



# Generalised Hydrodynamics of the KdV Soliton Gas.

Laboratoire de Physique Théorique de la Matière Condensée

Thibault Bonnemain, 18th February 2025

[Based off joint work with B. Doyon and G. El]

• Boltzmann 1868: micro-canonical ensemble in long time limit

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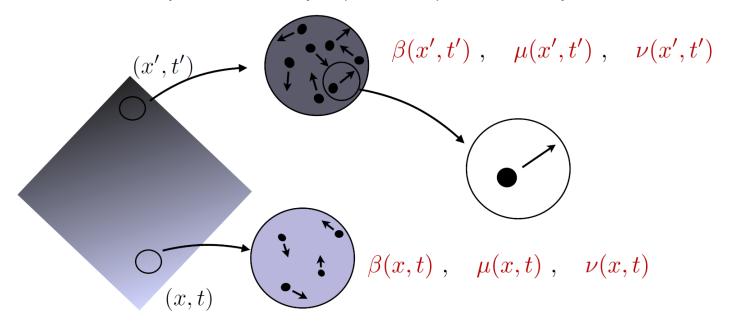
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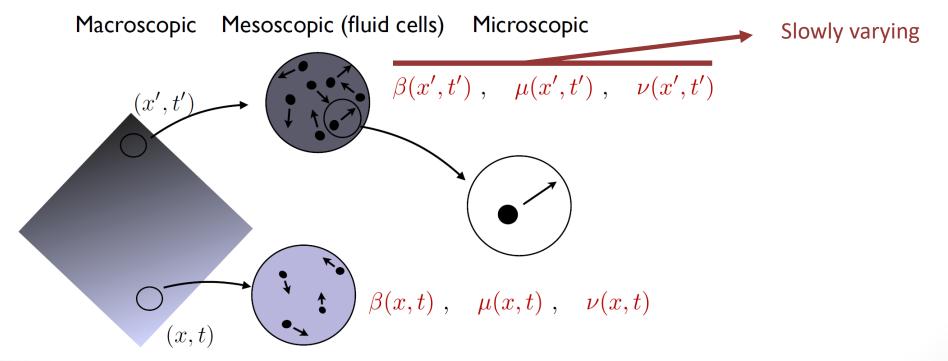
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 Gibbs ensembles (GE):  $\rho \propto \exp[-\beta (E - \mu N - \nu P)]$ 

• Hydrodynamic principle: separation of scales and propagation of local GE

Macroscopic Mesoscopic (fluid cells) Microscopic

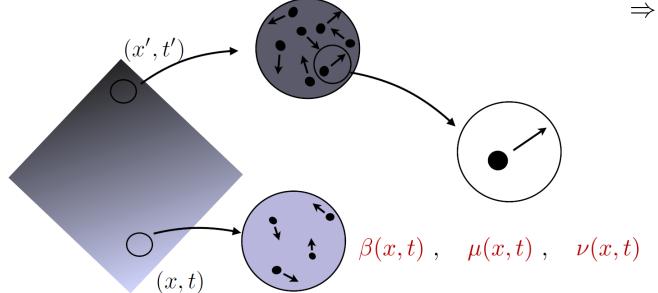


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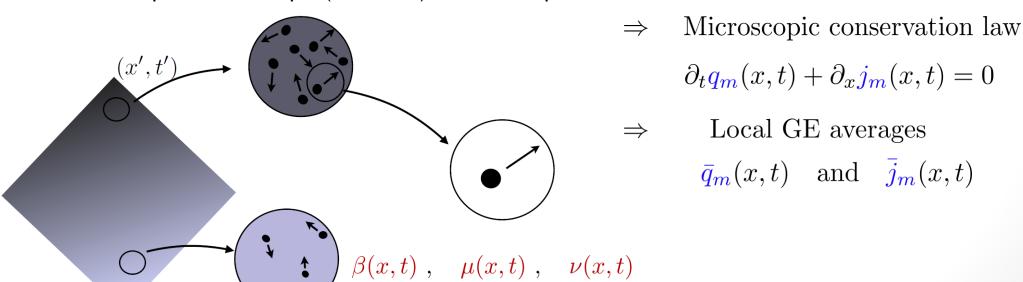


