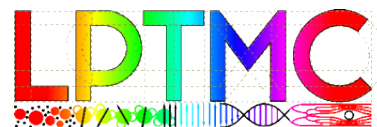


CECAM Workshop  
Mesoscopic Modeling in Physics of Molecular and Cell Biology  
Toulouse, october 10-13, 2016

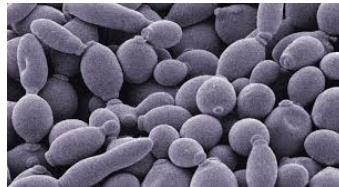
Finite-size scaling analysis of super-resolution  
imaging of epigenetic domains

Maria Barbi  
Antony Lesage, Jean-Marc Victor

LPTMC, University Pierre & Marie Curie  
and  
CNRS GDR 3536 « ADN »



# Chromatin organization



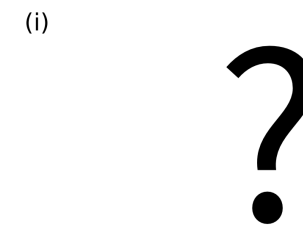
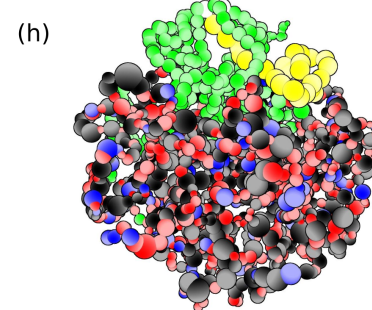
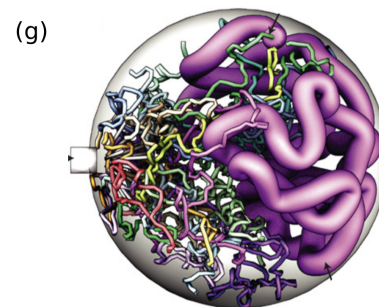
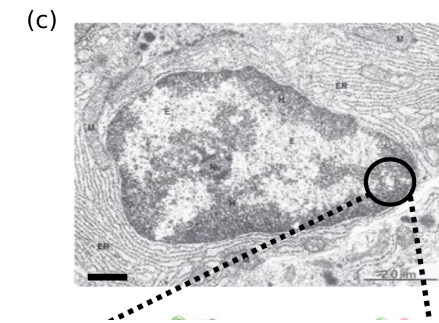
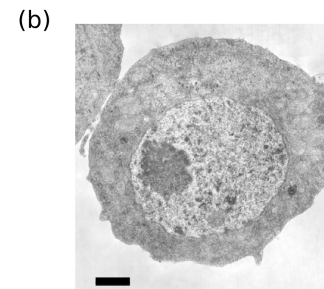
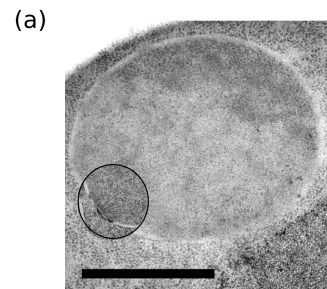
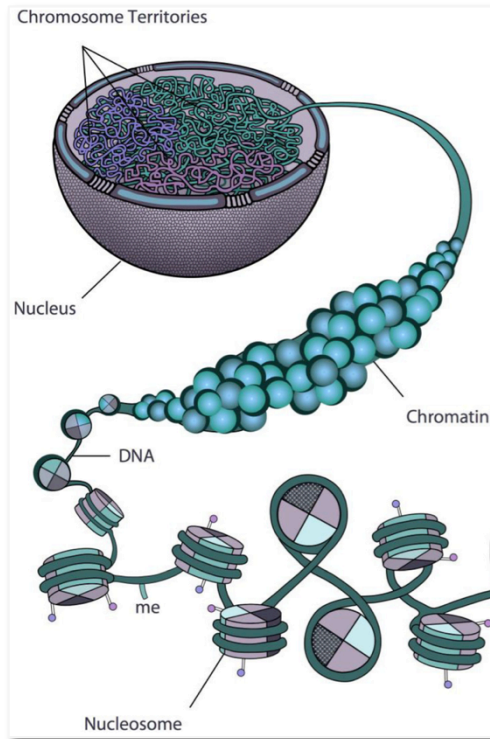
yeast



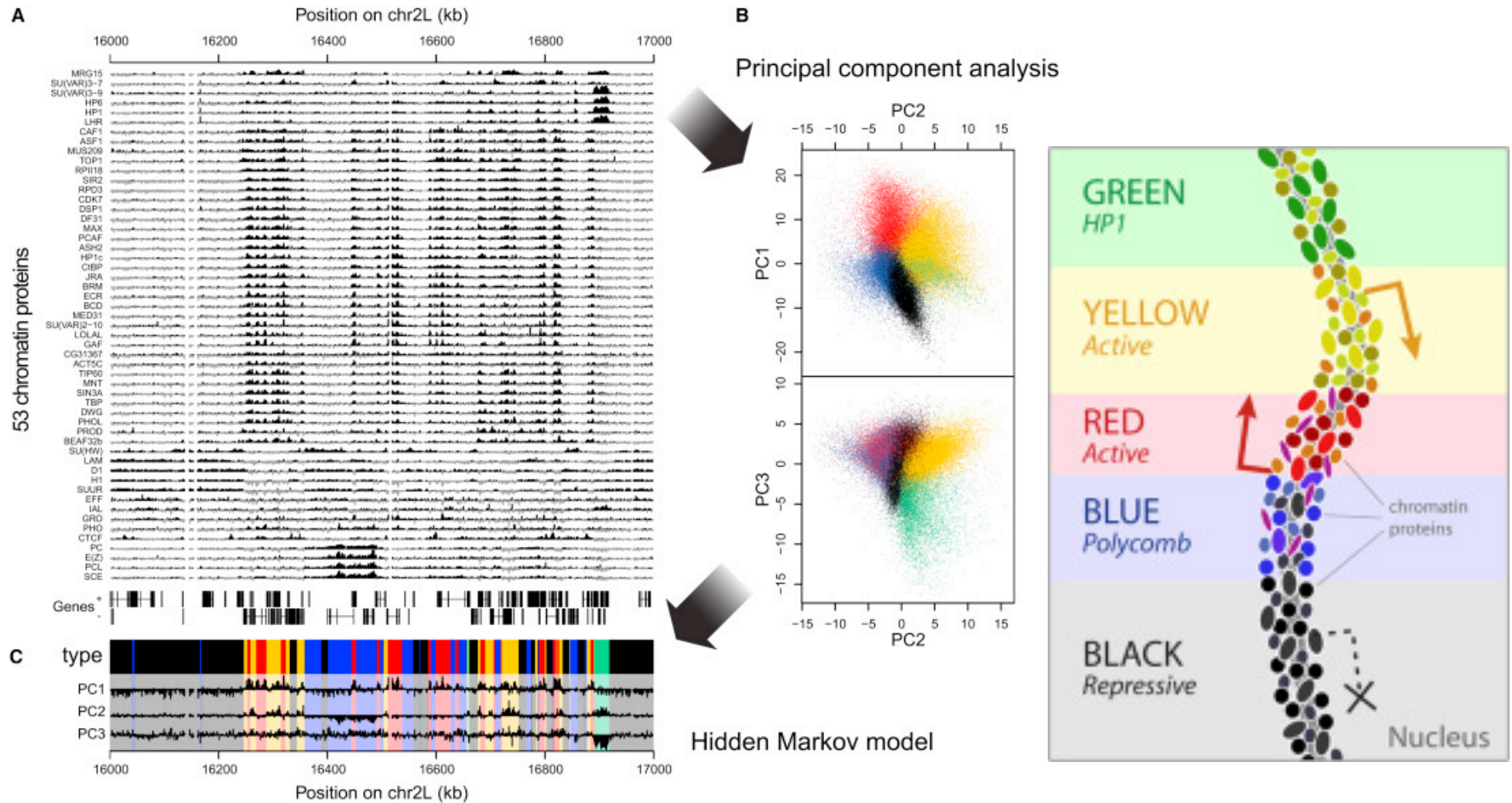
drosophila (fly)



human



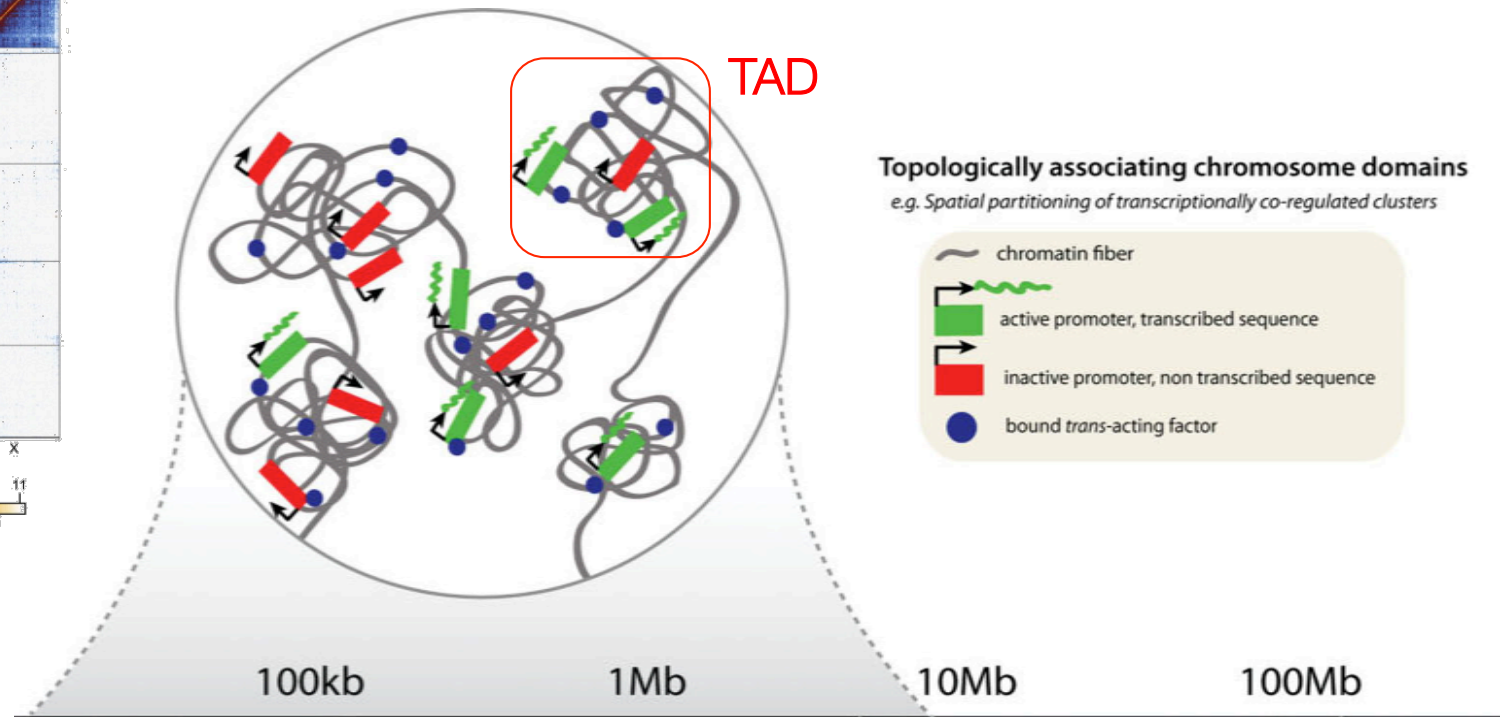
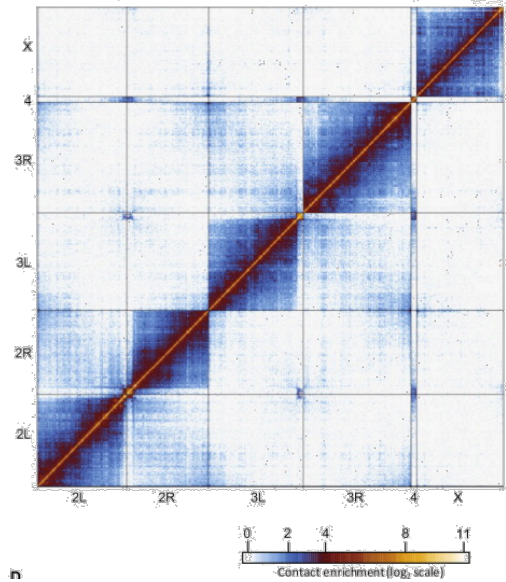
# Epigenetic colors



