

## Response and correlation functions

- **Linear response**

$$\begin{aligned}\hat{H}_{\text{tot}} &= \hat{H} - \int d^d r \sum_i h_i(\mathbf{r}, t) \hat{A}_i(\mathbf{r}) \\ \delta \langle \hat{A}_i(\mathbf{r}, t) \rangle &= \int_{-\infty}^{\infty} dt' \int d^d r' \sum_j \chi_{A_i A_j}^R(\mathbf{r}, t; \mathbf{r}', t') h_j(\mathbf{r}', t') + \mathcal{O}(h_j^2)\end{aligned}\quad (1)$$

$$\chi_{A_i A_j}^R(\mathbf{r}, t; \mathbf{r}', t') = i\theta(t - t') \langle [\hat{A}_i(\mathbf{r}, t), \hat{A}_j(\mathbf{r}', t')] \rangle \quad (\text{“retarded” response function or susceptibility})$$

where  $\hat{A}_i(\mathbf{r}, t) = e^{i\hat{H}t} \hat{A}_i(\mathbf{r}) e^{-i\hat{H}t}$  (Heisenberg representation) and  $\langle \dots \rangle = \text{Tr}(e^{-\beta\hat{H}} \dots) / \text{Tr} e^{-\beta\hat{H}}$ .

- **External field (adiabatically switched on)**

$$\begin{aligned}h_i(t) &= h_i(\omega) e^{-i(\omega+i\eta)t} + \text{c.c.} \quad (\eta \rightarrow 0^+) \\ \delta \langle \hat{A}_i(t) \rangle &= \sum_j \chi_{A_i A_j}^R(\omega) h_j(\omega) e^{-i\omega t + \eta t} + \text{c.c.} \\ \chi_{A_i A_j}^R(\omega) &= \int_{-\infty}^{\infty} dt \chi_{A_i A_j}^R(t) e^{i(\omega+i\eta)t}\end{aligned}\quad (2)$$

(To alleviate notations, we include the  $\mathbf{r}$  dependence in the  $i$  index.)

- 

$$\begin{aligned}\chi_{A_i A_j}^R(t; t') &= 2i\theta(t - t') \chi_{A_i A_j}''(t; t') \\ \chi_{A_i A_j}''(t; t') &= \frac{1}{2} \langle [\hat{A}_i(t), \hat{A}_j(t')] \rangle\end{aligned}\quad (3)$$

- **Dissipation**

$$\begin{aligned}\frac{d\overline{W}}{dt} &= \frac{1}{2} \sum_{i,j} h_i^*(\omega) \omega \chi_{A_i^\dagger A_j}''(\omega) h_j(\omega) \quad (\geq 0) \\ \{\omega \chi_{A_i^\dagger A_j}''(\omega)\} &\text{ positive matrix} \\ \omega \chi_{A_i^\dagger A}''(\omega) &\geq 0\end{aligned}\quad (4)$$

- **Spectral representation**

$$\begin{aligned}\chi_{A_i A_j}(z) &= \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\chi_{A_i A_j}''(\omega')}{\omega' - z} \quad (z \text{ complex frequency}) \\ \chi_{A_i A_j}^R(\omega) &= \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\chi_{A_i A_j}''(\omega')}{\omega' - \omega - i\eta} = \chi_{A_i A_j}(z = \omega + i\eta) \\ \chi_{A_i A_j}''(\omega) &= \frac{\pi}{Z} \sum_{n,m} (e^{-\beta\epsilon_n} - e^{-\beta\epsilon_m}) A_i^{nm} A_j^{mn} \delta(\omega + \epsilon_n - \epsilon_m) \quad (\text{spectral function})\end{aligned}\quad (5)$$

where  $\{|n\rangle, \epsilon_n\}$  form a complete set of eigenvectors of  $\hat{H}$  and  $A_i^{nm} = \langle n | \hat{A}_i | m \rangle$ .  $\chi_{A_i A_j}(z)$  is analytic in the whole complex plane except (possibly) on the real axis.

- $\chi''_{A^\dagger A}(\omega) = \Im[\chi^R_{A^\dagger A}(\omega)]$  is real.