Microscopic conservation law

$$\partial_t \mathbf{q_m}(x,t) + \partial_x \mathbf{j_m}(x,t) = 0$$

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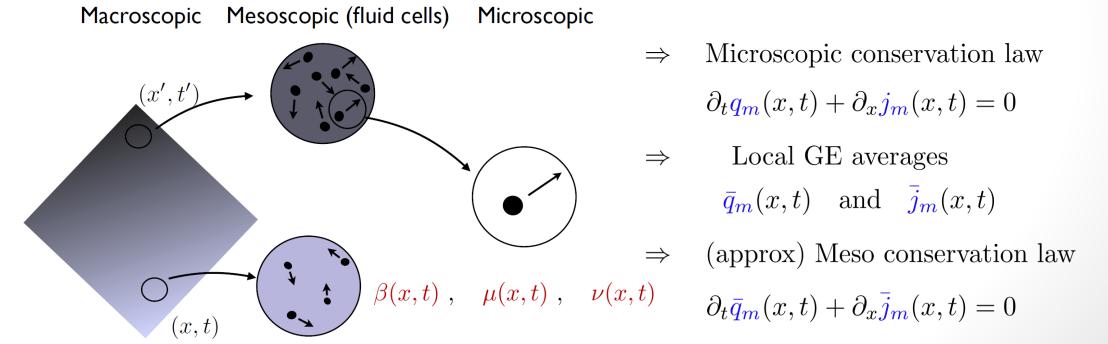
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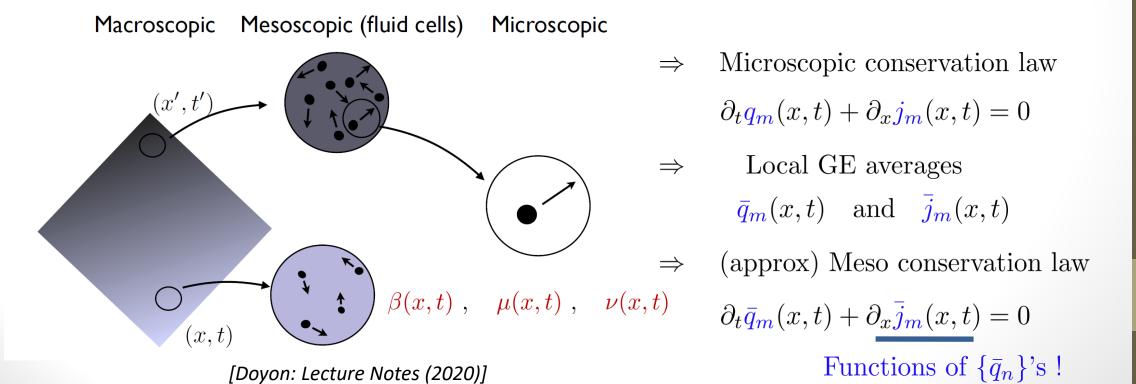
[Doyon: Lecture Notes (2020)]

(x,t)

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#### **Generalised Gibbs ensembles and GHD**

- Boltzmann 1868: micro-canonical ensemble in long time limit
  - $\hat{=}$  Generalised Gibbs ensembles (GGE):  $\rho \propto \exp \left[ -\sum_{n=0}^{\infty} \beta_n Q_n \right]$
- Hydrodynamic principle: separation of scales and propagation of local GGE

(x',t')  $\{\beta_n(x,t)\}$   $\{\beta_n(x,t)\}$ 

Macroscopic Mesoscopic (fluid cells) Microscopic

 $\Rightarrow$  Microscopic conservation law

$$\partial_t \mathbf{q}_m(x,t) + \partial_x \mathbf{j}_m(x,t) = 0$$

 $\Rightarrow$  Local GGE averages

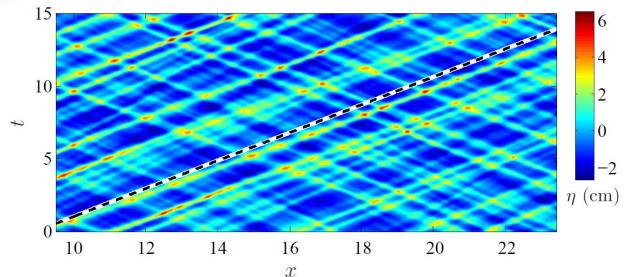
$$\bar{q}_m(x,t)$$
 and  $\bar{j}_m(x,t)$ 

 $\Rightarrow$  (approx) Meso conservation law

$$\partial_t \bar{q}_m(x,t) + \partial_x \bar{j}_m(x,t) = 0$$

Functions of  $\{\bar{q}_n\}$ 's!