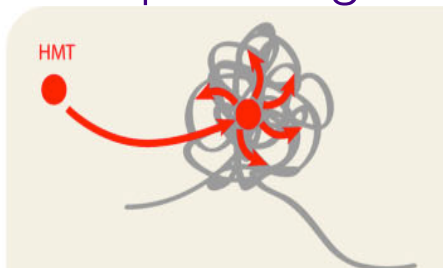
# « topologically associating domains » (TADs)

HiC contact frequency map

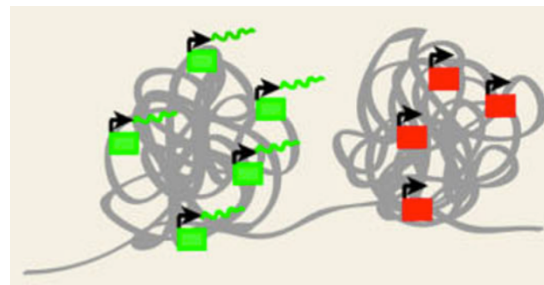
Sexton et al – Cell, 2012



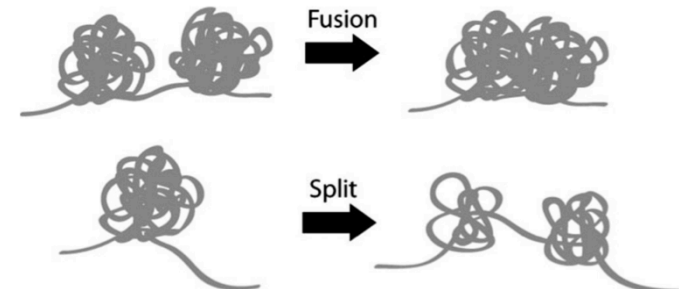
Epigenetic mark spreading



Transcription activity



Merge/split domains



Nora, Dekker, Heard - Bioessays, 2013

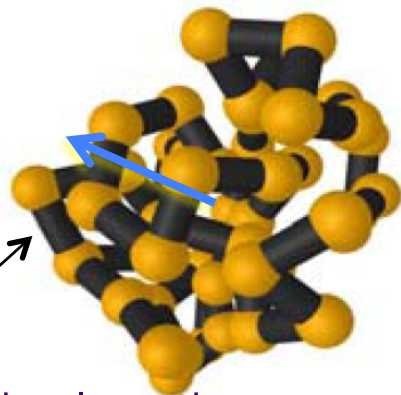
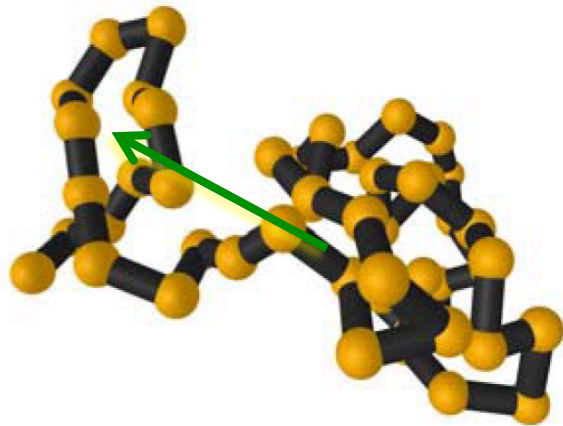
# Polymer models

Radius of gyration:

$$R_G^2 = \left\langle \frac{1}{2N} \sum_{i=1}^N \sum_{j=1}^N (r_i - r_j)^2 \right\rangle$$

Power law and Flory exponent:

$$R_G \sim N^\nu$$



( $K$  = Kuhn length = segment length)

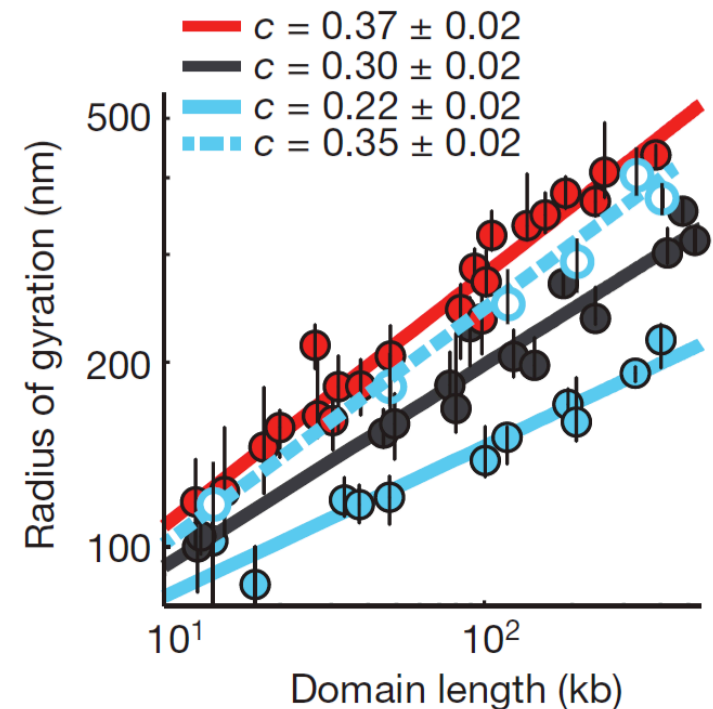
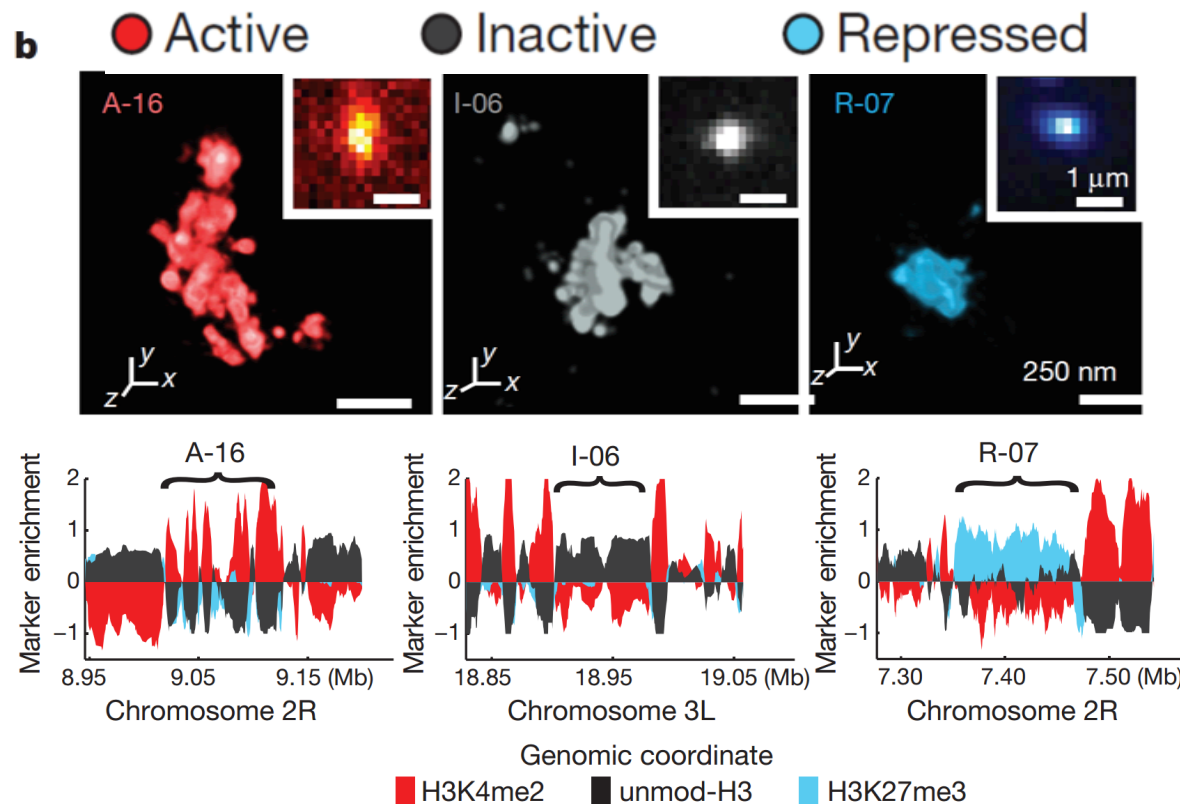
- Random Walk coil  $\nu = 1/2$
- Self-Avoiding Walk (SAW) coil  $\nu = 3/5$
- Equilibrium globule  $\nu = 1/3$
- Fractal globule  $\nu = 1/3$

## Super-resolution imaging reveals distinct chromatin folding for different epigenetic states

Alistair N. Boettiger<sup>1</sup>, Bogdan Bintu<sup>1</sup>, Jeffrey R. Moffitt<sup>1</sup>, Siyuan Wang<sup>1</sup>, Brian J. Beliveau<sup>2</sup>, Geoffrey Fudenberg<sup>3</sup>, Maxim Imakaev<sup>3</sup>, Leonid A. Mirny<sup>3</sup>, Chao-ting Wu<sup>2</sup> & Xiaowei Zhuang<sup>1</sup>

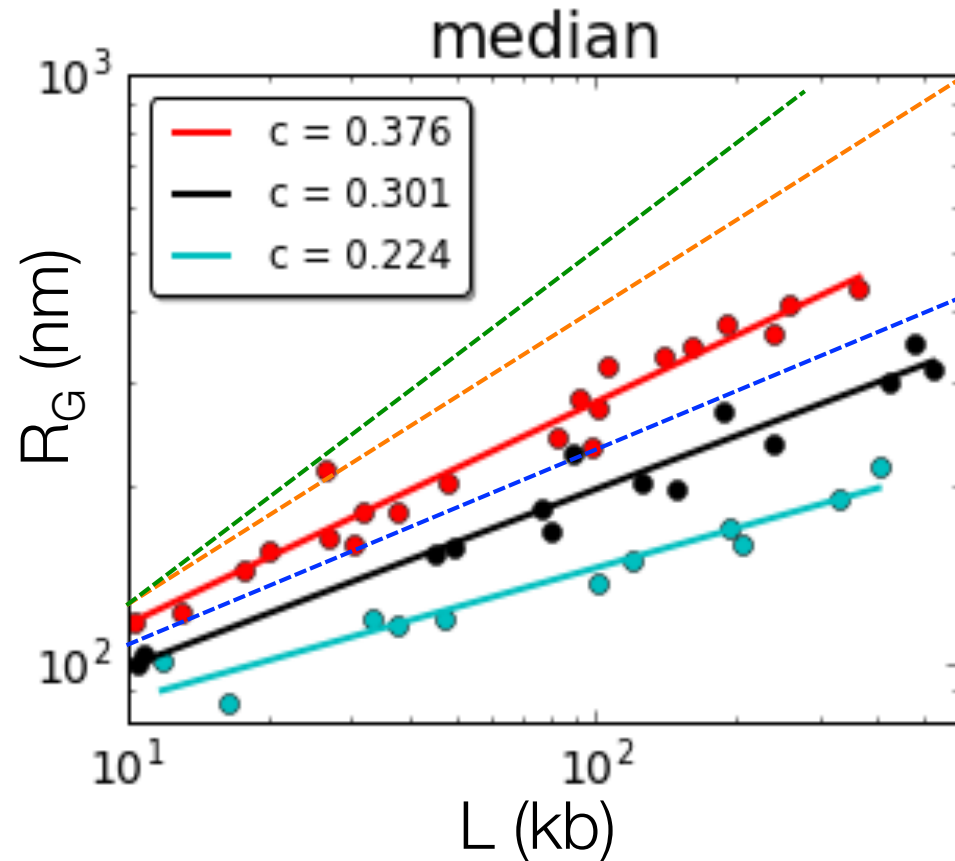
3 epigenetic states:

Rg scaling law :



# how to interpret the experimental scaling laws?

(Data (points) from Boettinger et al. 2016)



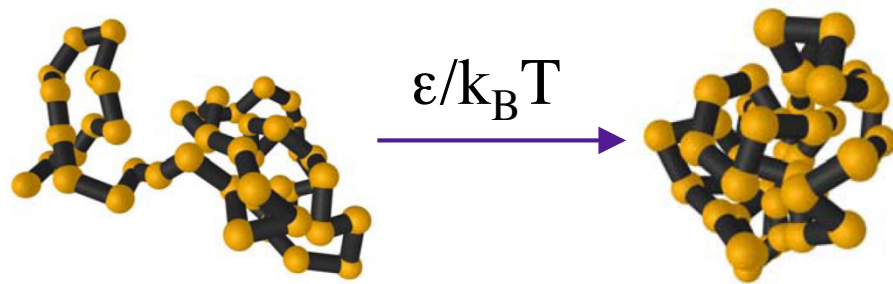
----- Random Walk  $\nu = 1/2$

----- SAW Coil  $\nu = 3/5$

----- Eq. Globule / Fractal Globule  $\nu = 1/3$

# The coil-globule transition: finite-size effects

A polymer with  $N$  identical monomers:



coil-globule transition

depending on the ratio between :

