, arXiv:2412.18490]
- **Four-loop elasticity renormalization of low-temperature flat polymerized membranes**
[SM, EPL, 2024, 10.1209/0295-5075/ad949a]
- **Critical Properties of Three-Dimensional Many-Flavor QEDs**
[SM & S. Teber, Symmetry, 2023, 10.3390/sym15091806]
- **Electron mass anomalous dimension at $O(1/N_f^2)$ in 3D $\mathcal{N}=1$ supersymmetric QED**
[SM & S. Teber, PLB, 2022, 838(2023)137729]
- **Flat polymerized membranes at three-loop order**
[SM, D. Mouhanna & S. Teber, J. Phys. Conf. Ser., 2022, 2438(2023)1:012141]
- **The flat phase of quenched disordered membranes at three-loop order**
[SM & D. Mouhanna, PRE, 2022, 106(6):064114]
- **Three-loop order approach to flat polymerized membranes**
[SM, D. Mouhanna & S. Teber, PRE Letter, 2022, 105(1):L012603]
- **Two-loop mass anomalous dimension in RQED and dynamical fermion mass generation**
[SM & S. Teber, JHEP, 2021 2021(9):107]
- **3D $\mathcal{N}=1$ supersymmetric QED at large- N_f and applications to super-graphene**
[A. James, SM & S. Teber, 2021 arXiv:2102.02722]