- **Lehmann representation**

$$\chi_{A_i A_j}(z) = \frac{1}{Z} \sum_{\substack{n,m \\ (\epsilon_n \neq \epsilon_m)}} \frac{e^{-\beta\epsilon_n} - e^{-\beta\epsilon_m}}{\epsilon_m - \epsilon_n - z} A_i^{nm} A_j^{mn} \quad (6)$$

- **Causality and Kramers-Kronig relations**

$$\begin{aligned} \Re[\chi^R(\omega)] &= \mathcal{P} \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\Im[\chi^R(\omega')]}{\omega' - \omega} \\ \Im[\chi^R(\omega)] &= -\mathcal{P} \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\Re[\chi^R(\omega')]}{\omega' - \omega} \end{aligned} \quad (7)$$

- **Fluctuation-dissipation theorem**

$$\begin{aligned} S_{A_i A_j}(t) &= \langle \hat{A}_i(t) \hat{A}_j \rangle - \langle \hat{A}_i \rangle \langle \hat{A}_j \rangle \quad (\text{structure factor}) \\ \chi''_{A_i A_j}(\omega) &= \frac{1}{2} (1 - e^{-\beta\omega}) S_{A_i A_j}(\omega) \\ \chi''_{A_i A_j}(\omega) &= \frac{\beta\omega}{2} S_{A_i A_j}(\omega) \quad \text{i.e.} \quad \chi_{A_i A_j}(t=0) = \beta S_{A_i A_j}(t=0) \quad \text{for } |\omega| \ll T \quad (\text{classical limit}) \end{aligned} \quad (8)$$

- **High-frequency limit**

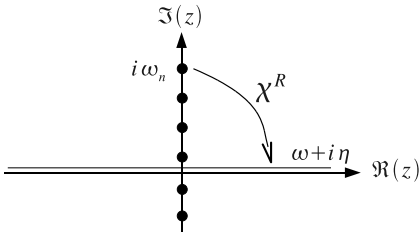
$$\begin{aligned} \chi_{A_i A_j}(z) &= -\frac{1}{z} \int_{-\infty}^{\infty} \frac{d\omega}{\pi} \chi''_{A_i A_j}(\omega) - \frac{1}{z^2} \int_{-\infty}^{\infty} \frac{d\omega}{\pi} \omega \chi''_{A_i A_j}(\omega) + \mathcal{O}(z^{-3}) \\ &= -\frac{1}{z} \langle [\hat{A}_i, \hat{A}_j] \rangle - \frac{1}{z^2} \langle [[\hat{A}_i, \hat{H}], \hat{A}_j] \rangle + \mathcal{O}(z^{-3}) \end{aligned} \quad (9)$$

- **Imaginary-time correlation functions**

$$\begin{aligned} \chi_{A_i A_j}(\tau) &= \langle T_\tau \hat{A}_i(\tau) \hat{A}_j \rangle - \langle \hat{A}_i \rangle \langle \hat{A}_j \rangle \\ \chi_{A_i A_j}(i\omega_n) &= \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\chi''_{A_i A_j}(\omega')}{\omega' - i\omega_n} = \chi_{A_i A_j}(z = i\omega_n) \end{aligned} \quad (10)$$

(the second equation is not always true for  $i\omega_n = 0$ ) with  $\hat{A}_i(\tau) = e^{\tau\hat{H}} \hat{A}_i e^{-\tau\hat{H}}$ .

- **Analytic continuation**

$$\chi^R_{A_i A_j}(\omega) = \chi_{A_i A_j}(i\omega_n \rightarrow \omega + i\eta)$$


- **One-particle Green function**

$$\begin{aligned} G(\mathbf{k}, \tau) &= -\langle T_\tau \hat{\psi}(\mathbf{k}, \tau) \hat{\psi}^\dagger(\mathbf{k}, 0) \rangle \\ G(\mathbf{k}, i\omega_n) &= \int_{-\infty}^{\infty} d\omega \frac{A(\mathbf{k}, \omega)}{i\omega_n - \omega} \\ A(\mathbf{k}, t) &= \frac{1}{2\pi} \langle [\hat{\psi}(\mathbf{k}, t), \hat{\psi}^\dagger(\mathbf{k}, 0)]_{-\zeta} \rangle \\ A(\mathbf{k}, \omega) &= -\frac{1}{\pi} \Im[G^R(\mathbf{k}, \omega)] \end{aligned} \quad (11)$$

$\zeta = 1$  (bosons) or  $-1$  (fermions).