$\epsilon$  interaction energy per monomer

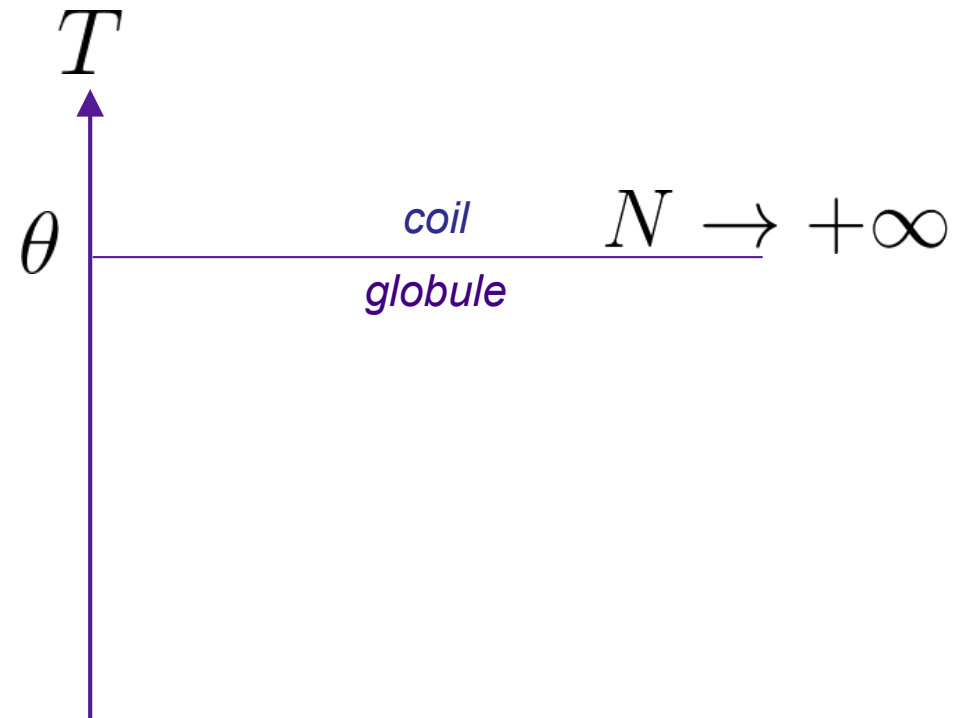
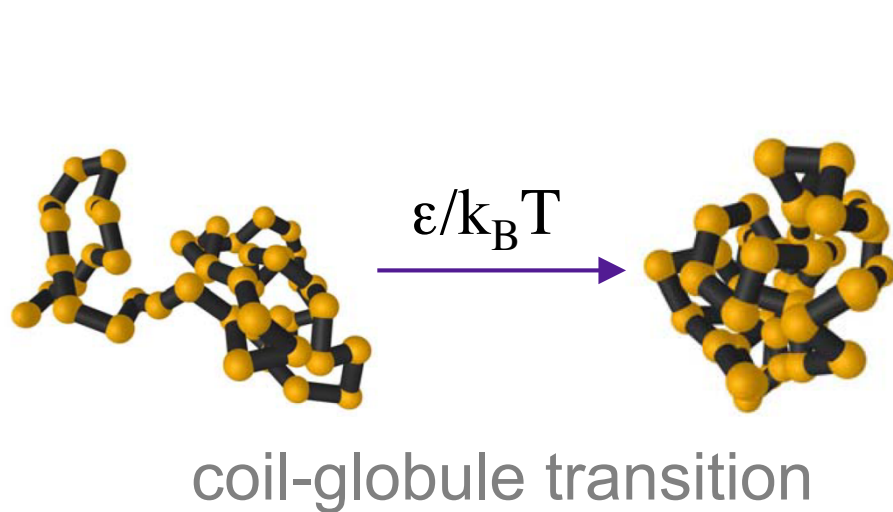
$k_B T$  thermal energy

( $K$  Kuhn length = segment length)



# The coil-globule transition: finite-size effects

A polymer with  $N$  identical monomers:



depending on the ratio between :

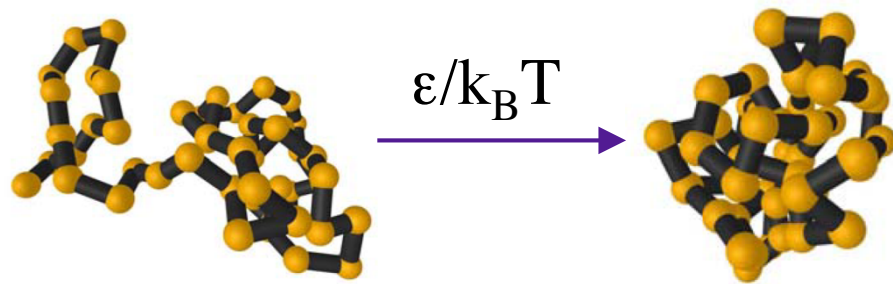
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# The coil-globule transition: finite-size effects

A polymer with  $N$  identical monomers:



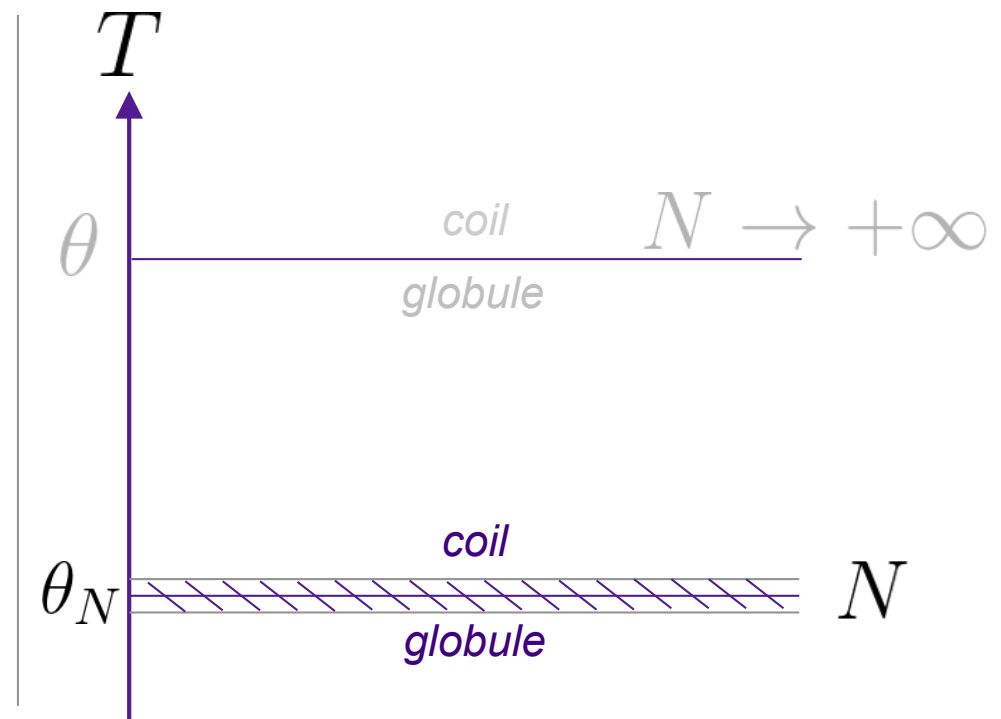
coil-globule transition

depending on the ratio between :

$\epsilon$  interaction energy per monomer

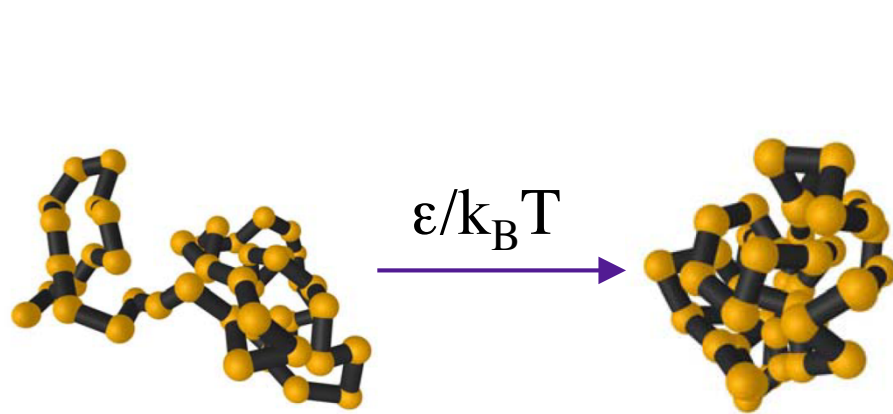
$k_B T$  thermal energy

( $K$  Kuhn length = segment length)



# The coil-globule transition: finite-size effects

A polymer with  $N$  identical monomers:



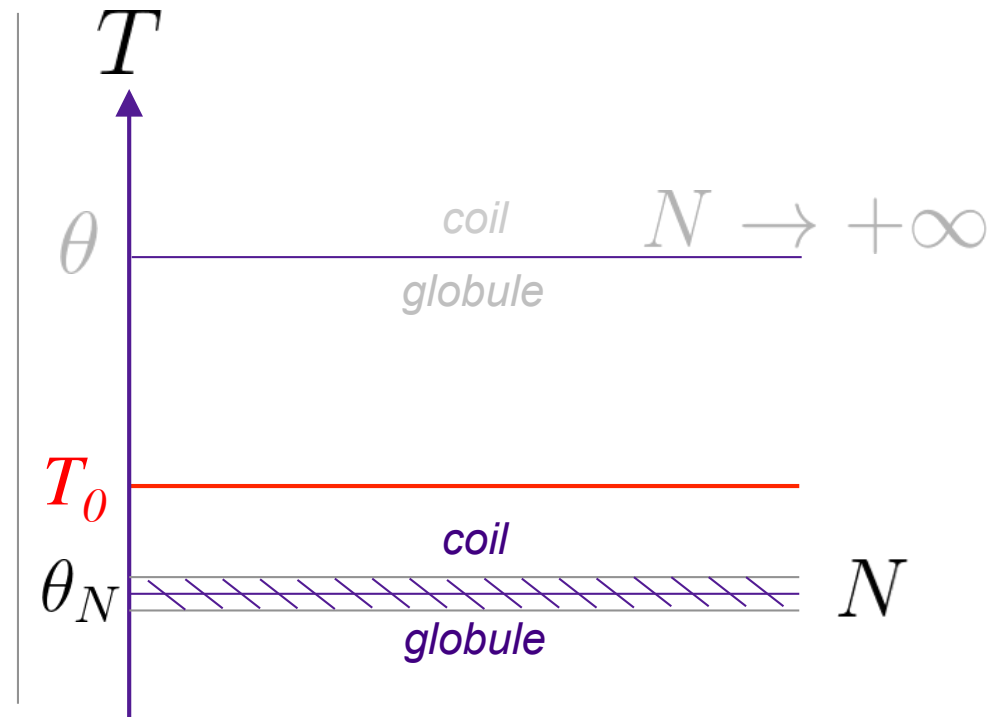
coil-globule transition

depending on the ratio between :

$\epsilon$  interaction energy per monomer

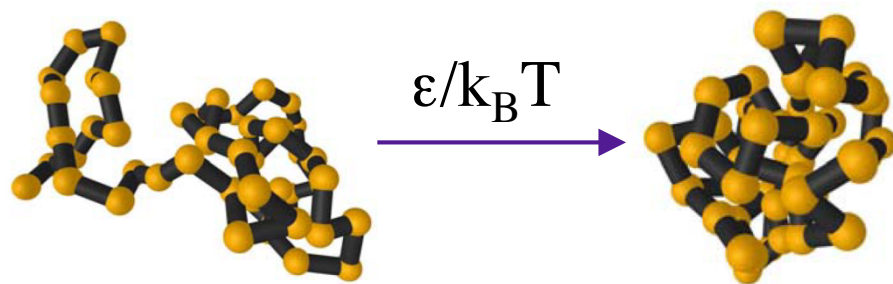
$k_B T$  thermal energy

( $K$  Kuhn length = segment length)



# The coil-globule transition: finite-size effects

A polymer with  $N$  identical monomers:



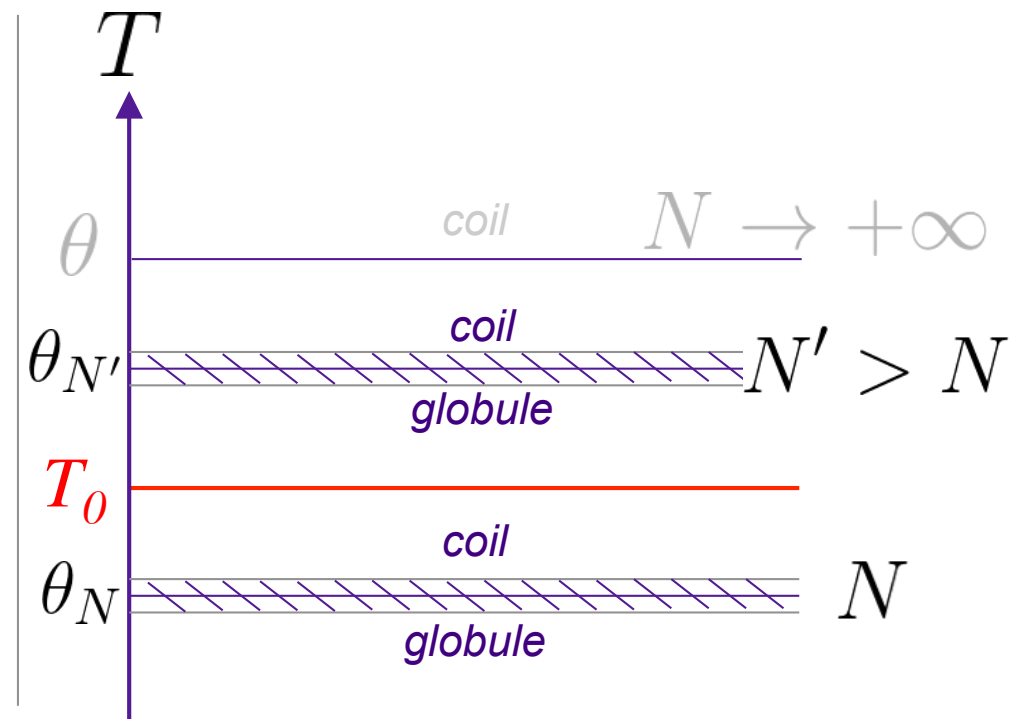
coil-globule transition

depending on the ratio between :

$\epsilon$  interaction energy per monomer

$k_B T$  thermal energy

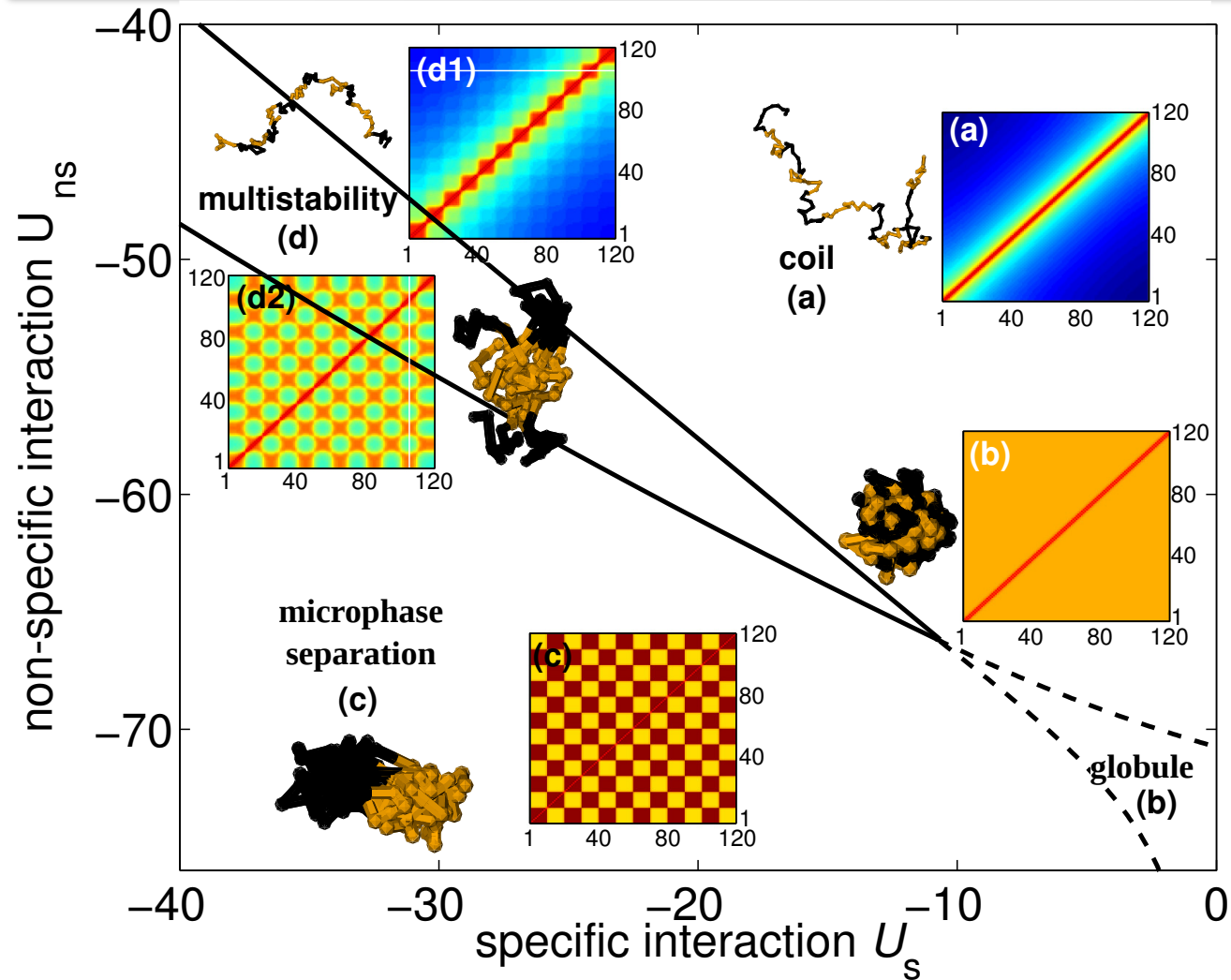
( $K$  Kuhn length = segment length)



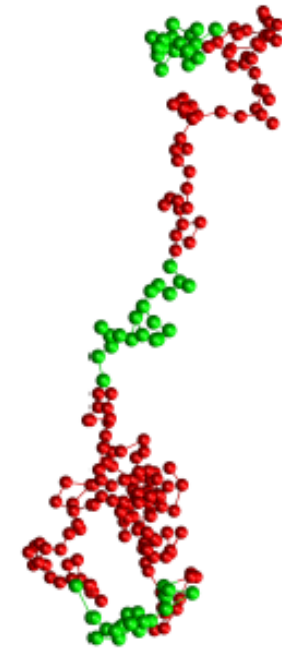
- Small polymers are coil
- Big polymers are globule

# Modeling epigenome folding: formation and dynamics of topologically associated chromatin domains

Daniel Jost<sup>1</sup>, Pascal Carrivain<sup>2</sup>, Giacomo Cavalli<sup>2,\*</sup> and Cédric Vaillant<sup>1,\*</sup>



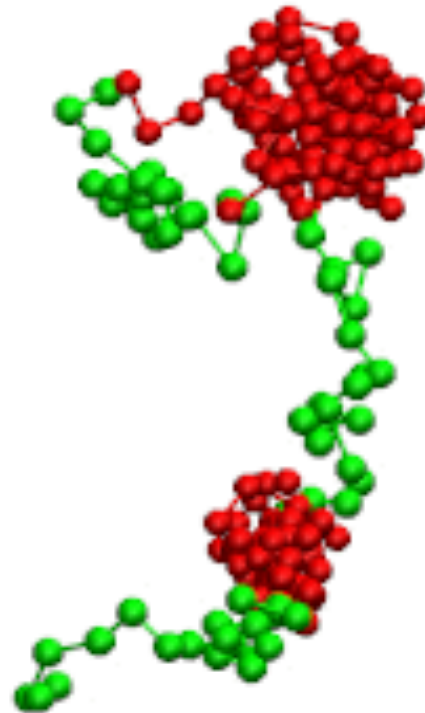
block co-polymer



## Chromatin epigenomic domain folding: size matters

Bertrand R. Caré<sup>1,2,3</sup>, Pierre-Emmanuel Emeriau<sup>1,2,3</sup>, Ruggero Cortini<sup>1,2,3</sup>, Jean-Marc Victor<sup>1,2,3</sup>, , 

20 / 55 ; intra-color interaction 
$$U = \begin{cases} -\varepsilon \left[ 1 - e^{-a(r-r_0)^2} \right] & \text{if } 0 \leq r \leq r_{\max} \\ 0 & \text{if } r > r_{\max} \end{cases}$$

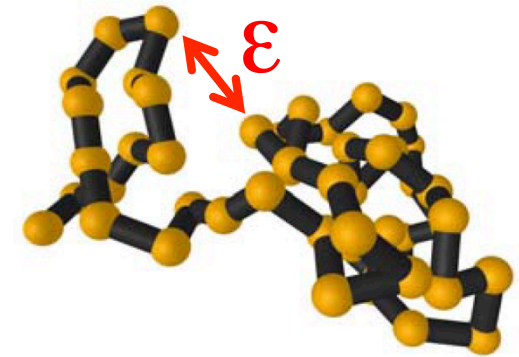


# The coil-globule transition: model

«Interacting Self Avoiding Walk » (ISAW) model

Order parameter:  $\hat{t} = N^{1/2} \left( \frac{N}{r_G^3} \right)^{5/4}$

Dimensionned order parameter:  $s = L^{1/2} \left( \frac{L}{R_G^3} \right)^{5/4}$



Free energy :  $\beta F_L(s) = A_1 N^{1/2} s + A_2 s^2 + A_3 N^{-1/3} s^{-2/3} + 1.13 \ln s$

$A_1$   $A_2$  and  $A_3$  depends on the 3 physical parameters :  
the energy  $\epsilon$ , the Kuhn length  $K$  and the compaction  $\alpha$

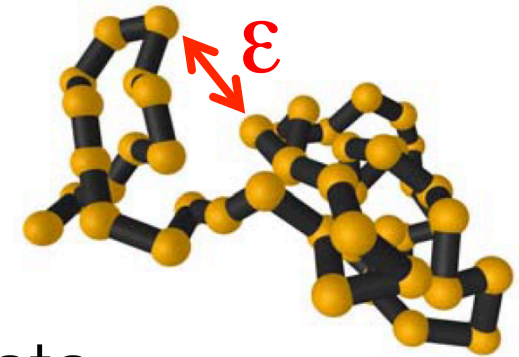
# The coil-globule transition: model

## Working hypothesis:

$\epsilon$ ,  $K$ ,  $\alpha$  are specific for a given state

## Goal:

Fit  $A_1$ ,  $A_2$  and  $A_3$  over the whole set of data,  
for each state



$$\text{Free energy : } \beta F_L(s) = A_1 N^{1/2} s + A_2 s^2 + A_3 N^{-1/3} s^{-2/3} + 1.13 \ln s$$