#### Soliton gases in experiments

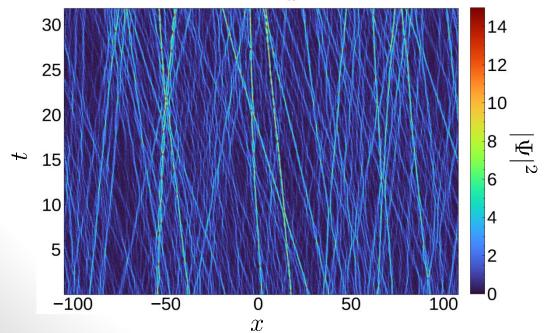


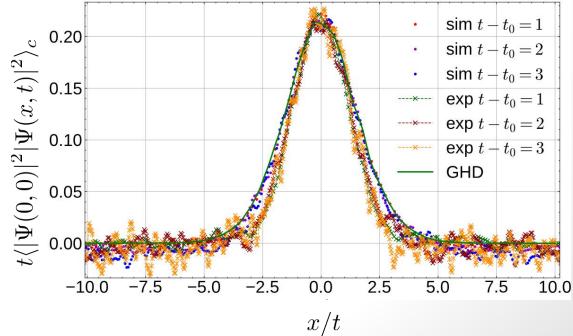
Left: Soliton gas in a water tank.

[Redor et al. (2019)]

Below: Soliton gas in an optical fiber and intensity correlations.

[Curtesy of Elias Charnay]





# The Korteweg-de Vries equation

• KdV: integrable, nonlinear, dispersive PDE

$$\partial_t u + 6u\partial_x u + \partial_x^3 u = 0.$$

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• Infinite set of conservation laws

Time conserved "charges" 
$$Q_n = \int \mathrm{d}x \; q_n(x,t) \;, \quad \text{and} \quad J_n = \int \mathrm{d}t \; j_n(x,t) \;, \quad \text{conserved conserved conserved formula}$$

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• Exactly solvable via Inverse Scattering Transform (IST).

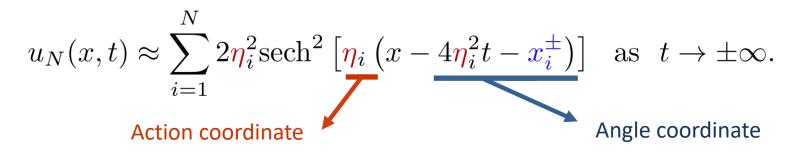
#### Some properties of N-soliton solutions

• Long time asymptotics of N-soliton solutions

$$u_N(x,t) \approx \sum_{i=1}^N 2\eta_i^2 \operatorname{sech}^2\left[\eta_i\left(x - 4\eta_i^2 t - x_i^{\pm}\right)\right] \text{ as } t \to \pm \infty.$$

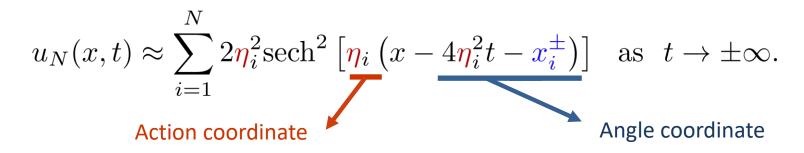
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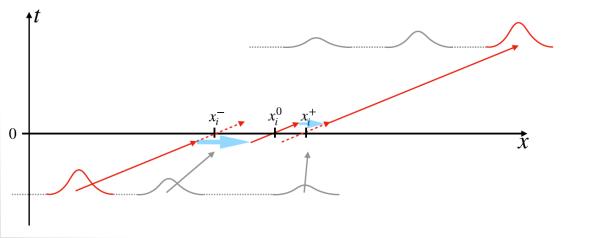


