

## Comment on “Microscopic Theory of Network Glasses”

In a recent Letter, Hall and Wolynes [1] (HW) ask whether a microscopic theory of network glasses can be developed starting from a model of dense spherical fluids. To do so, they constrain the number of nearest neighbors and count their central force interactions separately. They obtain the dynamical transition temperature  $T_A$  (below which the system is nonergodic and the motion is landscape determined), and the entropy crisis (Kauzman) temperature  $T_K$  as functions of  $n_b$ , the average number of nearest neighbor bonds/atom. A Lindemann melting criterion on the amplitude of nearest neighbor vibrations defines the glass transition temperature  $T_G$ . The model shows that  $T_A/T_G$  and  $T_K/T_G$  monotonically increase and decrease, respectively, as functions of  $n_b$ . For  $n_b = 2.4$ ,  $T_A/T_G = 100$ , an unreasonable result. Mode coupling theory defines a critical nonergodicity temperature  $T_c$  beyond which a radical change in the long time limit of the density-density correlation function occurs.  $T_c$  has been plausibly estimated for vitreous silica in molecular dynamics simulations [2] as  $T_c/T_G = 2$ . Even this temperature is presently outside the reach of experimental investigations [2].

How do these results compare with experiments? There are standard procedures for extrapolating specific heats to obtain  $T_K$ , but identifying the onset of nonergodicity at  $T_A$  (or  $T_c$ ) is much more difficult. Chalcogenide glasses are ideal test systems because one can synthesize them over a wide range of  $n_b$  by chemical alloying group IV additives in Se base glass. Fortunately, glass transitions of these systems have been recently reinvestigated [3] using  $T$ -modulated scanning calorimetry (MDSC), a new method which permits separating the usual DSC heat flow endotherm  $\dot{H}_T$  into a reversing part  $\dot{H}_{rev}$  which is ergodic (and which follows the modulated  $T$  profile) from the nonreversing part  $\dot{H}_{nr}$  which is nonergodic (arising from underlying temperature dependent activated processes) as illustrated in Fig. 1(a). MDSC permits one to establish this temperature  $T_A$  at which dynamics become landscape dominated (i.e., in MDSC language dominated by a  $T$ -dependent  $\dot{H}_{nr}$ ) in contrast to the linear response regime (i.e., at high temperatures when the heat flow is  $\dot{H}_{rev}$  dominated by a constant activation energy).

In binary  $\text{Ge}_x\text{Se}_{1-x}$  glasses, observed variations in  $T_A/T_G$  and  $T_K/T_G$  as a function of  $n_b = 2 + 2x$  are compared to HW predictions in Fig. 1(b). The  $T_K/T_G(n_b)$  results were obtained from a Vogel-Fulcher analysis of viscosity measurements [4]. One finds the observed and predicted variations in  $T_K/T_G(n_b)$  ratio to be in reasonable accord with each other showing a general reduction starting from a value of about 0.9 at  $n_b = 2$  to a value of 0.6 near  $n_b = 2.7$ . Note, however, that the broad global minimum in the observed  $T_K/T_G$  ratio near  $n_b = 2.4$  is not reproduced by the HW approach. More serious is the fact that ob-

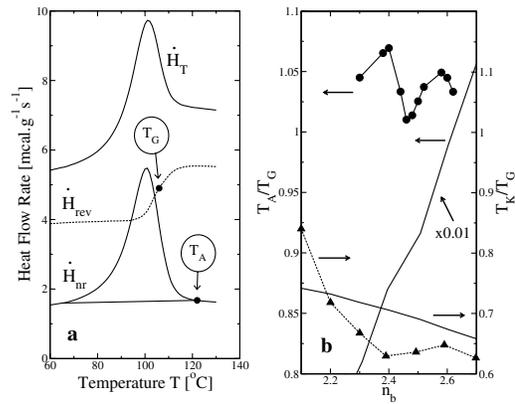


FIG. 1. (a) MDSC scan of  $\text{As}_{20}\text{Se}_{80}$  bulk glass;  $T_G$  is defined as the inflexion point of  $\dot{H}_{rev}$  while  $T_A$  is the end point of the  $\dot{H}_{nr}$  endotherm. (b)  $T_A/T_G$  (●) and  $T_K/T_G$  (▲) for Ge-Se glasses. Lines without symbols are HW predictions [1].

served  $T_A/T_G$  values are (i) 2 orders of magnitude lower than the HW predictions and (ii) show a global minimum near  $n_b = 2.4$  that is in sharp contrast to the almost linear increase [Fig. 1(b)] predicted by HW. The global minima in  $T_K/T_G$  and  $T_A/T_G$  ratios are features related to self-organization of glasses that are missing in the HW approach. Clearly, features such as inclusion of local structures [3,5], structural self-organization [5,6], and noncentral forces (bond bending) are missing in the theory. Bond-bending forces constitute  $(4n_b - 6)/(5n_b - 6)$  of the global number of network constraints (e.g., 0.6 at the ideal  $n_b = 2.4$ ) indicating that the noncentral (angular) forces have to be included in a successful theory of network glasses.

M. Micoulaut<sup>1</sup> and P. Boolchand<sup>2</sup>

<sup>1</sup>LPTL, Université Paris VI  
4 Place Jussieu 75252, Paris CEDEX 05, France

<sup>2</sup>Department of ECECS  
University of Cincinnati  
Cincinnati, Ohio 45221-0030, USA

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