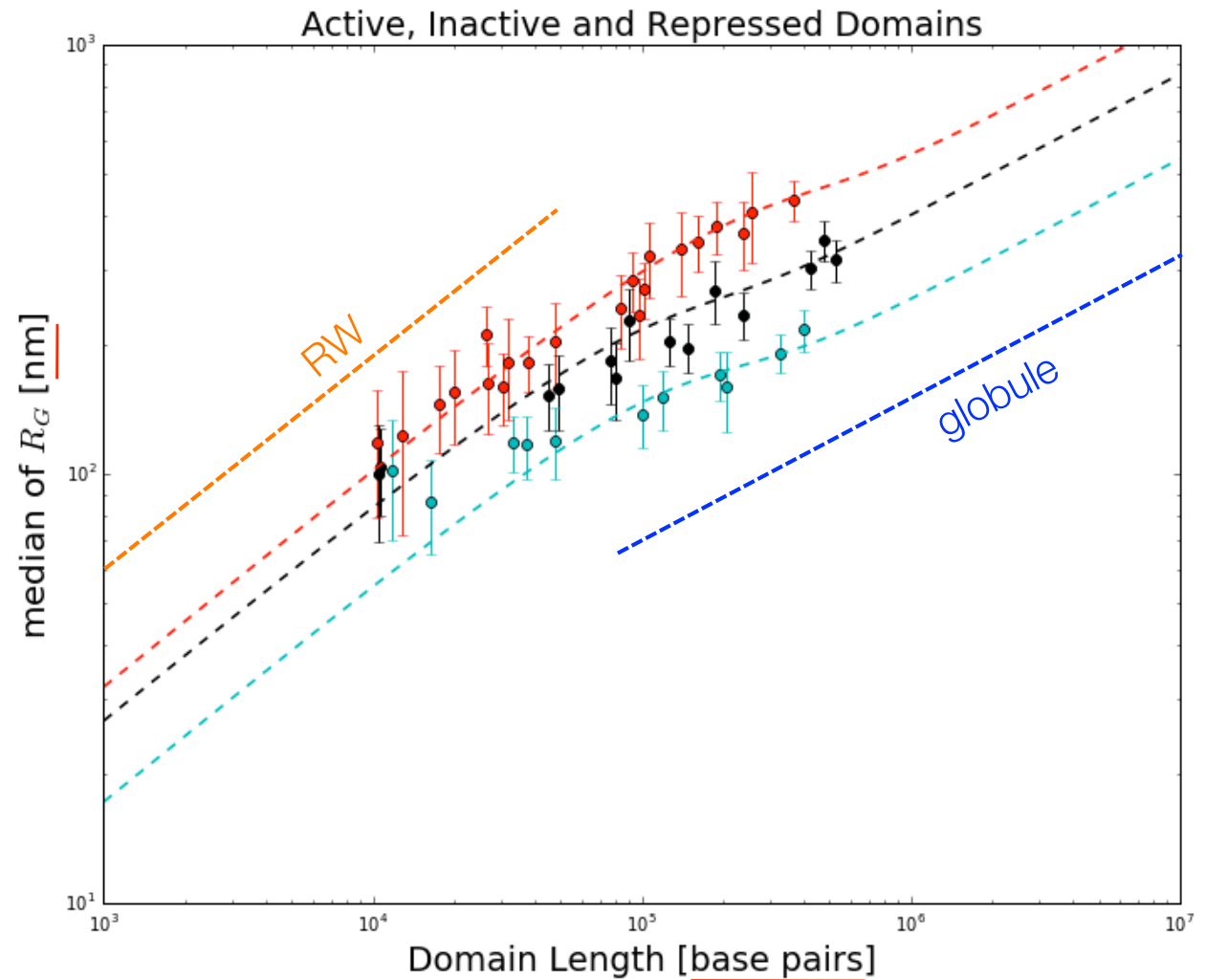
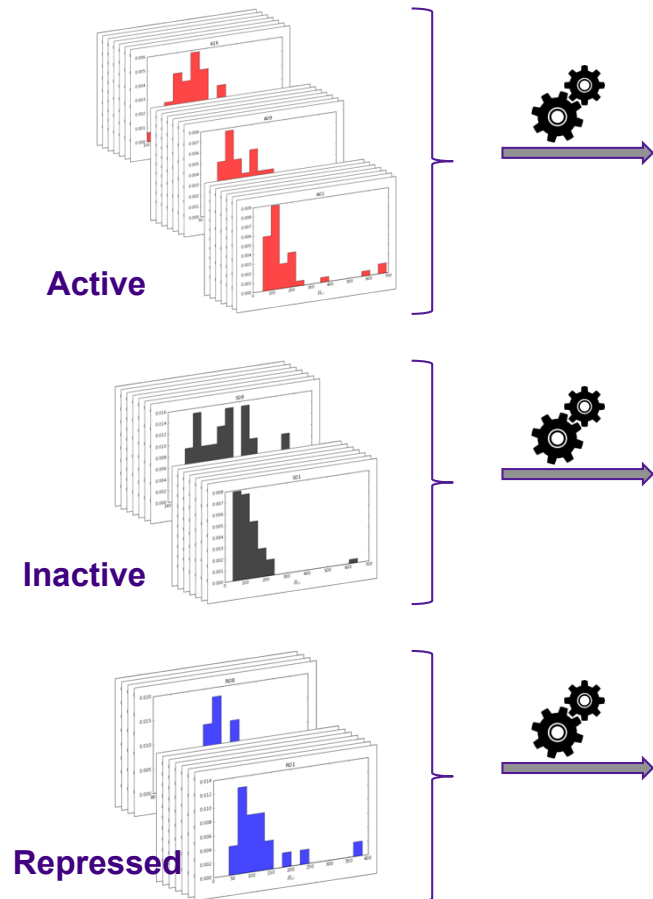
$A_1$ ,  $A_2$  and  $A_3$  depends on the 3 physical parameters :  
the energy  $\epsilon$ , the Kuhn length  $K$  and the compaction  $\alpha$



# Fitting approach and result



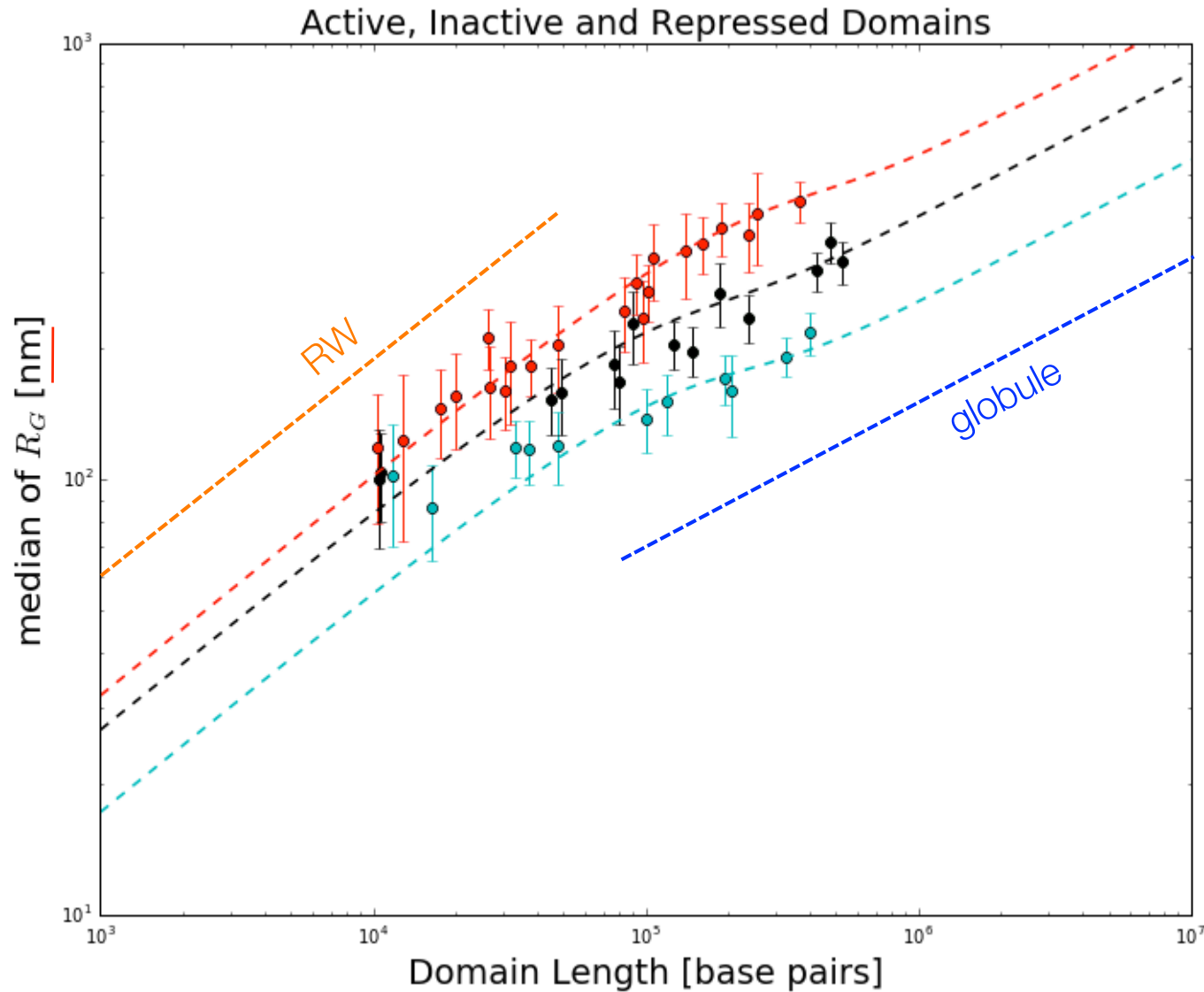
= maximum likelihood + Bayesian inference



(Data (points) from Boettinger et al. 2016)

# Results (0)

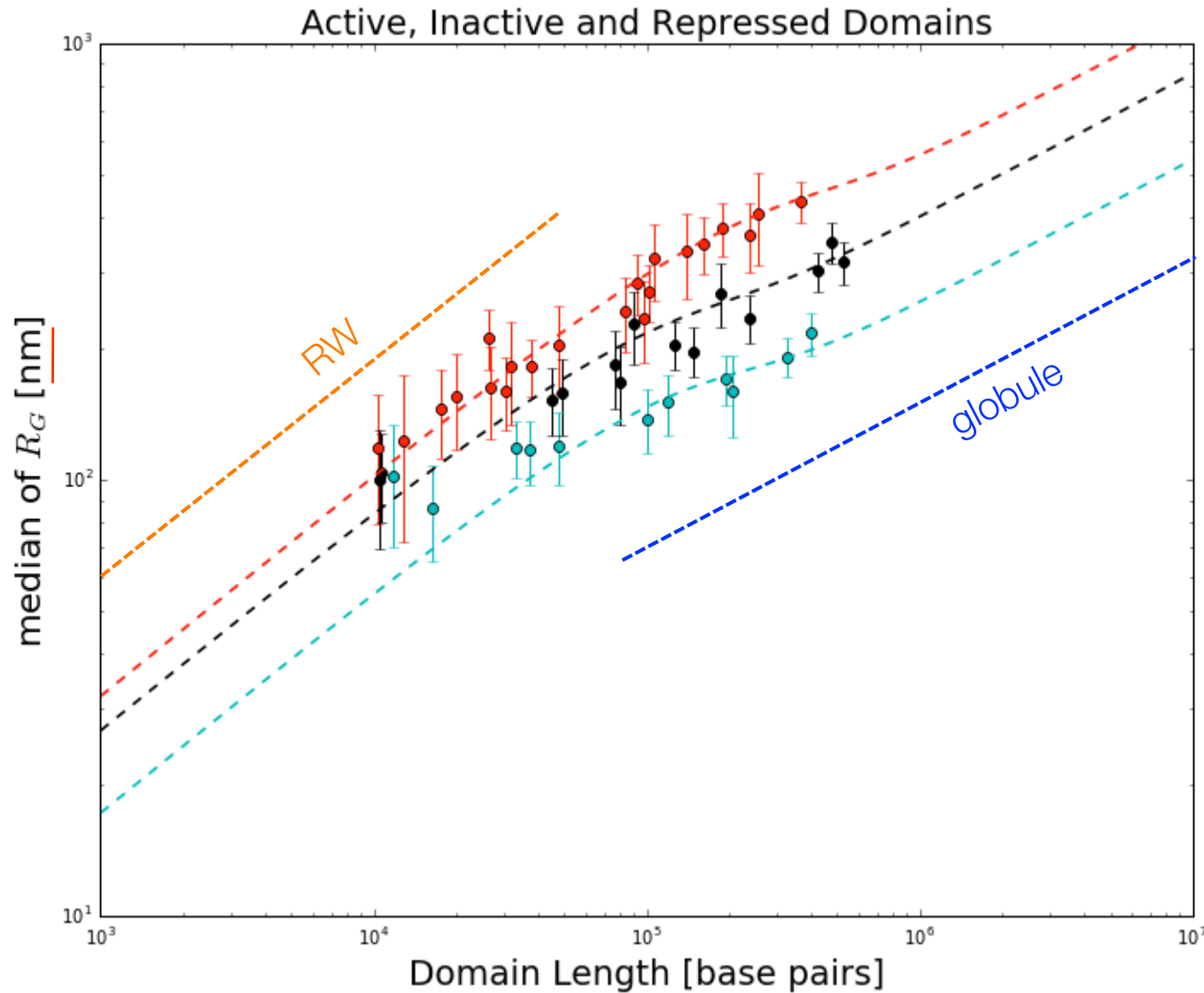
It is possible to fit the data with this model



(Data (points) from Boettinger et al. 2016)