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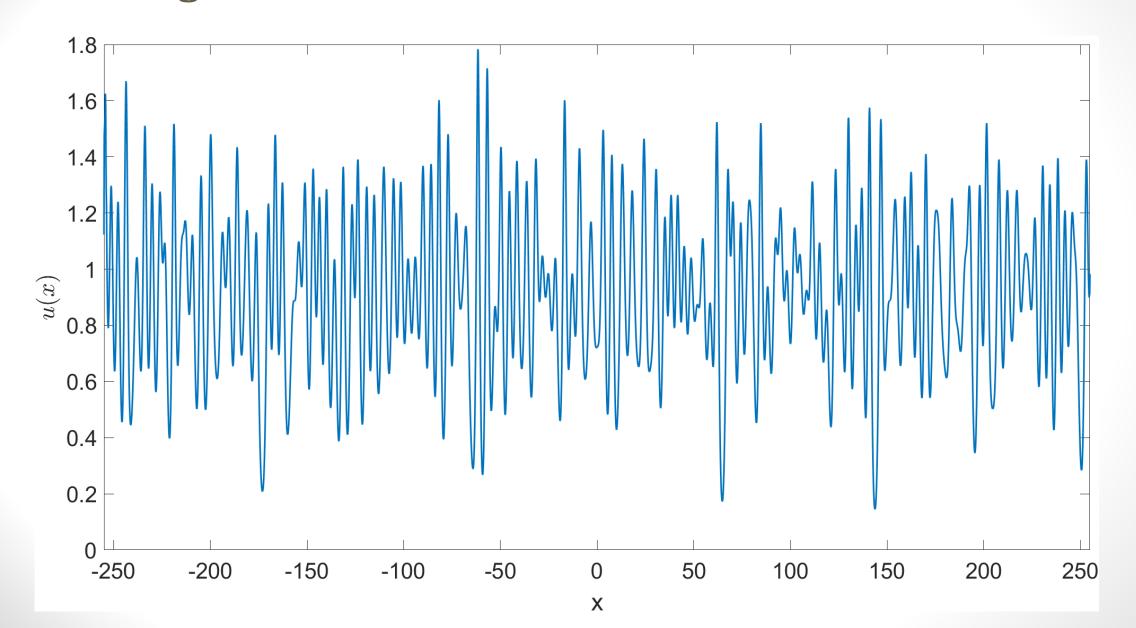


• Relation between asymptotic states given by scattering shift

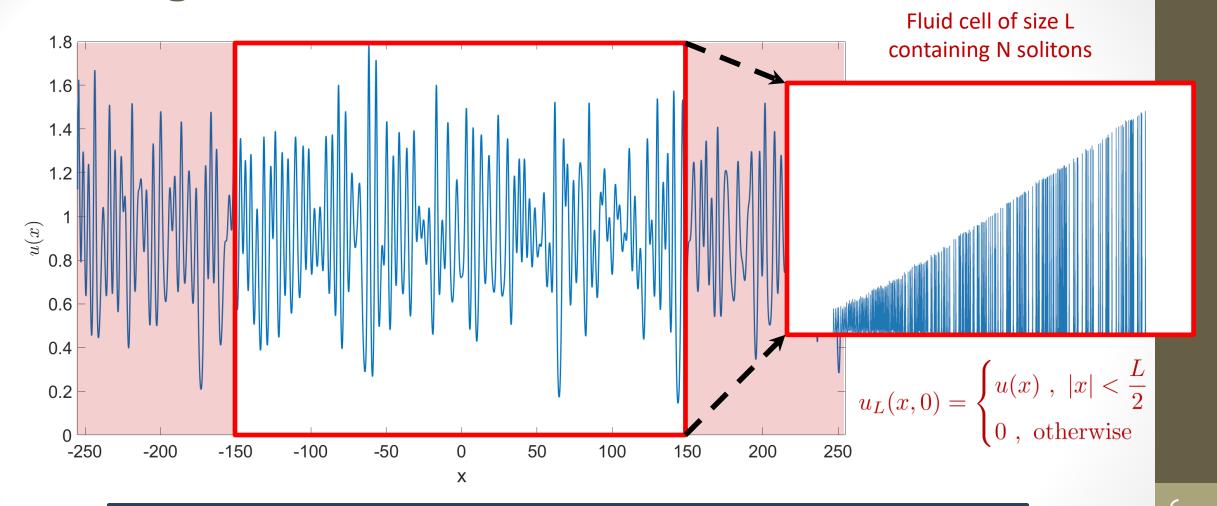


$$x_i^+ - x_i^- = \sum_i \frac{\operatorname{sgn}(\eta_i - \eta_j)}{\eta_i} \ln \left| \frac{\eta_i + \eta_j}{\eta_i - \eta_j} \right| .$$

# Soliton gas: basic idea and motivations



#### Soliton gas: basic idea and motivations



Asymptotically: 
$$u_L(x,t) \approx \sum_{i=1}^{N} 2\eta_i^2 \operatorname{sech}^2\left[\eta_i\left(x - 4\eta_i^2 t - x_i^{\pm}\right)\right]$$
 as  $t \to \pm \infty$ .