# Results (1)

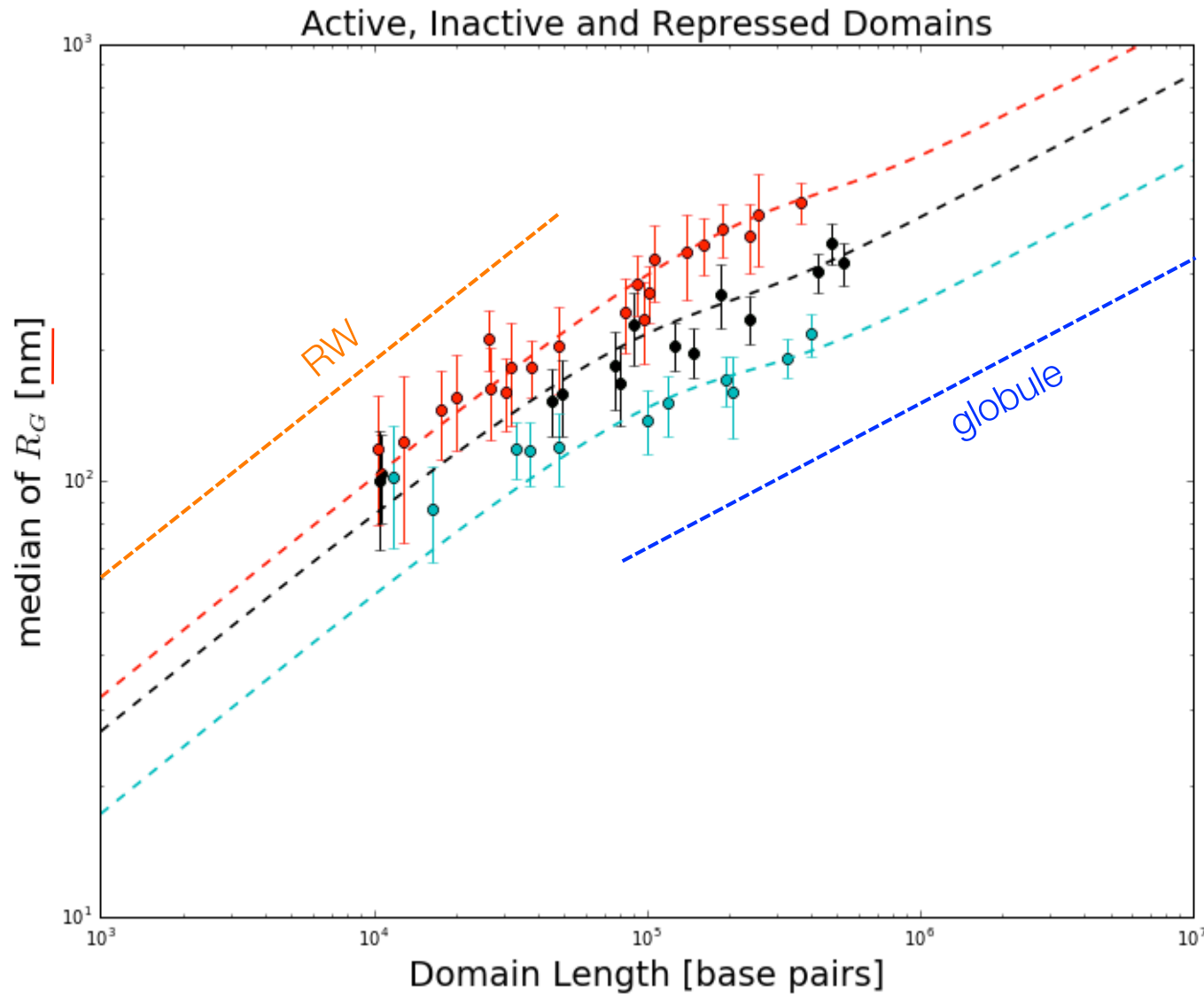
All three states are at the coil-globule transition



(Data (points) from Boettinger et al. 2016)

# Results (2)

Fits give  $\varepsilon \sim 0.3 k_B T$  per Kuhn segment (for all...)




(Data (points) from Boettinger et al. 2016)

## Results (3)

Fits give  $\varepsilon \sim 0.3 k_B T$  per Kuhn segment (for all...)

But the Kuhn length  $K$  is different for different states:



<i>State</i>	$K$ [bp]	$\alpha$ [bp/nm]
<i>Active</i>	$203 \pm 93$	$6.2 \pm 3.1$
<i>Inactive</i>	$154 \pm 50$	$6.5 \pm 2.3$
<i>Repres.</i>	$50 \pm 8$	$5.8 \pm 1.1$

smaller Kuhn length.. (similar compaction)

...hence larger energy/bp from active to repressed

# Results (4)

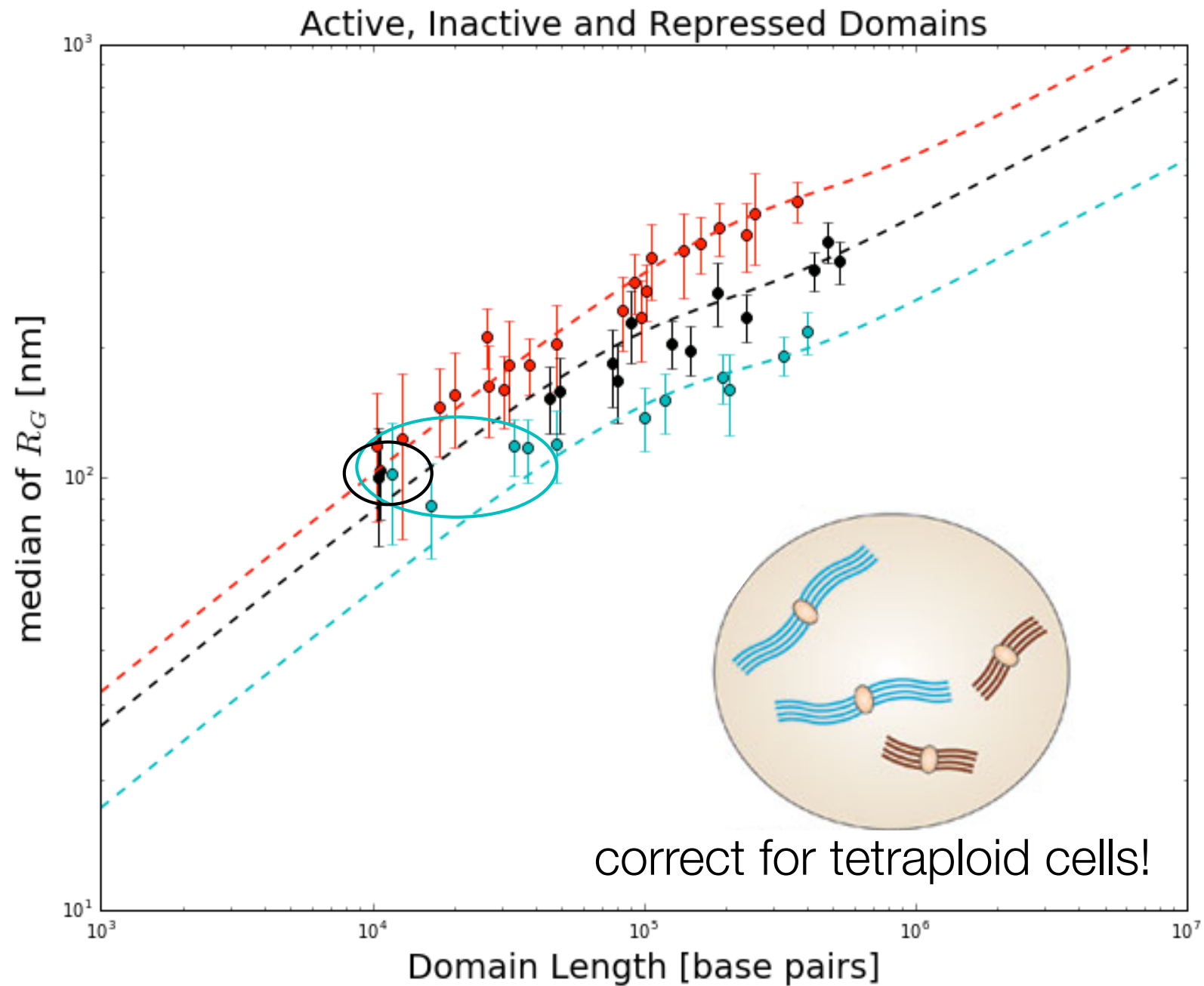
Up to a VERY small persistence length!

$$(K = 2 L_p)$$

<i>State</i>	<i>K</i> [bp]	$\alpha$ [bp/nm]
<i>Active</i>	203 ± 93	6.2 ± 3.1
<i>Inactive</i>	154 ± 50	6.5 ± 2.3
<i>Repres.</i>	50 ± 8	5.8 ± 1.1

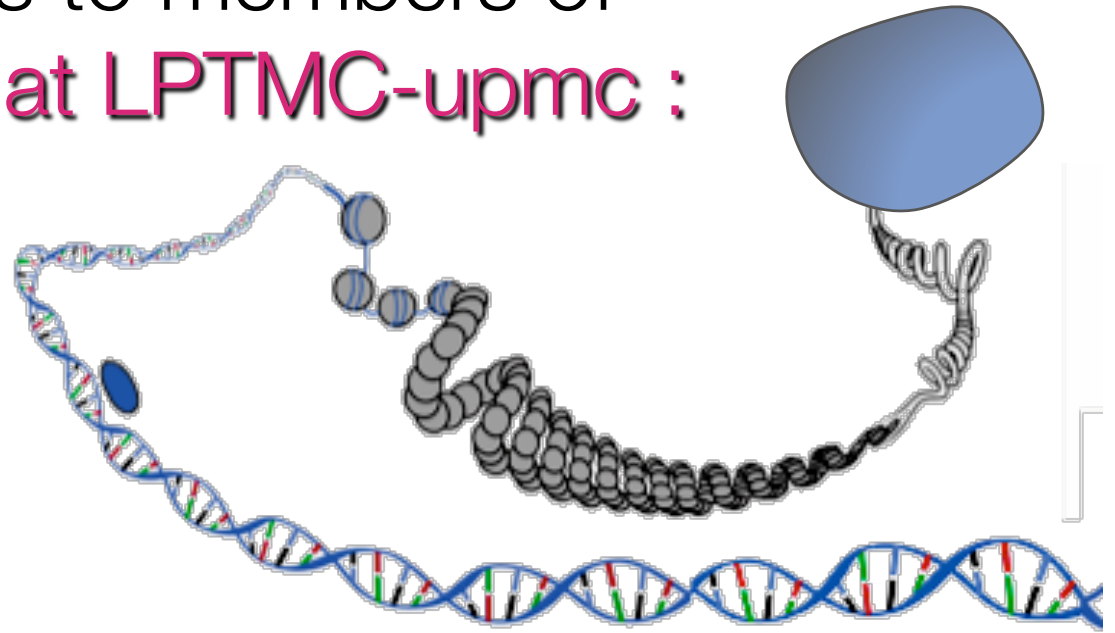
Less than one nucleosome repeat...  
... to be confirmed!

# Work in progress

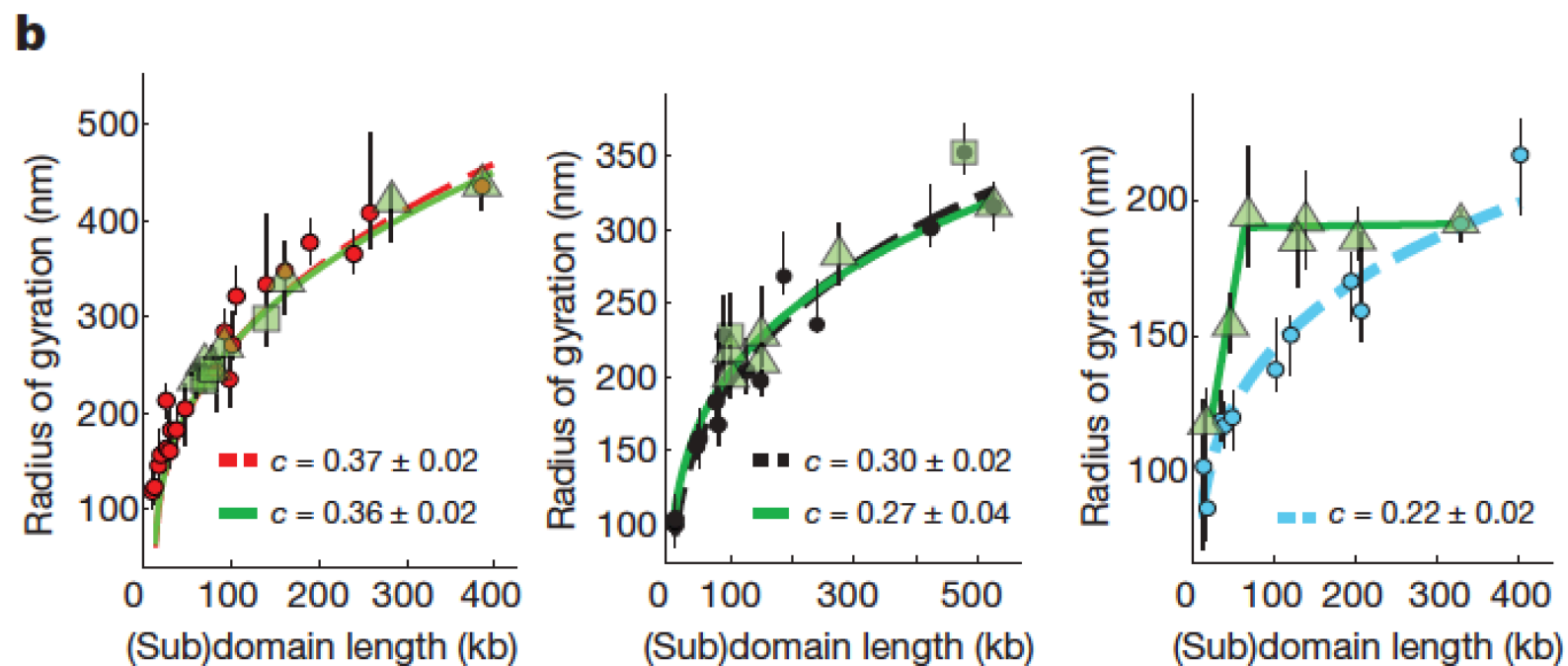
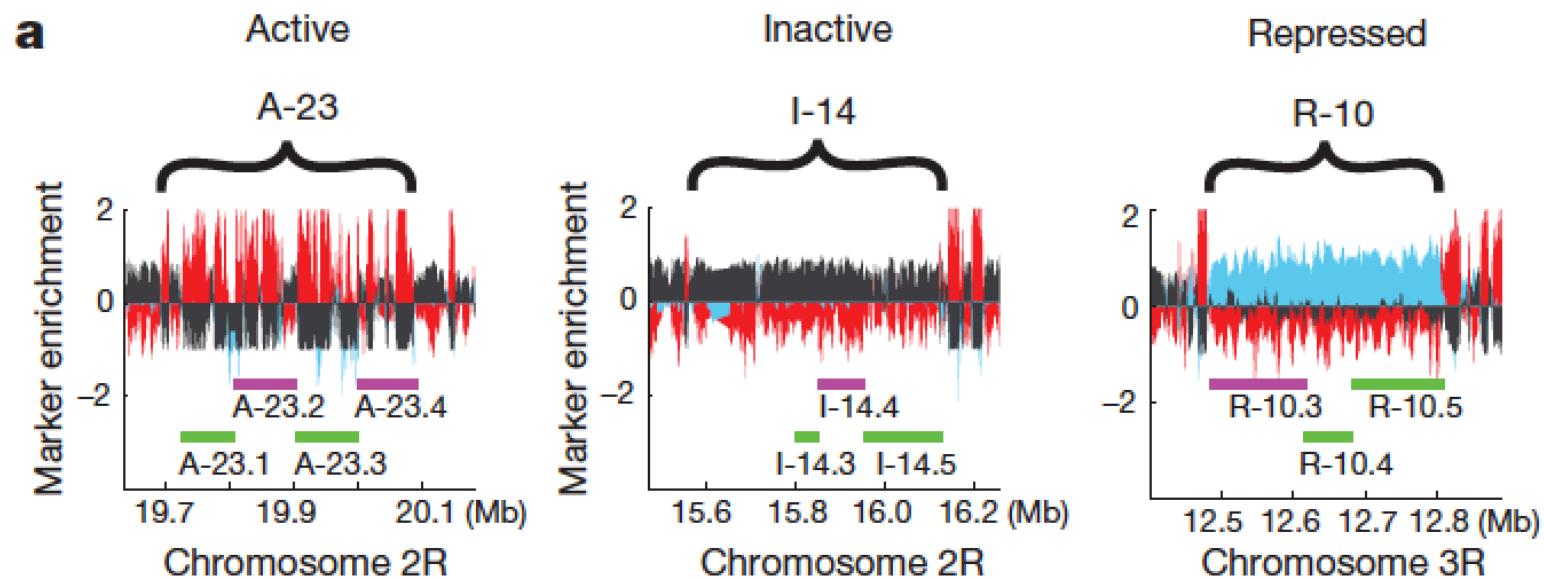


Thank you for your attention,  
and thanks to members of  
**M3V Group at LPTMC-upmc :**

- Jean-Marc Victor
- Annick Lesne
- Julien Mozziconacci
- Bertrand Caré (postdoc ANR/INCa, 2013-2015)
- Ruggero Cortini (postdoc ANR/INCa, 2013-2015)
- **Antony Lesage** (Master student, starting PhD) << who did the work!

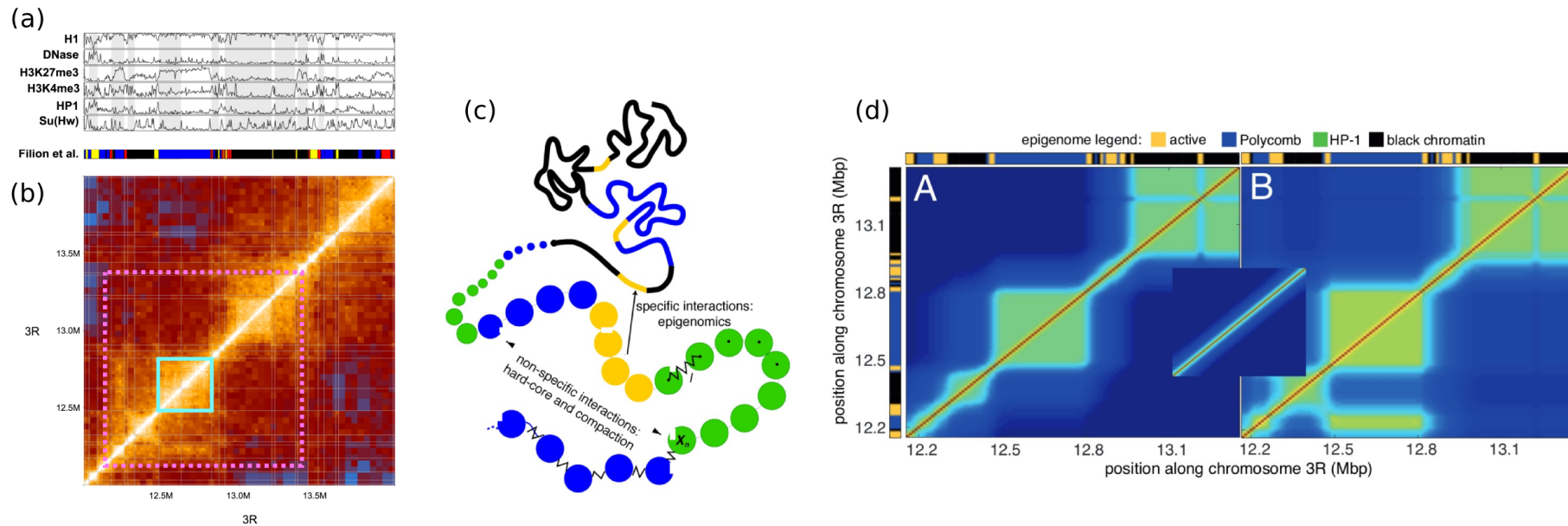






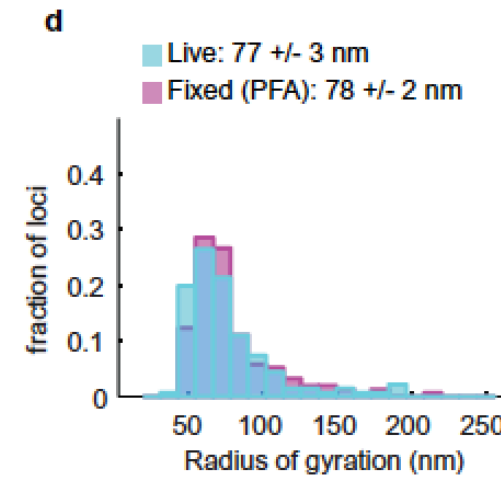
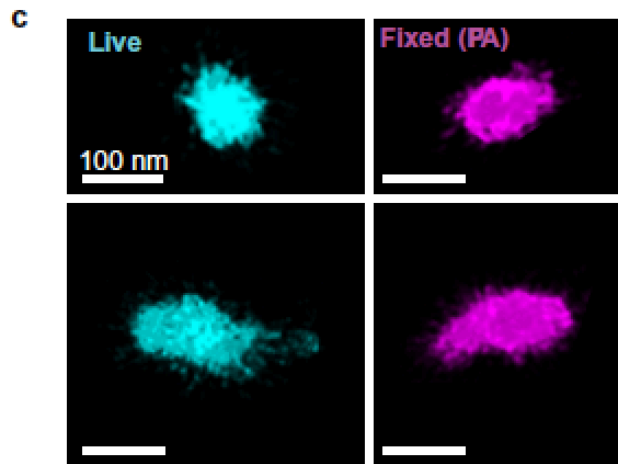
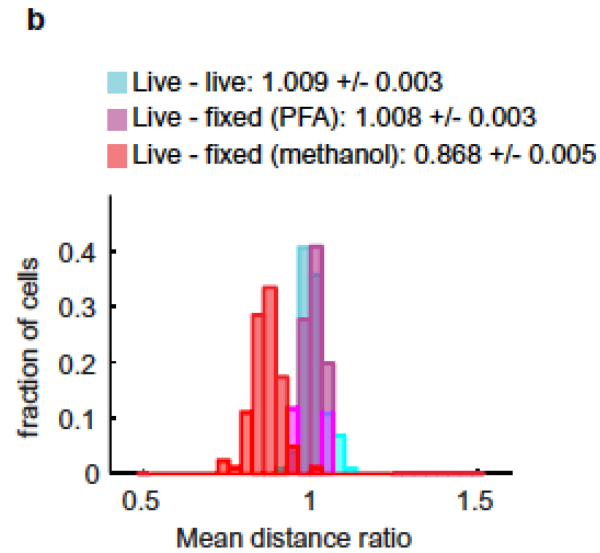
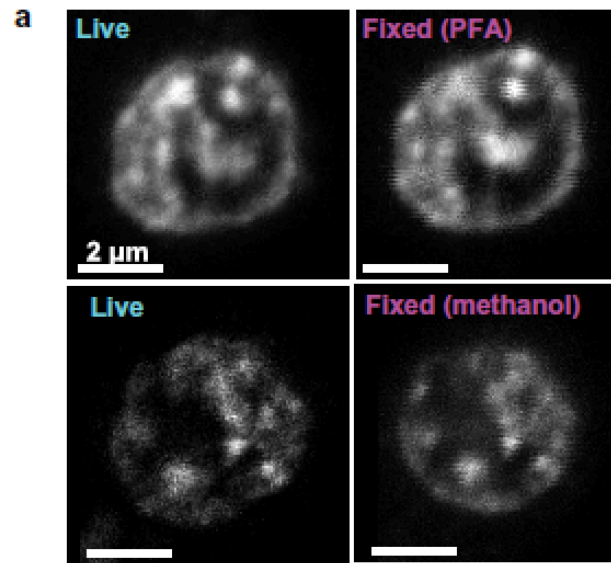
# The drosophila nucleus is a copolymer brush

Jost et al, NAR 2014



$$U_{i,j} = (U_{ns} + U_s) \exp\left(-\frac{r_{i,j}^2}{a^2}\right) \quad \text{if } i \text{ and } j \text{ are of the same color}$$

$$U_{i,j} = U_{ns} \exp\left(-\frac{r_{i,j}^2}{a^2}\right) \quad \text{if } i \text{ and } j \text{ are of different colors}$$



osmotically balanced methanol-free formaldehyde in PBS

# how to interpret the experimental scaling laws?

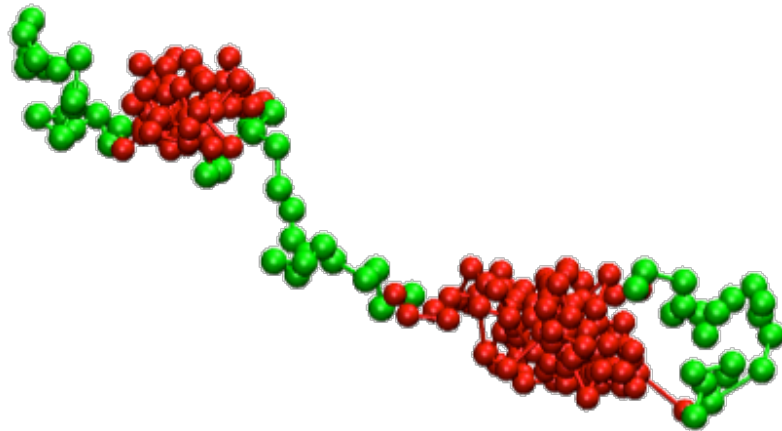
AIMS Biophysics, 2015, 2(4): 517-530. doi: 10.3934/biophy.2015.4.517