• Partition function

$$\mathcal{Z}_L = \sum_{N=0}^{\infty} \frac{1}{N!} \int_{\Gamma^N \times \mathbb{R}^N} \prod_{i=1}^N \frac{\mathrm{d}p(\eta_i)}{2\pi} \mathrm{d}x_i^- \exp\left[-\sum_{i=1}^N w(\eta_i)\right] \chi\left(u_N(x, t=0) < \epsilon_x, x \notin [0, L]\right)$$

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Soliton bare velocity

$$p(\eta) = 4\eta^2$$

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Soliton bare velocity

Generalised
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Constraint / Entropy

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ullet Thermodynamic limit  $L o \infty$ : large deviations theory [Varadhan (1966), Touchette (2009)]

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**Occupation function** 

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Occupation function

• Density of states  $\rho(\eta)$ 

$$\frac{\rho(\eta)}{\frac{n}{n}(\eta)} = \eta - \int_{\Gamma} d\mu \, \rho(\mu) \log \left| \frac{\eta + \mu}{\eta - \mu} \right|$$

 $\rho(\eta) d\eta dx = \# \text{ of solitons in } [x, x + dx] \times [\eta, \eta + d\eta]$ 

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Change of metric

• Occupation function

$$\log \mathbf{n}(\eta) = -w(\eta) + \int_{\Gamma} d\mu \mathbf{n}(\eta) \log \left| \frac{\eta - \mu}{\eta + \mu} \right|, \qquad \mathcal{F} = -\int_{\Gamma} \eta \mathbf{n}(\eta) d\eta,$$

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• Density of states

$$\rho(\eta) = \frac{\delta \mathcal{F}}{\delta w(\eta)} ,$$

• Thermodynamic averages

$$\langle q_n \rangle = \frac{\partial \mathcal{F}}{\partial \beta_n} = \int_{\Gamma} d\eta \ \rho(\eta) h_n(\eta) \ ,$$

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• Static covariance matrix

$$\mathsf{C}_{ab} \equiv \int_{\Gamma} \mathrm{d}x \left( \langle q_a(x) q_b(0) \rangle - \langle q_a(x) \rangle \langle q_b(0) \rangle \right) = -\frac{\partial^2 \mathcal{F}}{\partial \beta_a \partial \beta_b} ,$$

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	GHD	Simulations
$C_{00}^{\mathrm{DC}}$	0.0235	$0.022 \pm 0.003$
$C_{01}^{\mathrm{DC}}$	0.027	$0.024 \pm 0.004$
$C_{11}^{\mathrm{DC}}$	0.042	$0.039 \pm 0.005$
$C_{\mathrm{D}}^{00}$	0.22	$0.2 \pm 0.03$
$C_{01}^{\mathrm{U}}$	0.28	$0.23 \pm 0.04$
$C_{11}^{\mathrm{U}}$	0.39	$0.36 \pm 0.05$
$C^\mathrm{L}_{00}$	0.2	$0.2 \pm 0.01$
$C^{\mathrm{L}}_{01}$	0.25	$0.23 \pm 0.01$
$C_{11}^{L}$	0.36	$0.34 \pm 0.02$

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$$\langle o(x,t)\rangle \approx \langle o\rangle_{\{\beta_n(x,t)\}} \equiv \bar{o}_n(x,t)$$
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#### **Derivation of GHD equations**

[Based on: Doyon, Spohn, Yoshimura (2017)]

• Asymptotic dynamics

$$x_{j}^{-}(t) = x_{j}^{-}(0) + 4\eta_{j}^{2}t ,$$

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• Continuity equation for the DOS

$$\partial_t \rho(\eta; x, t) + \partial_x \left[ \rho(\eta; x, t) v^{\text{eff}}(\eta; x, t) \right] = 0$$
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