Research article **Special Issues**

## Chromatin epigenomic domain folding: size matters

Bertrand R. Caré<sup>1,2,3</sup>, Pierre-Emmanuel Emeriau<sup>1,2,3</sup>, Ruggero Cortini<sup>1,2,3</sup>, Jean-Marc Victor<sup>1,2,3</sup>, , 

co-polymer model



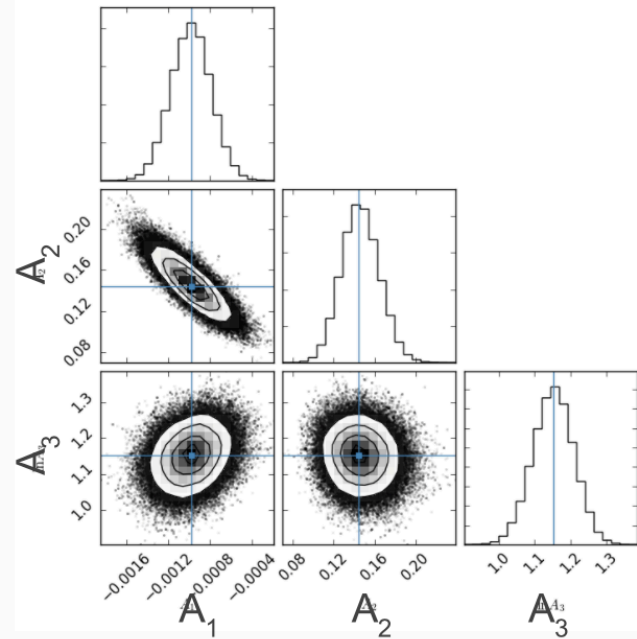
$$\theta(N) = \theta_{\infty} \left( 1 - \frac{c}{N^{1/2}} \right)$$

$$\Delta T(N) \sim \frac{1}{N^{1/2}}$$

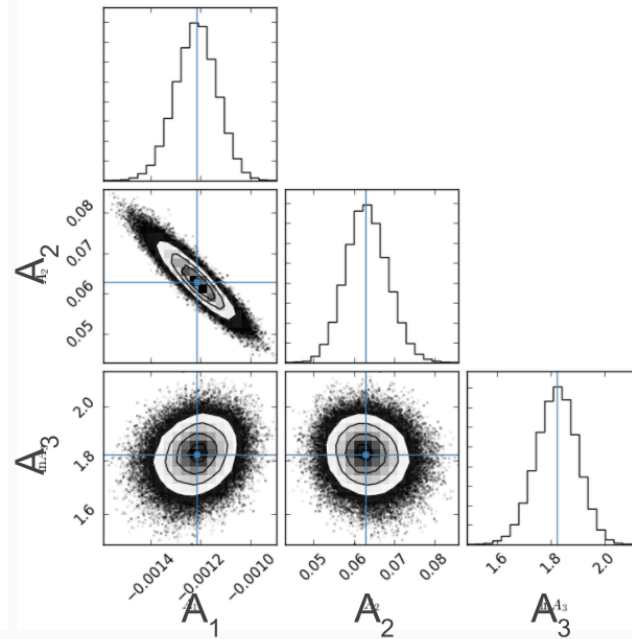
$$U = \begin{cases} -\varepsilon \left[ 1 - e^{-a(r-r_0)^2} \right] & \text{if } 0 \leq r \leq r_{\max} \\ 0 & \text{if } r > r_{\max} \end{cases}$$

« A key prediction of the model is that chromatin compaction should increase with block size for a given epigenomic state (i.e. for a given interaction between beads of chromatin in a block). »

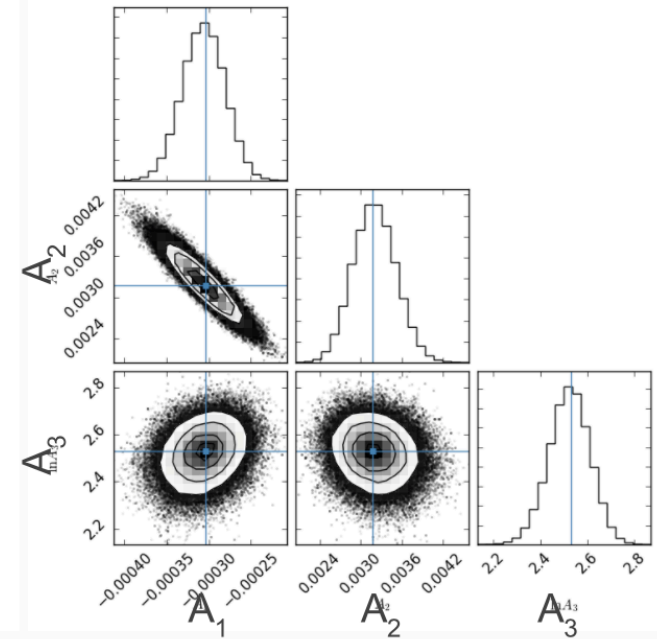
# Bayesian inference



**Active**



**Inactive**



**Repressed**

# coming back to the adimensional model

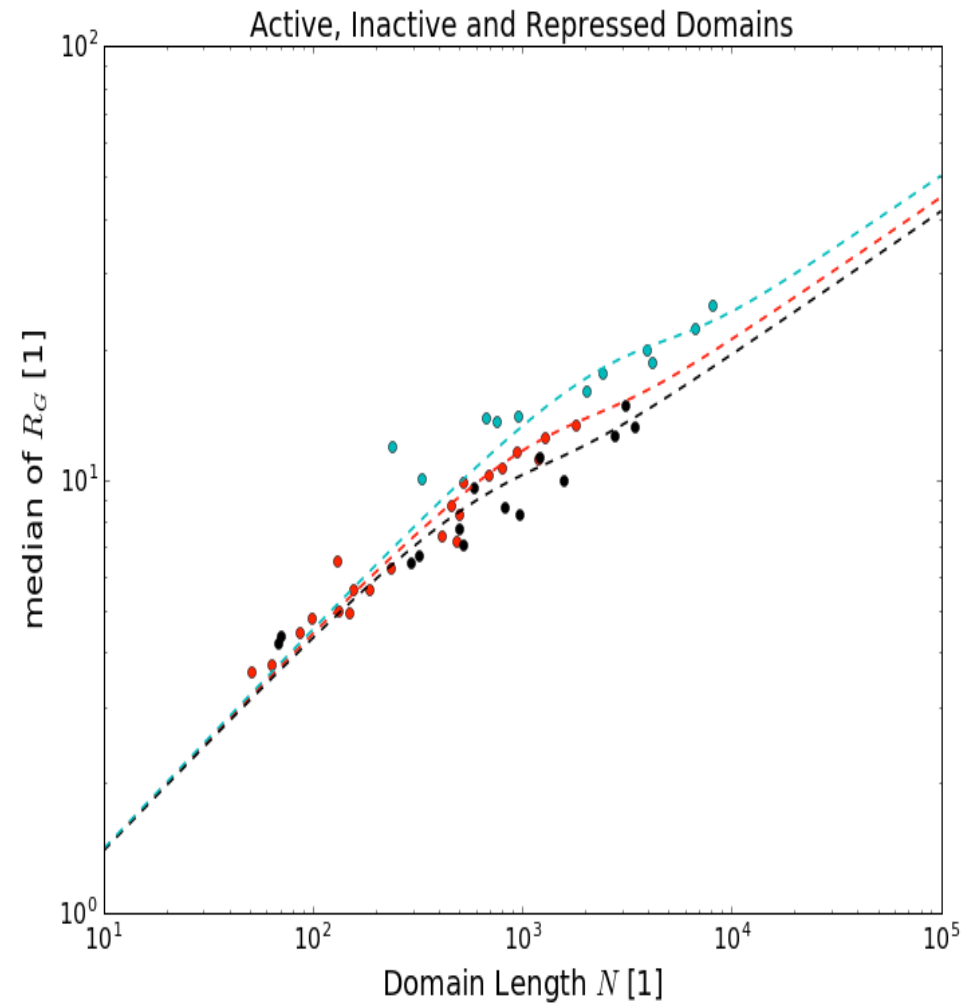
$a_2, a_3$  constants, determined by simulations

$$a_1(U/k_B T)$$

$$a_1 = A_1 \left( \frac{A_3}{a_3} \right)^{\frac{3}{2}} \rightarrow \beta \varepsilon$$

$$K_{bp} = \frac{A_2}{a_2} \left( \frac{A_3}{a_3} \right)^3$$

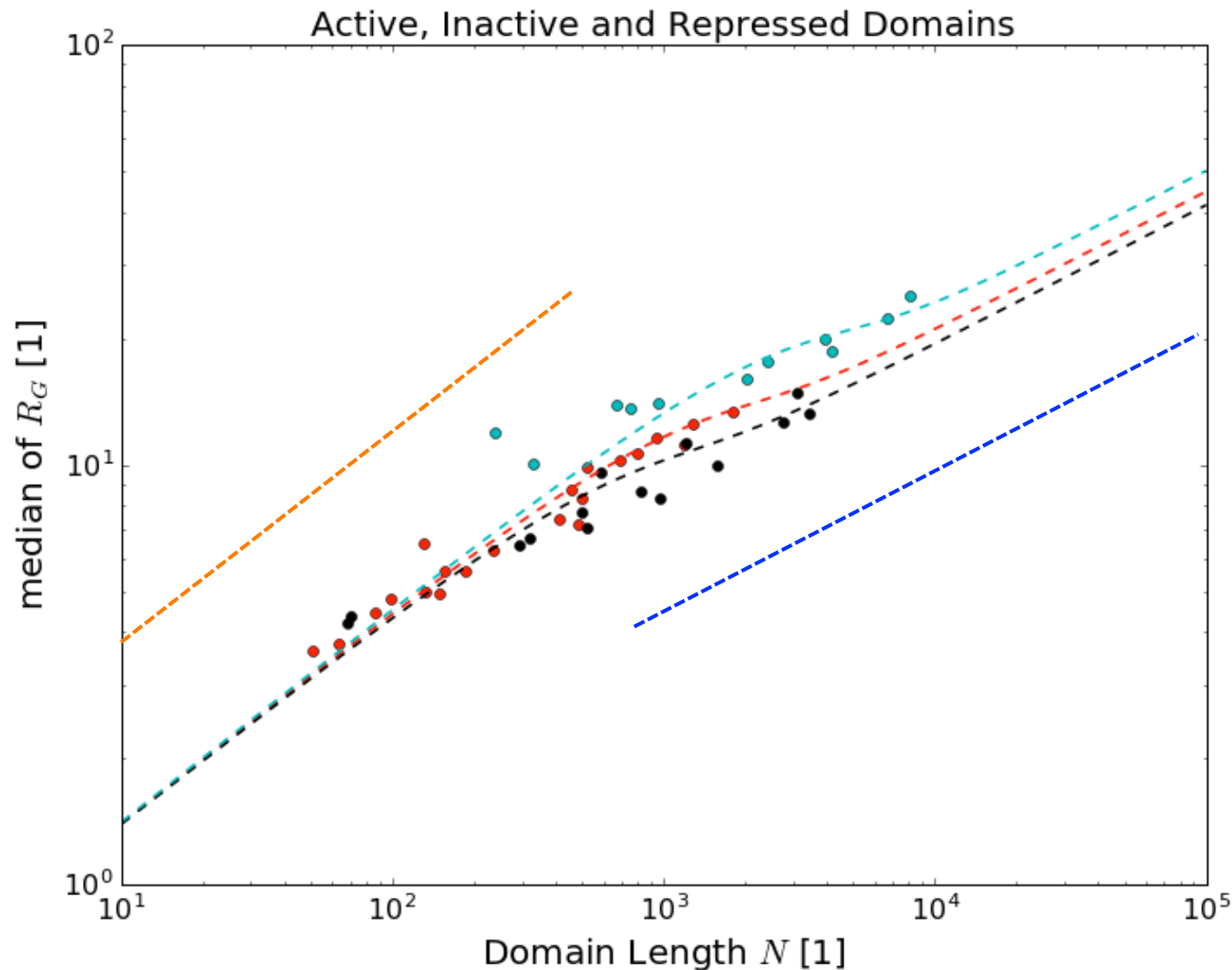
$$K_{nm} = \left( \frac{A_2}{a_2} \right)^{\frac{3}{5}} \left( \frac{A_3}{a_3} \right)^{\frac{7}{5}}$$



# coming back to the adimensional model

physical data → dimensionless quantities

→ can be compared with simulations of model polymers



$$N = \frac{L}{K_{bp}}$$

$$r = \frac{R_G}{K_{nm}}$$